

Beyond DMFT: Spin fluctuations, Pseudogaps and Superconductivity

Karsten Held (TU Wien)

Jülich, Oct. 7th 2022

- Motivation
- Synopsis: **dynamical vertex approximation (DΓA)**
- 2D Hubbard model, cuprates and nickelates
- Spin fluctuations
- Pseudogap
- Superconductivity

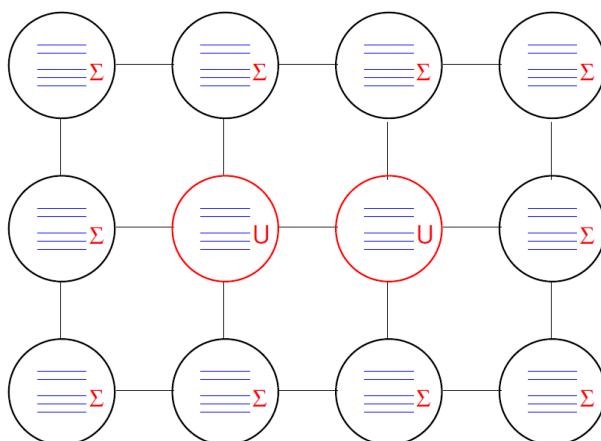
Please ask questions

Correlations beyond DMFT:

d-, p-wave superconductivity, pseudogaps,
(para-)magnons, quantum criticality



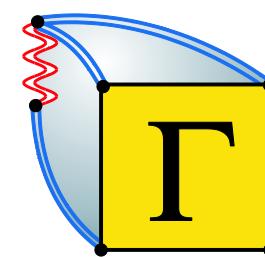
cluster extensions



Hettler et al.'98, Lichtenstein, Katsnelson'00

Kotliar et al.'01, Potthoff et al.'03, Maier et al.'05

diagrammatic extensions



D Γ A: Toschi, Katanin, KH'07

DF: Rubtsov et al.'08

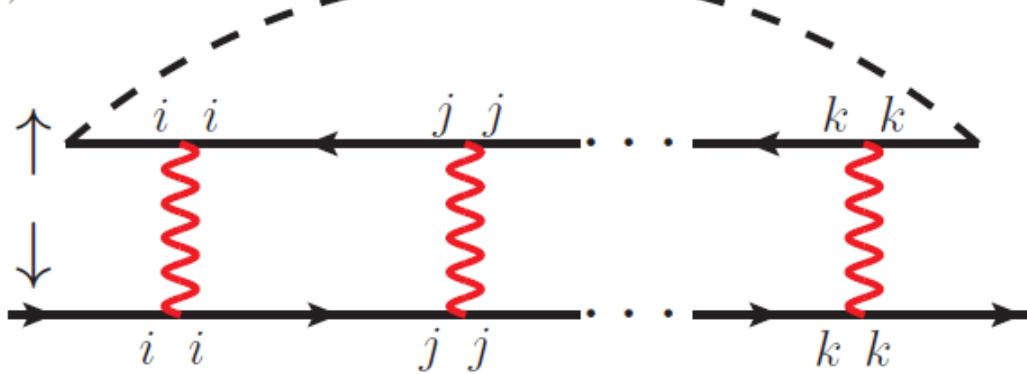
1PI, TRILEX ...

RMP'18

Take spin fluctuations

(a)

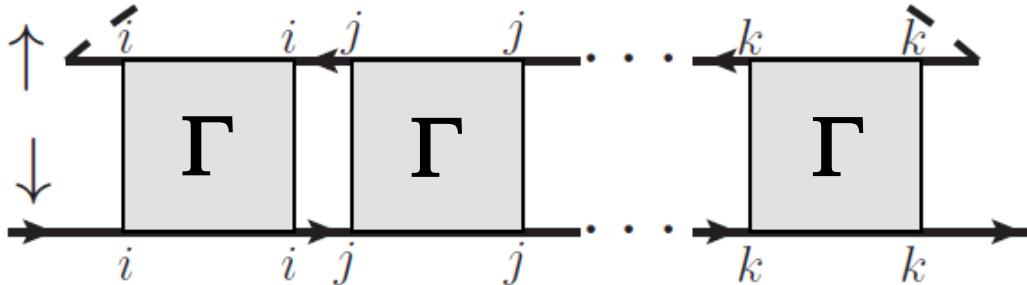
magnon self-energy Edwards-Hertz
Hertz-Millis-Moriya theory for QCP



$$\chi_{q\omega} = \chi^0_{q\omega} / (1 - \mathbf{U} \chi^0_{q\omega})$$

(b)

magnons but now including
all local DMFT physics



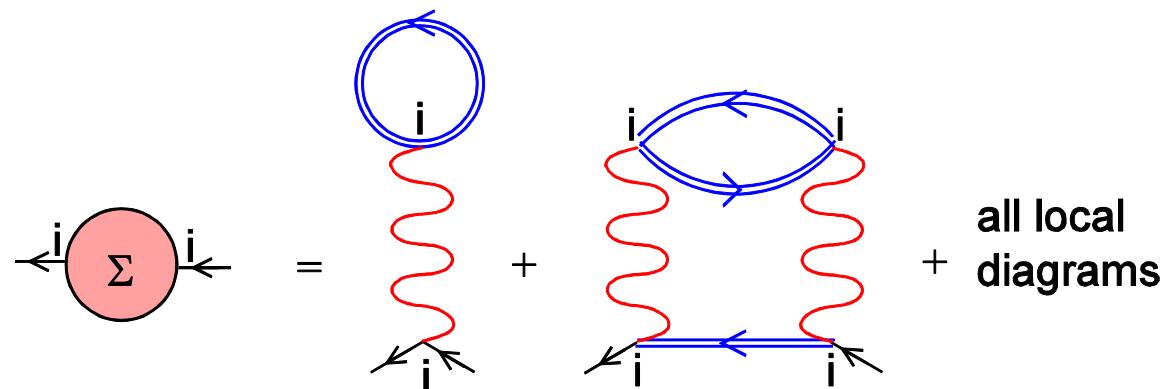
$$\chi_{q\omega} = \chi^0_{q\omega} (1 - \Gamma_q \chi^0_{q\omega})^{-1}$$

matrices in v, v'

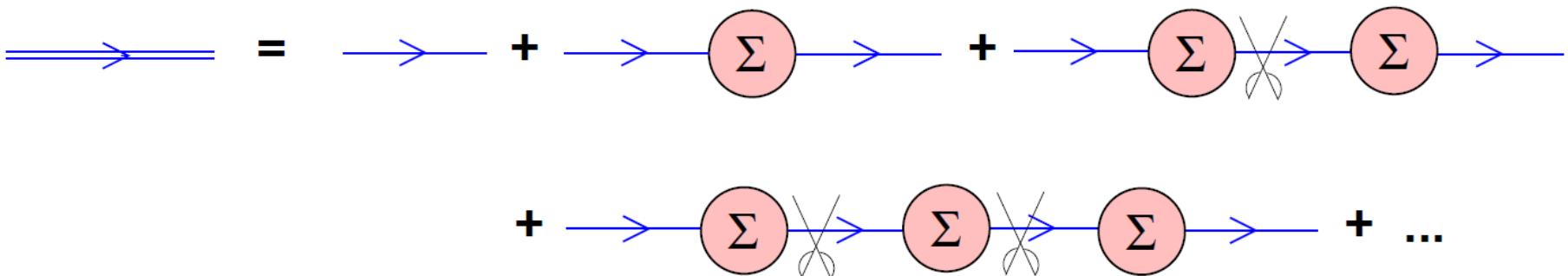
Resummation of Feynman diagrams in terms of locality
 n-particle fully irreducible vertex *approximated as local*

n=1: DMFT

- local 1-particle fully irreducible vertex Σ



Σ : one-particle irreducible one-particle vertex



Resummation of Feynman diagrams in terms of **locality**
 n-particle fully irreducible vertex ***approximated as local***

n=1: DMFT

- local 1-particle fully irreducible vertex Σ
- local 2-particle fully irreducible vertex Λ

