

Density functional theory & practice

1. Introduction & Perspective:

electronic structure methods
Born-Oppenheimer approximation
units

2. Theorems:

HK: Density as the basic variable
KS: mapping to single-electron problem

3. Practice:

practical density functionals
self-consistency
basis sets
large systems: $\mathcal{O}(N)$

4. Applications:

ground state: energy, structure
forces: molecular dynamics
excited states?
parameters for correlated systems

References

W. Kohn:

Electronic structure of matter —
wave functions and density functionals
Rev. Mod. Phys. **71**, 1253-1266 (1998)

K. Capelle:

A bird's eye view of density-functional theory
[cond-mat/0211443](https://arxiv.org/abs/cond-mat/0211443)

U. von Barth:

Basic Density-Functional Theory – An Overview
www.fysik4.fysik.uu.se/~thor/vonBarth_lecture.pdf

K. Burke: The ABC of DFT

dft.rutgers.edu/kieron/beta

R.M. Martin: Electronic Structure
Basic Theory and Practical Methods
Cambridge University Press, 2003

R.M. Dreizler and E.K.U. Gross:
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Springer, 1990

U. Scherz: Quantenmechanik (Kap. 7–9)
Teubner, 1999

www.mpi-stuttgart.mpg.de/andersen/users/koch/DFT.html

$$\begin{aligned}
\hat{H} = & - \frac{\hbar^2}{2m_e} \sum_i \nabla_i^2 + \frac{1}{2} \sum_{i \neq j} \frac{e^2}{4\pi\epsilon_0 |\mathbf{r}_i - \mathbf{r}_j|} \\
& + \sum_{i,I} \frac{Z_I e^2}{4\pi\epsilon_0 |\mathbf{r}_i - \mathbf{R}_I|} \\
& - \sum_I \frac{\hbar^2}{2M_I} \nabla_I^2 + \frac{1}{2} \sum_{I \neq J} \frac{Z_I Z_J e^2}{4\pi\epsilon_0 |\mathbf{R}_I - \mathbf{R}_J|}
\end{aligned}$$

The underlying laws necessary for the mathematical theory of a large part of physics and the whole of chemistry are thus completely known, and the difficulty is only that exact applications of these laws lead to equations which are too complicated to be soluble.

P.M.A. Dirac, Proc. Roy. Soc. **A 123**, 714 (1929)

More is different

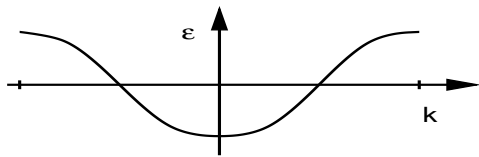
[...] the reductionist hypothesis does not by any means imply a “constructionist” one: The ability to reduce everything to simple fundamental laws does not imply the ability to start from those laws and reconstruct the universe. Sometimes, as in the case of superconductivity, the new symmetry — now called broken symmetry because the original symmetry is no longer evident — may be of an entirely unexpected kind and extremely difficult to visualize. In the case of superconductivity, 30 years elapsed between the time when physicists were in possession of every fundamental law necessary for explaining it and the time when it was actually done.

P.W. Anderson, *Science* **177**, 393 (1972)

Mott transition

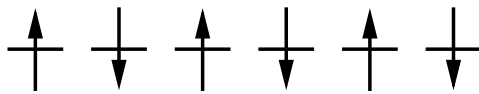
$$\begin{aligned} H &= -t \sum_{\langle ij \rangle \sigma} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow} \\ &= \sum_{k\sigma} \varepsilon_k c_{k\sigma}^\dagger c_{k\sigma} + \frac{U}{M} \sum_{k,k',q} c_{k\uparrow}^\dagger c_{k-q\uparrow} c_{k'\downarrow}^\dagger c_{k'+q\downarrow} \end{aligned}$$