n=2: DΓA

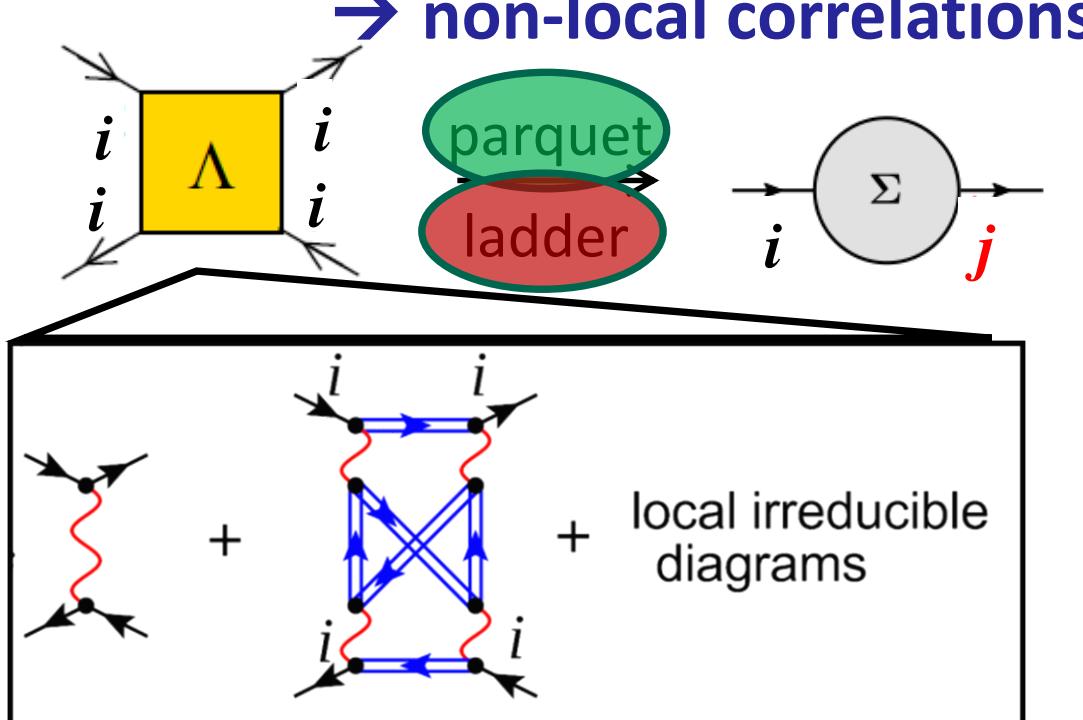
→ non-local correlations

n=3: error estimate

Ribic et al. PRB'17, 18

...

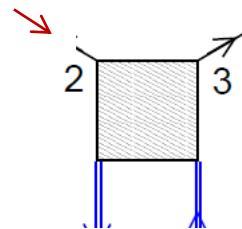
n → ∞: exact



Two-particle irreducibility

$$\text{full 2-particle vertex } \mathbf{F} = -\chi - \text{full 2-particle vertex } \mathbf{F}$$

3 ways (channels) to cut two lines



$$\text{full vertex } \mathbf{F} = \Lambda + \Phi_{ph} + \Phi_{\bar{ph}} + \Phi_{pp}$$

parquet equation

Parquet equations

$$\begin{array}{c} \text{Diagram: } \Phi_{ph}^{(2)} = F^{(2)} \Lambda^{(3)} + F^{(2)} \Phi_{ph}^{(3)} + F^{(2)} \Phi_{pp}^{(3)} \\ \text{Legend: } \text{Blue arrow} \rightarrow, \text{Blue arrow} \leftarrow, \text{Yellow box} = \end{array}$$

$$\begin{array}{c} \text{Diagram: } \Phi_{ph}^{(2)} = F^{(2)} + \Phi_{ph}^{(2)} + \Phi_{pp}^{(2)} \\ \text{Legend: } \text{Blue arrow} \uparrow, \text{Blue arrow} \downarrow, \text{Yellow box} = \end{array}$$

$$\begin{array}{c} \text{Diagram: } \Phi_{pp}^{(2)} = F^{(2)} \Lambda^{(3)} + F^{(2)} \Phi_{ph}^{(3)} + F^{(2)} \Phi_{ph}^{(3)} \\ \text{Legend: } \text{Blue arrow} \swarrow, \text{Blue arrow} \nwarrow, \text{Yellow box} = \end{array}$$

$$\begin{array}{c} \text{Diagram: } F^{(2)} = \Lambda^{(2)} + \Phi_{ph}^{(2)} + \Phi_{ph}^{(2)} + \Phi_{pp}^{(2)} \\ \text{Legend: } \text{Blue arrow} \leftarrow, \text{Blue arrow} \leftarrow, \text{Yellow box} = \end{array}$$

3 Bethe-Salpeter eq.



Λ given

$\rightarrow F, \Phi_{ph}, \Phi_{ph}, \Phi_{pp}, \Sigma, G$

$\Lambda = U$ (parquet approx.)

$\Lambda = \Lambda_{loc}$ ($D\Gamma A$)

Schwinger-Dyson

$$\begin{array}{c} \text{Diagram: } \Sigma = - F + \text{loop} \\ \text{Legend: } \text{Red wavy line}, \text{Blue arrow} \leftarrow, \text{Blue arrow} \rightarrow, \text{Yellow box} = \end{array}$$

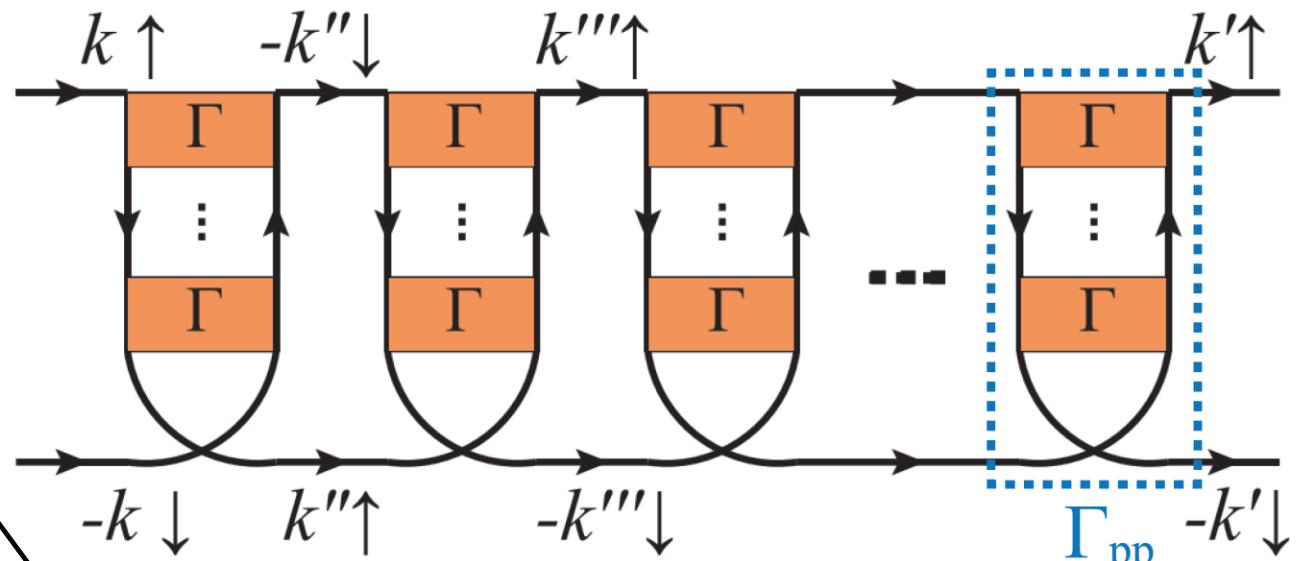
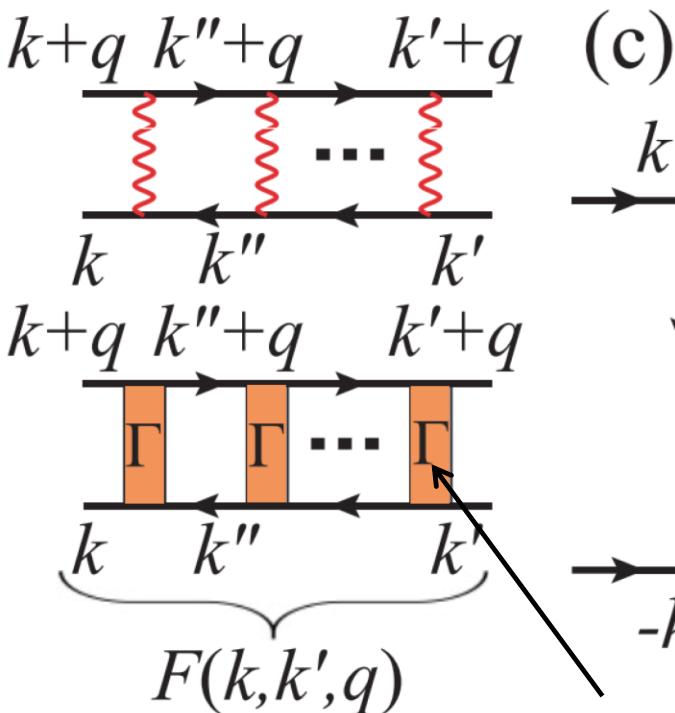
Dyson equation

$$\begin{array}{c} \text{Diagram: } \text{Blue arrow} \rightarrow = \text{Blue arrow} \rightarrow + \Sigma \rightarrow + \Sigma \Sigma \rightarrow \dots \\ \text{Legend: } \text{Blue arrow} \rightarrow, \text{Yellow box} = \end{array}$$

Calculations using ladder D Γ A (with λ -correction)

details: Kitatani ... KH arXiv:2203.12844,

J. Phys. Mater. 5 034005 (2022)



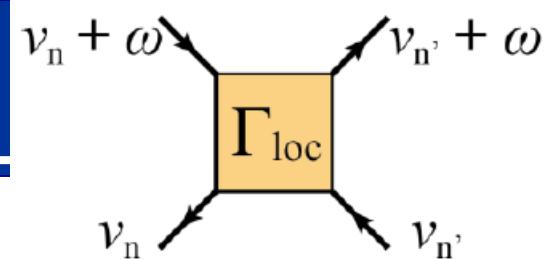
local from AIM, fully frequency dependent

Unbiased: spin and charge fluctuations, quantum critical fluct.

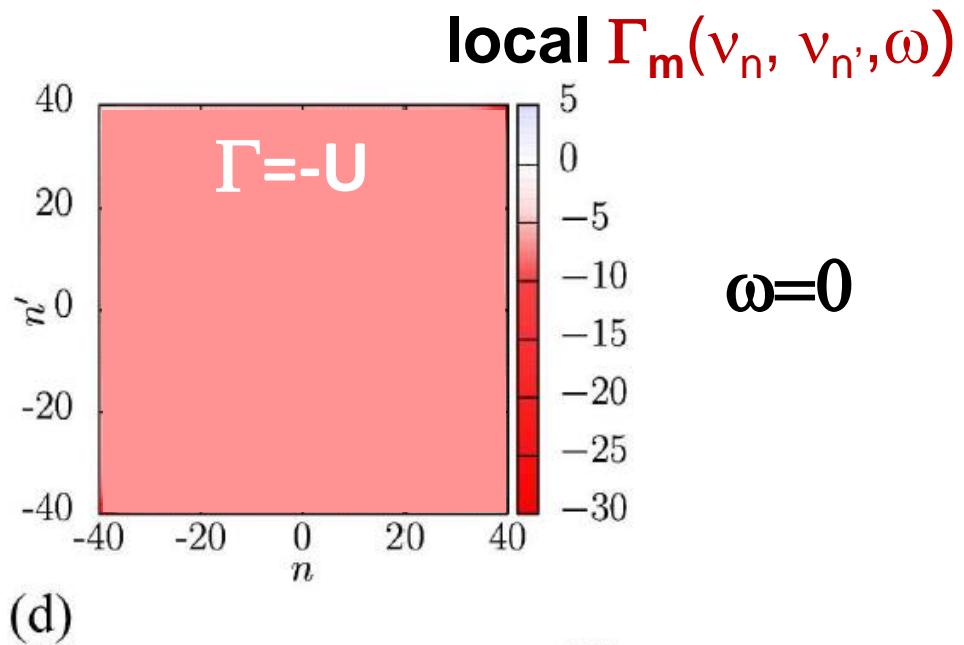
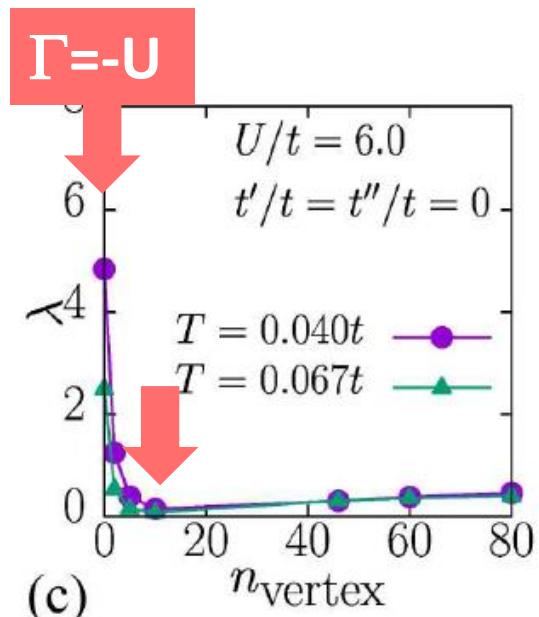
excellent agreement with diag. MC ...

Schäfer et al. PRX'21

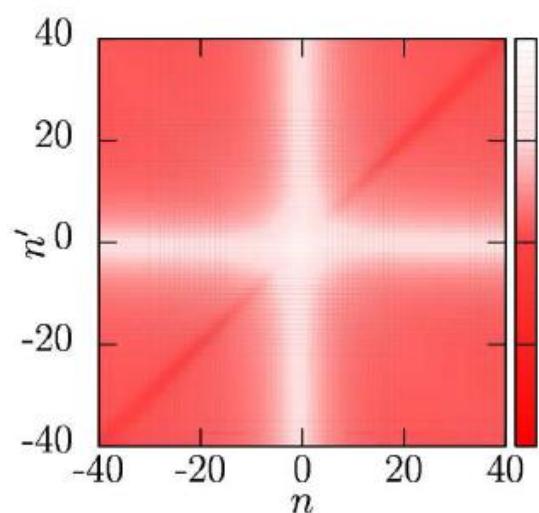
Importance of vertex dynamics



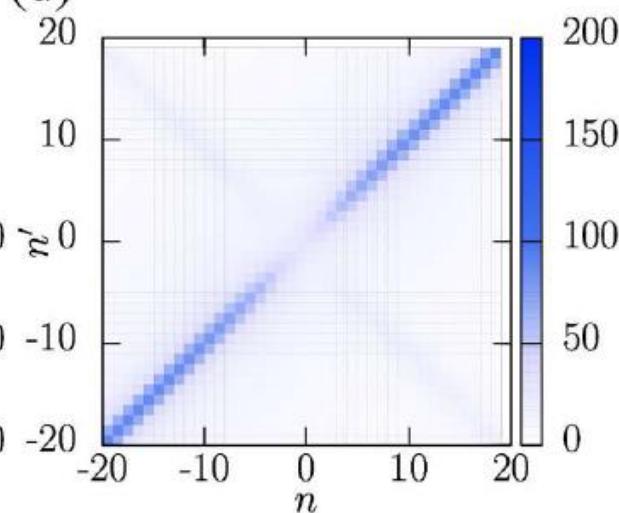
$U=6t$



nonlocal
F



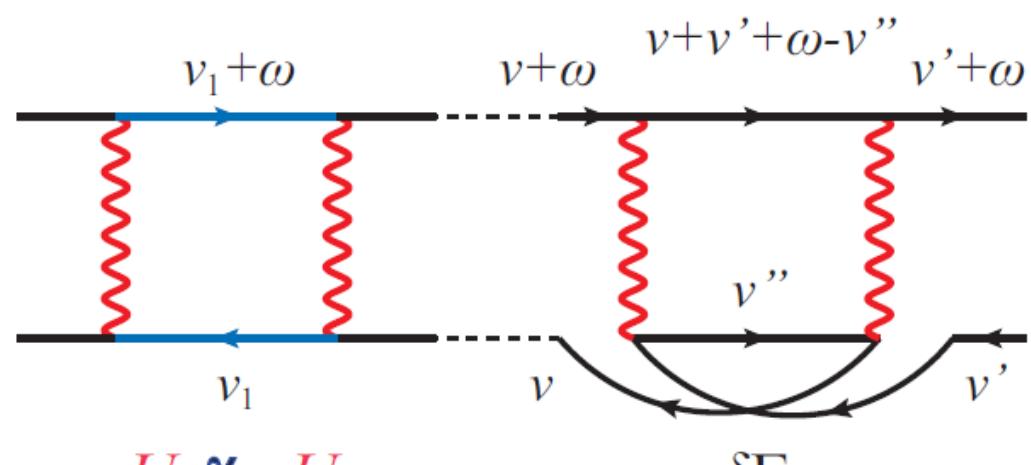
pairing
interaction
 Γ_{pp}



Importance of vertex dynamics

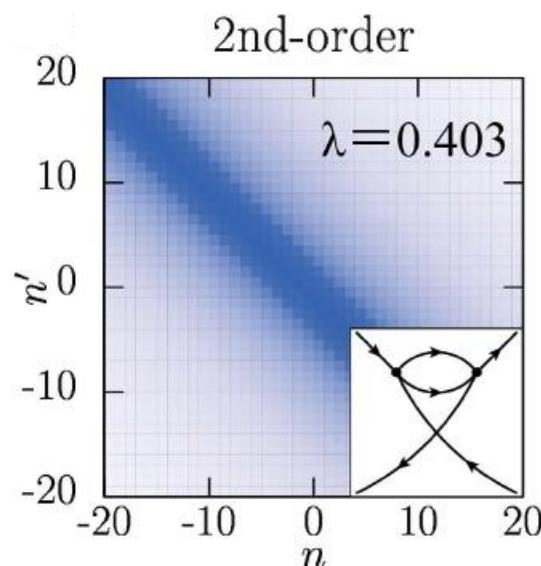
spin fluctuations

$$\chi = \chi_0 + \chi_0 U \chi_0 + \chi_0 U \chi_0 U \chi_0 + \dots$$
$$= \chi_0 / (1 - \chi_0 U)$$