$t \gg U$: band-limit



filling of a band
 \rightsquigarrow metal

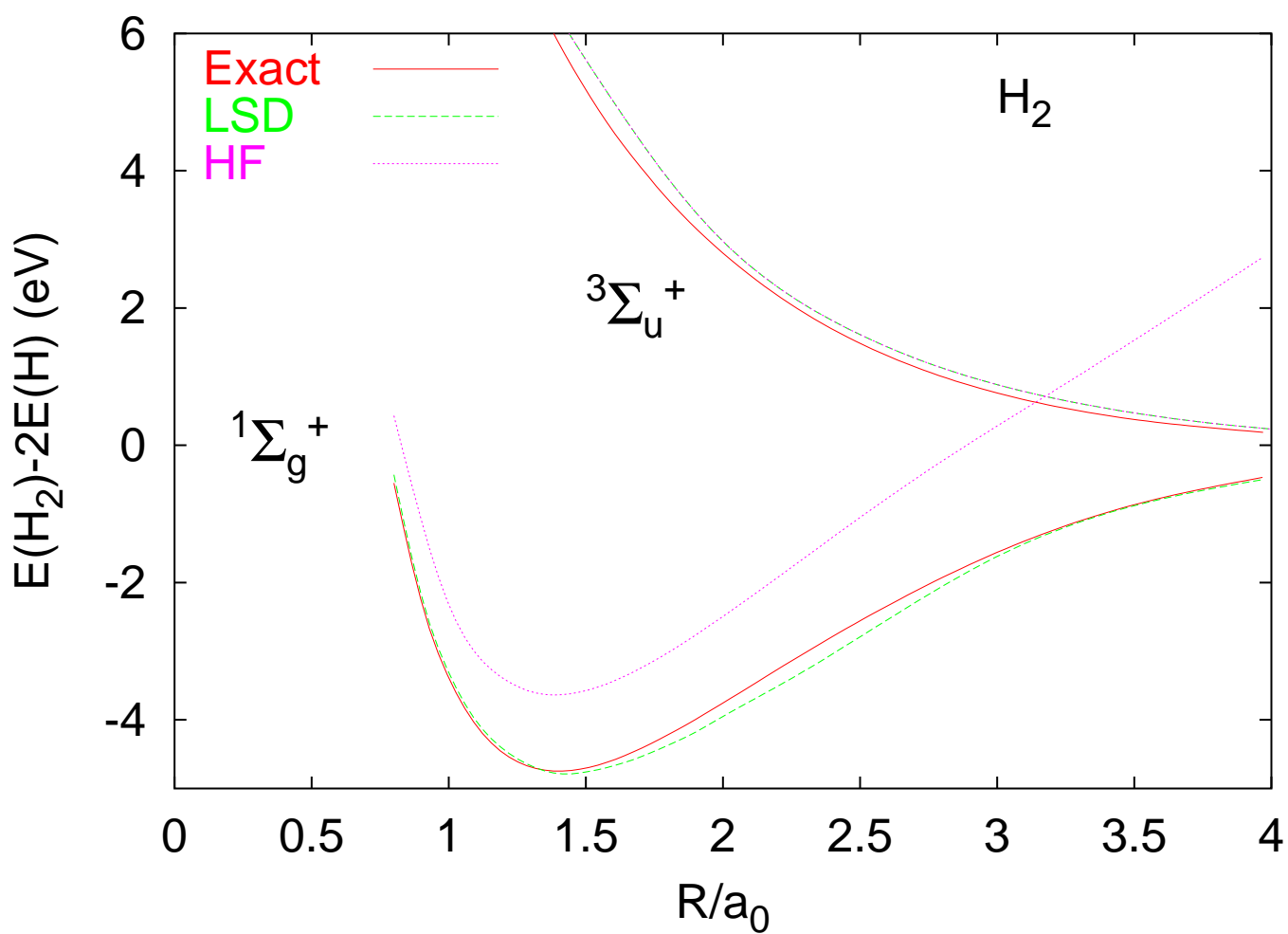
$t \ll U$: atomic limit



no hopping for
integer filling
 \rightsquigarrow insulator

metal-insulator transition for integer fillings

Born-Oppenheimer surfaces for H₂



atomic units

$$\begin{aligned}
 \hbar &= 1.0546 \cdot 10^{-34} \text{ Js} & [ML^2T^{-1}] \\
 m_e &= 9.1094 \cdot 10^{-31} \text{ kg} & [M] \\
 e &= 1.6022 \cdot 10^{-19} \text{ C} & [Q] \\
 4\pi\epsilon_0 &= 1.1127 \cdot 10^{-10} \text{ F/m} & [M^{-1}L^{-3}T^2Q^2]
 \end{aligned}$$

<http://physics.nist.gov/cuu/>

$$\begin{array}{l}
 \text{solve} \\
 \hbar = 1 a_0^2 m_e / t_0 \\
 m_e = 1 m_e \\
 e = 1 e \\
 4\pi\epsilon_0 = 1 t_0^2 e^2 / a_0^3 m_e
 \end{array}
 \quad \text{to obtain}$$

$$\begin{aligned}
 1 \text{ a.u. length} &= a_0 = \frac{4\pi\epsilon_0 \hbar^2}{m_e e^2} \approx 5.2918 \cdot 10^{-11} \text{ m} \\
 1 \text{ a.u. mass} &= m_e = \approx 9.1095 \cdot 10^{-31} \text{ kg} \\
 1 \text{ a.u. time} &= t_0 = \frac{(4\pi\epsilon_0)^2 \hbar^3}{m_e e^4} \approx 2.4189 \cdot 10^{-17} \text{ s} \\
 1 \text{ a.u. charge} &= e = \approx 1.6022 \cdot 10^{-19} \text{ C}
 \end{aligned}$$

$$\Psi(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N)$$

The tabulation of one variable requires a page, of two variables a volume, and of three variables a library; but the full specification of a single wavefunction of neutral Fe is a function of seventy eight variables. It would be rather crude to restrict to ten the number of values at which to tabulate this function, but even so, full tabulation of it would require 10^{78} entries, and even if this number could be reduced somewhat from considerations of symmetry, there would still not be enough atoms in the whole solar system to provide the material for printing such a table.

D.R. Hartree