suppressed by pp channel $\delta\Gamma_m$

2nd order



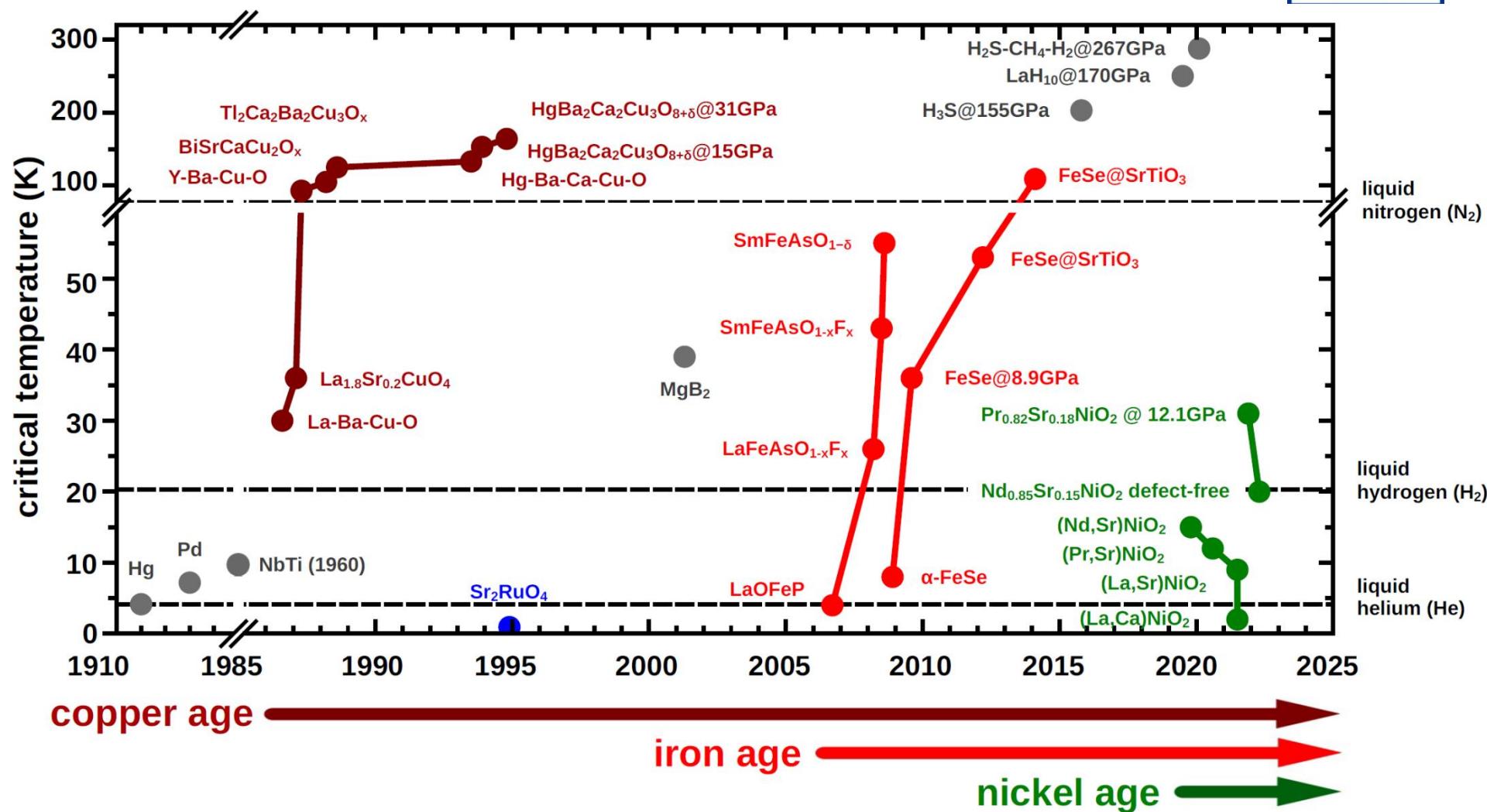
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New age of superconductivity: nickel age

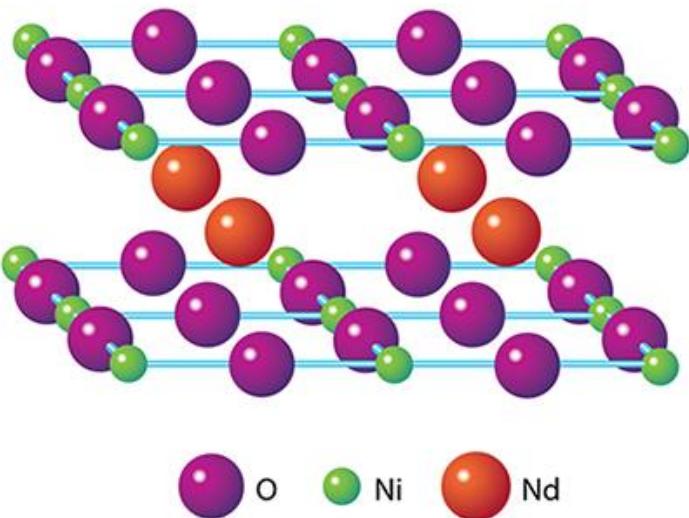


KH et al. Physik Journal Feb. 2022

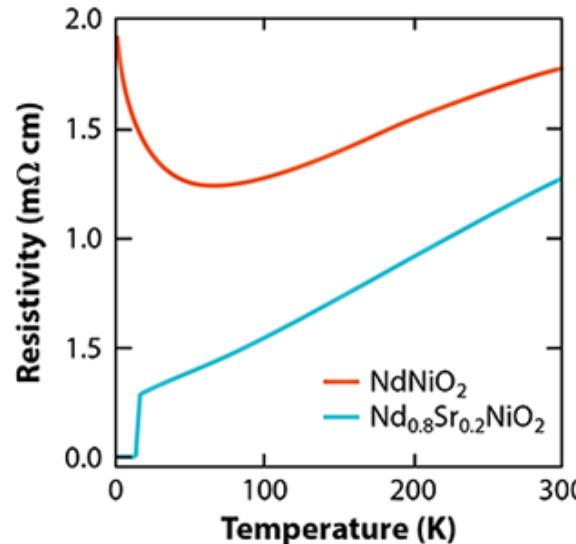
similar & different to cuprates → better understand high T_c supercond.

Nickelates & cuprates - similarities & differences

Anisomov et al. PRB'99 DFT+U: nickelates are cuprate analog



isoelectric: $3d^9$
isostructural
to $\text{Ca}^{2+}\text{Cu}^{2+}\text{O}_2^{4-}$



Li et al.
Nature'19



Zeng et al. PRL'20



Osada et al. Nano Lett'20



Zeng Sci. Adv.'22



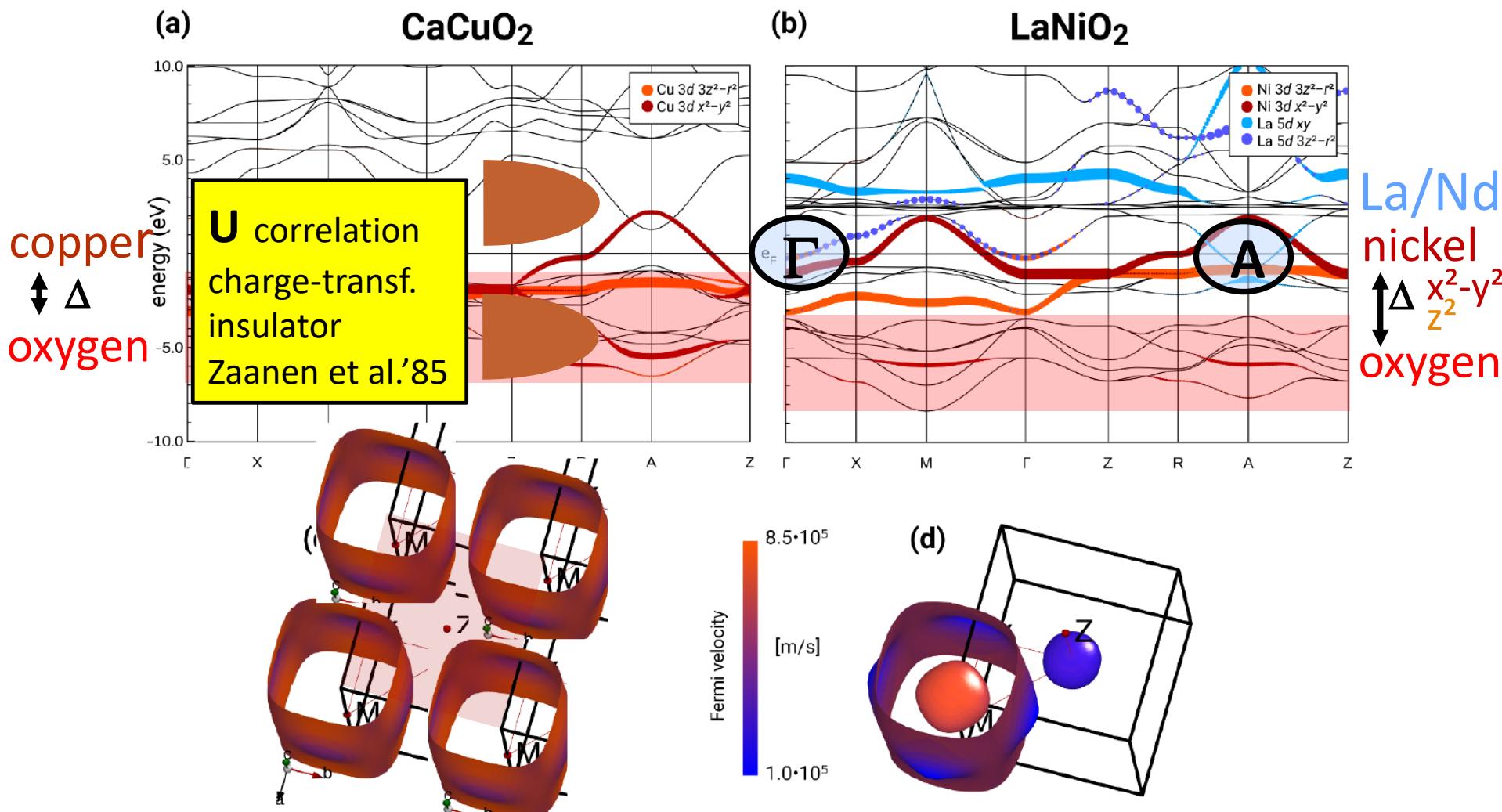
Osada et al. Adv. Mat.'21



Pan et al. Nature Phys.'21

high quality films, pressure ...

Nickelates & cuprates - similarities & differences

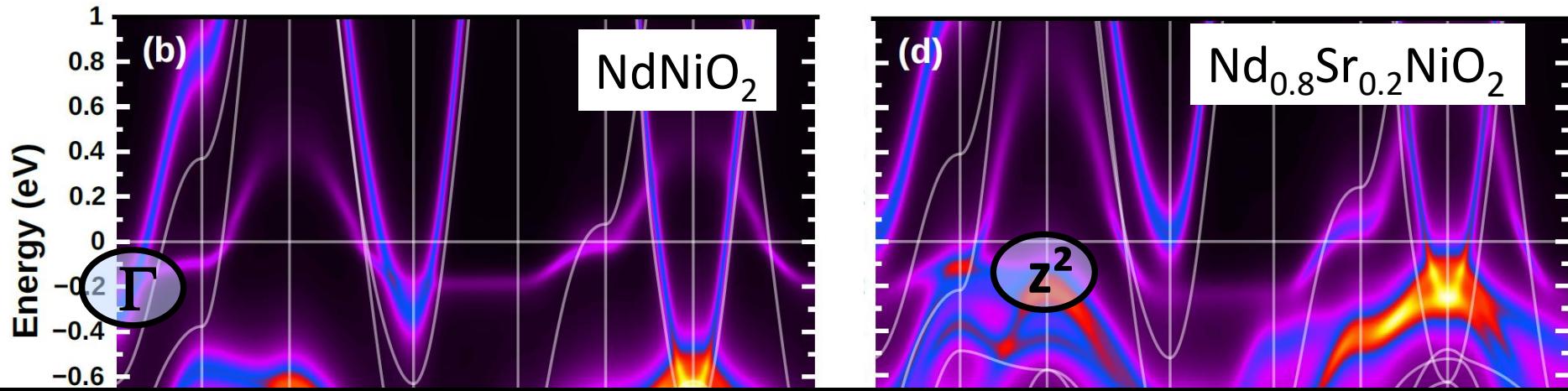


DFT bandstructure

cf. Botana Norman PRX'20 Sakakibara et al. PRB'20

Wu et al. PRB'20, Zhang et al PRB'20, Werner, Hoshino PRB'20 ...

DFT+DMFT: Ni x^2-y^2 + A pocket



Occam's razor: Ni x^2-y^2 + A pocket



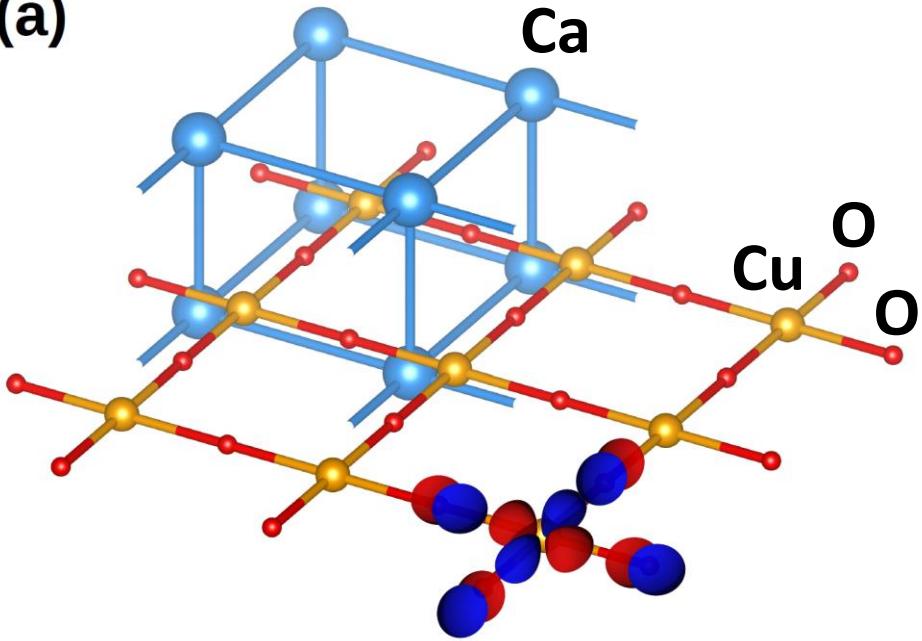
inverse Occam's razor

Mazin Nature Phys.'22

Cuprates & nickelates - similarities & differences

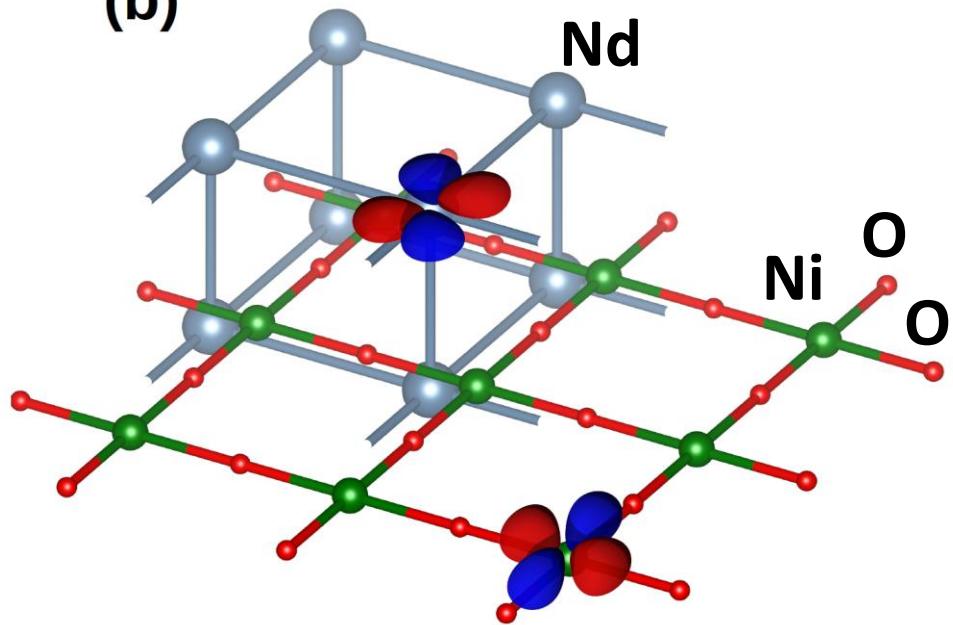
cuprates

(a)



nickelates

(b)

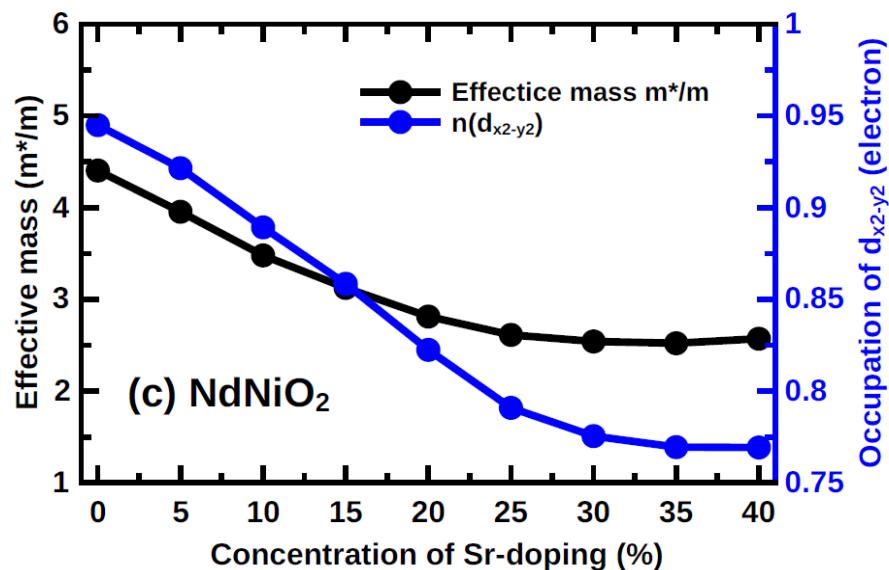
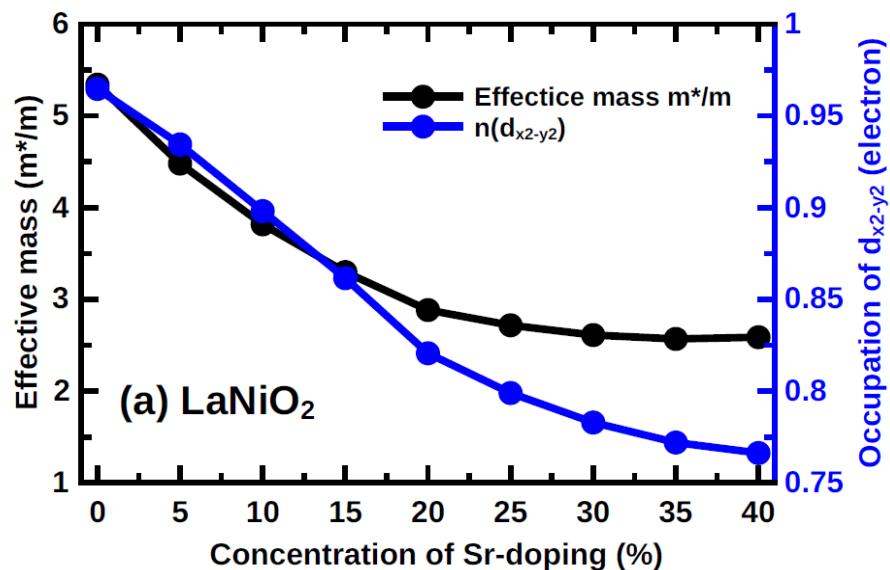


oxygen orbitals important
holes → oxygen orbitals
Emery model

(Hubbard model at most
mimics Zhang-Rice singlet)

A pocket only electron reservoir
holes $\frac{1}{2}$ A pocket $\frac{1}{2}$ $Ni x^2-y^2$

Hubbard model
w properly calculated doping



NdNiO_2 :

1-band Hubbard model with

$$U=8t \quad (\text{cRPA: } U = 2.6\text{eV} = 6.7 t)$$

$$t = 395 \text{ meV}, \quad t'/t = -0.25, \quad t''/t = 0.12$$

and **doping** from above Fig.

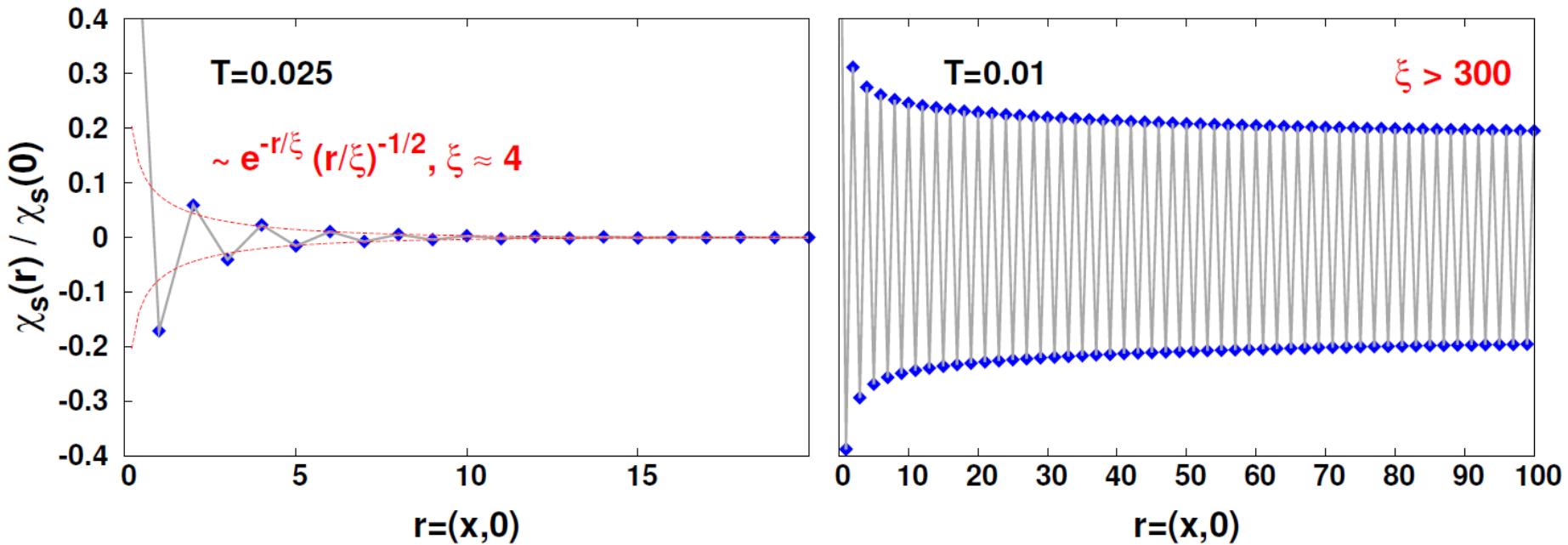
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Spin fluctuations

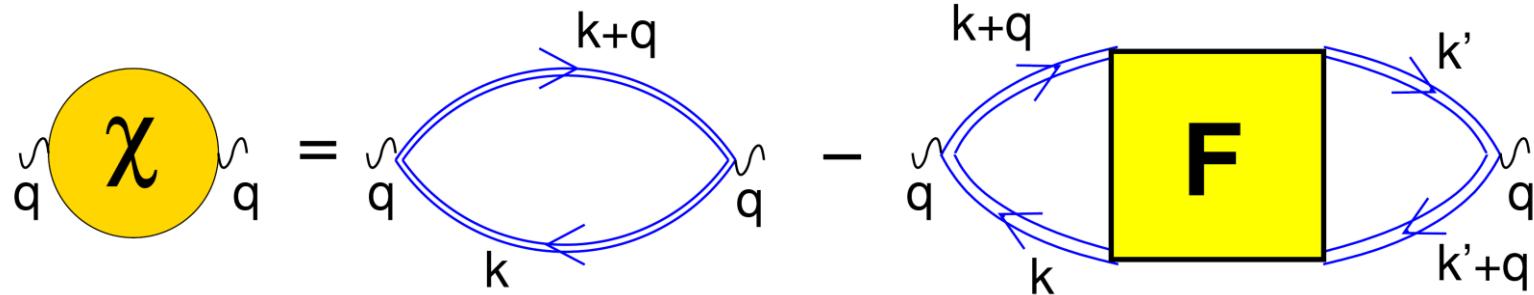


$$\chi(r) = \langle S(r)S(0) \rangle \sim (r/\xi)^{-1/2} e^{-r/\xi}$$

in practice: Ornstein-Zernike form

$$\chi_{\omega=0} q \sim 1/[(q-Q)^2 + \xi^{-2}]$$

Spin fluctuations

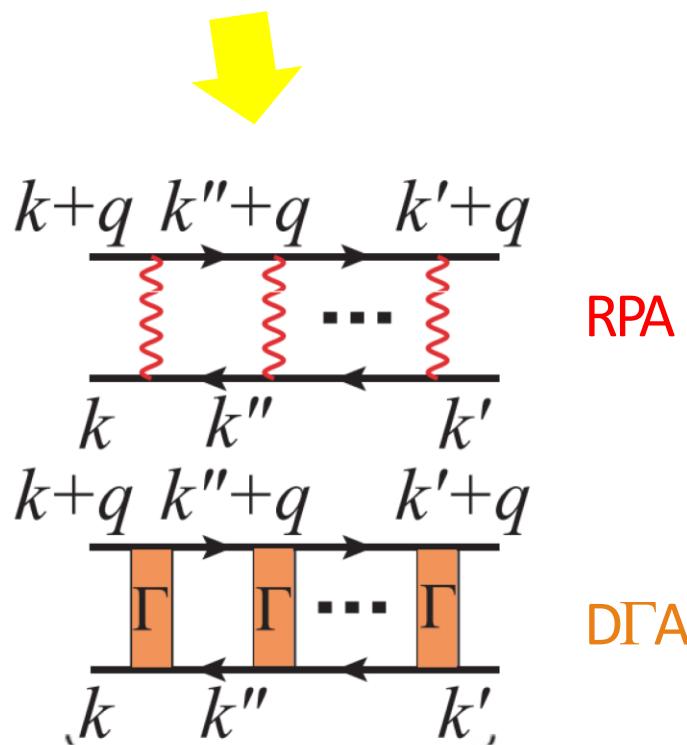


$$\chi_{r,\mathbf{q}\omega} = \underbrace{\sum_{\mathbf{k}\nu} G_{\mathbf{k}\nu} G_{(\mathbf{k}+\mathbf{q})(\nu+\omega)}}_{\equiv \chi_{\mathbf{q}\omega}^0}$$

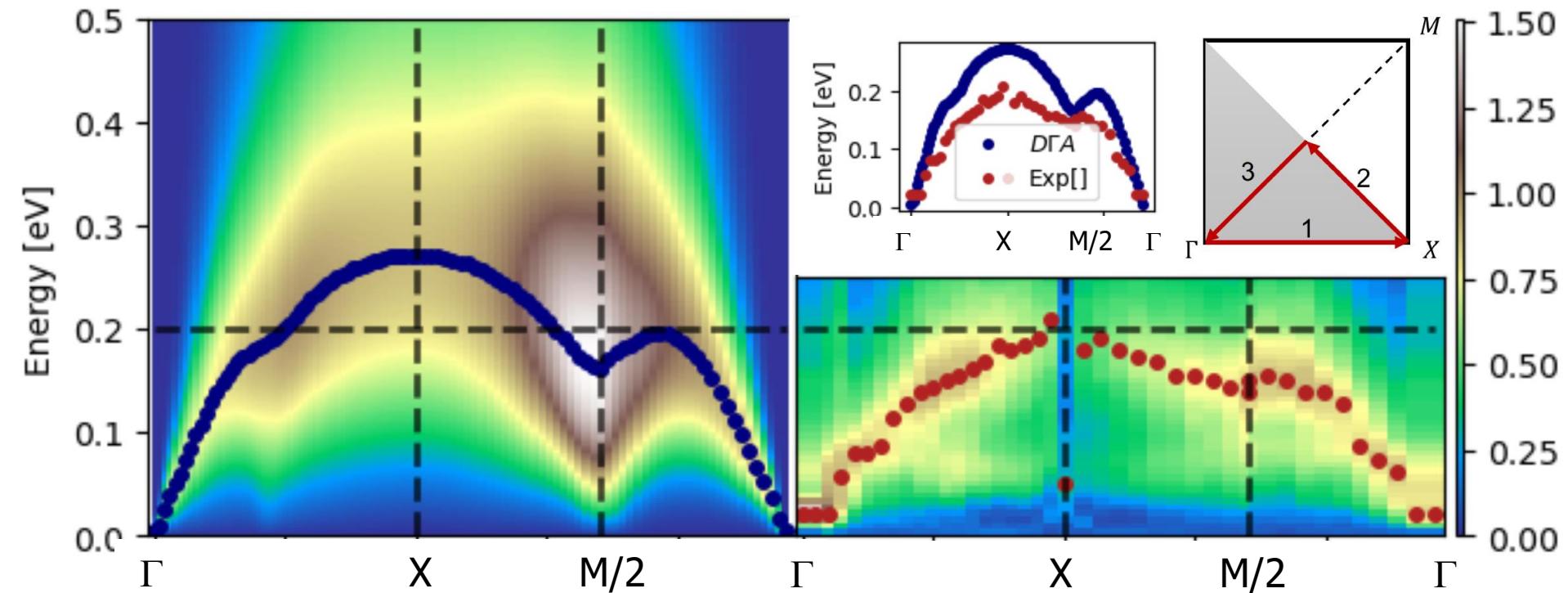
$$- \sum_{\mathbf{k}\mathbf{k}'\nu\nu'} G_{\mathbf{k}\nu} G_{(\mathbf{k}+\mathbf{q})(\nu+\omega)} F_{r,\mathbf{k}+\mathbf{k}'\mathbf{q}}^{\nu\nu'\omega} G_{\mathbf{k}'\nu} G_{(\mathbf{k}'+\mathbf{q})(\nu+\omega)}$$

$$\chi_{\mathbf{q}\omega} = \chi_{\mathbf{q}\omega}^0 / (1 - \mathbf{U} \chi_{\mathbf{q}\omega}^0) \quad \text{RPA}$$

$$\chi_{\mathbf{q}\omega} = \chi_{\mathbf{q}\omega}^0 (1 - \mathbf{\Gamma}_{\mathbf{q}} \chi_{\mathbf{q}\omega}^0)^{-1} \quad \text{DΓA matrices in } \nu, \nu'$$



Paramagnon spectrum NdNiO_2



D Γ A Theory

Worm ... KH unpublished'22

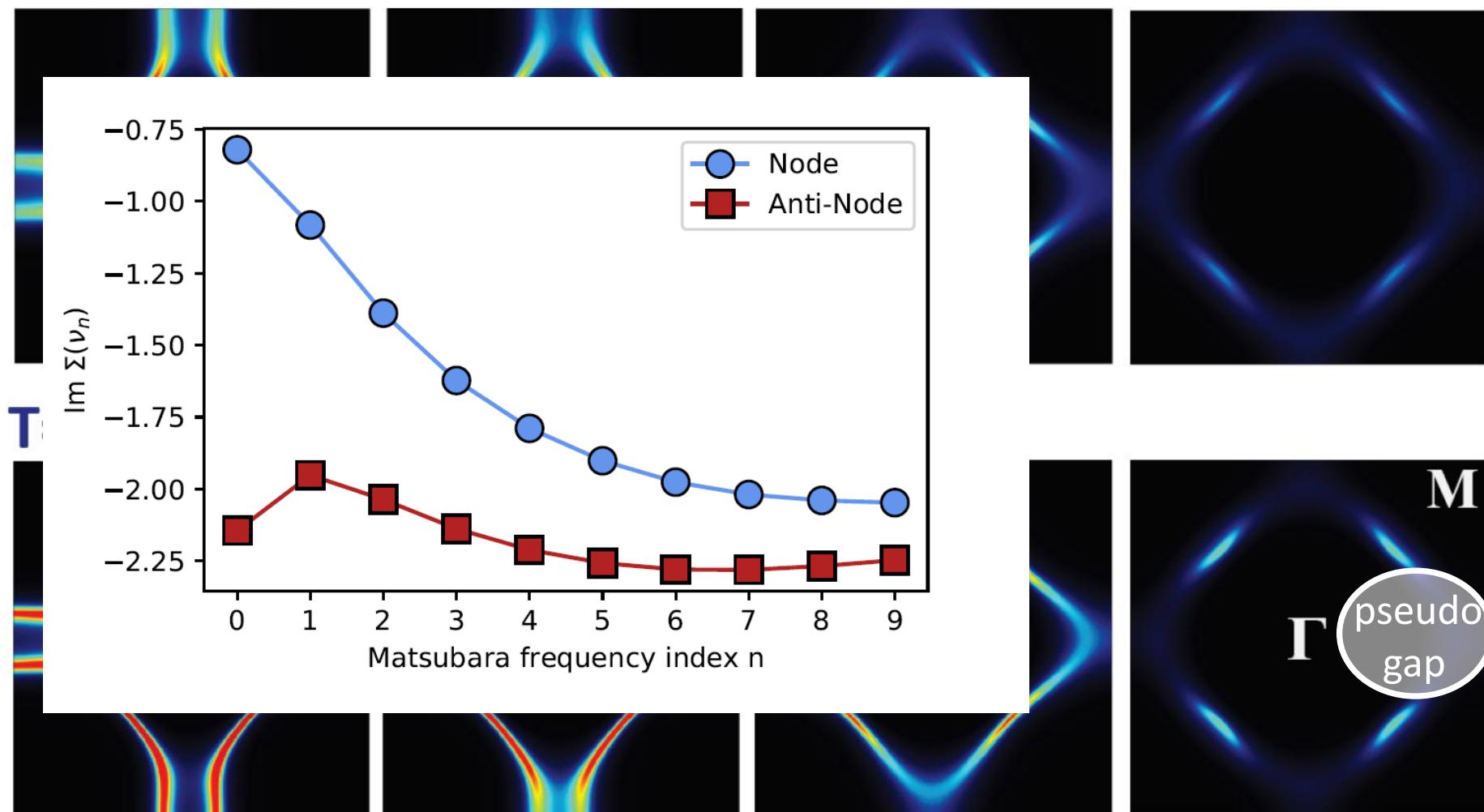
RIXS exp. Lu et al. Science'21

NdNiO_2 on STO (with stacking faults)
better films → better agreement?

Pseudogap: Nickelates

T=92K

D Γ A spectrum for nickelate Hubbard model



$n_{dx^2-y^2}=0.80$

$n_{dx^2-y^2}=0.85$

$n_{dx^2-y^2}=0.90$

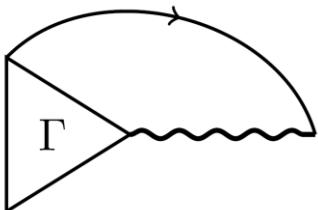
$n_{dx^2-y^2}=0.95$

Pseudogap: Im Γ mechanism

Krien, Worm ... KH arXiv:2101.06529

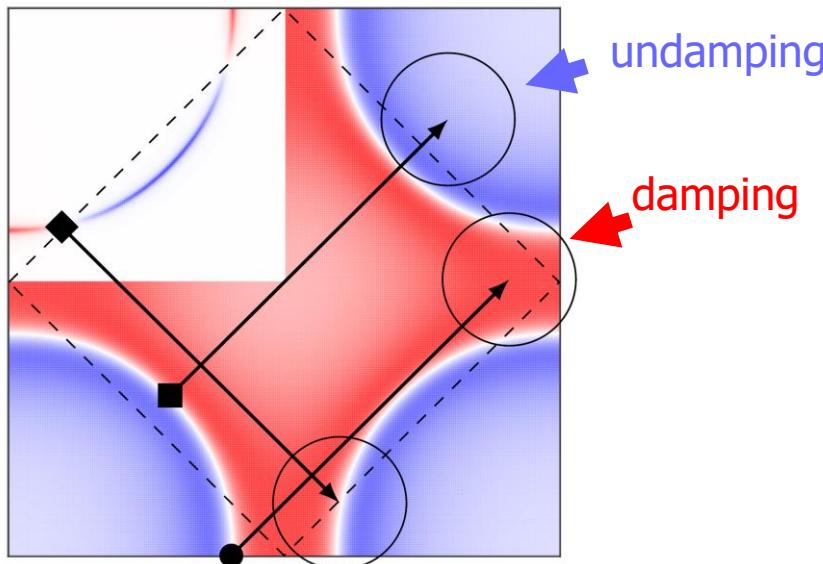
Comm. Phys in print

We can rewrite self-energy as **spin-fermion vertex**



$$\Sigma_{\text{sp}}(k, q) \propto -G_{k+q} W_q \Gamma_{kq}$$

at strong coupling: contribution from **Im Γ (!)**

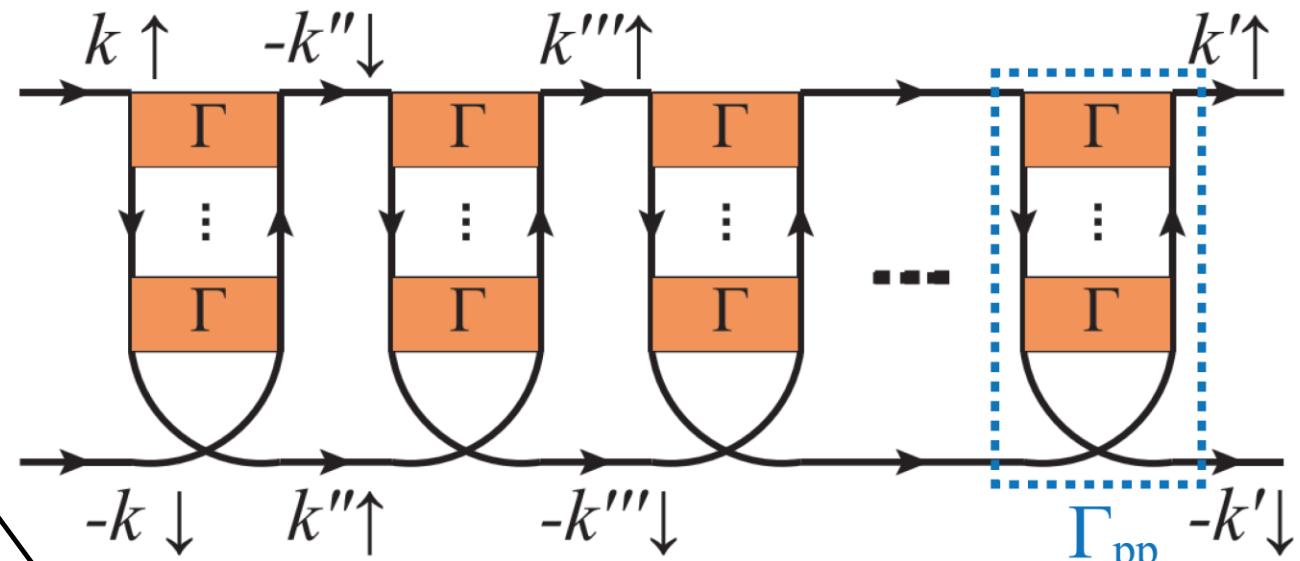
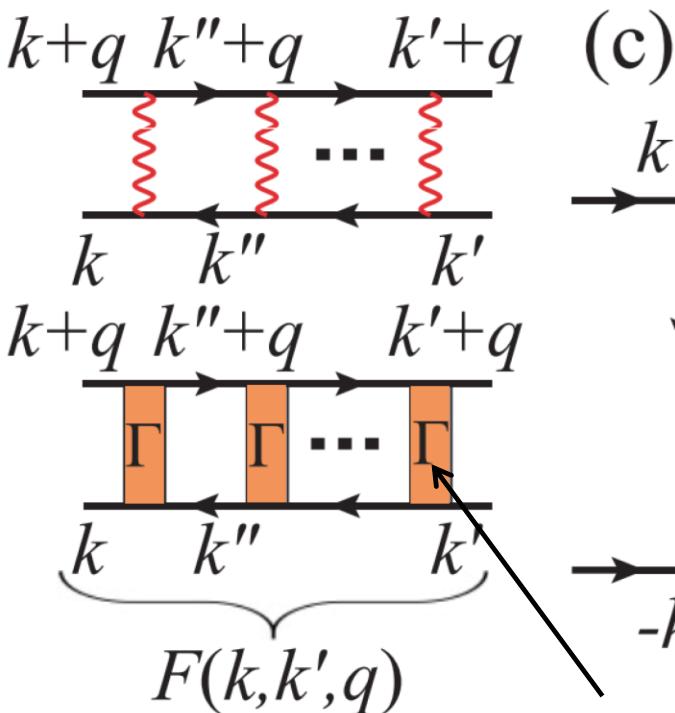


Superconductivity: nickelates

Calculations using ladder D Γ A (with λ -correction)

details: Kitatani ... KH arXiv:2203.12844,

J. Phys. Mater. 5 034005 (2022)



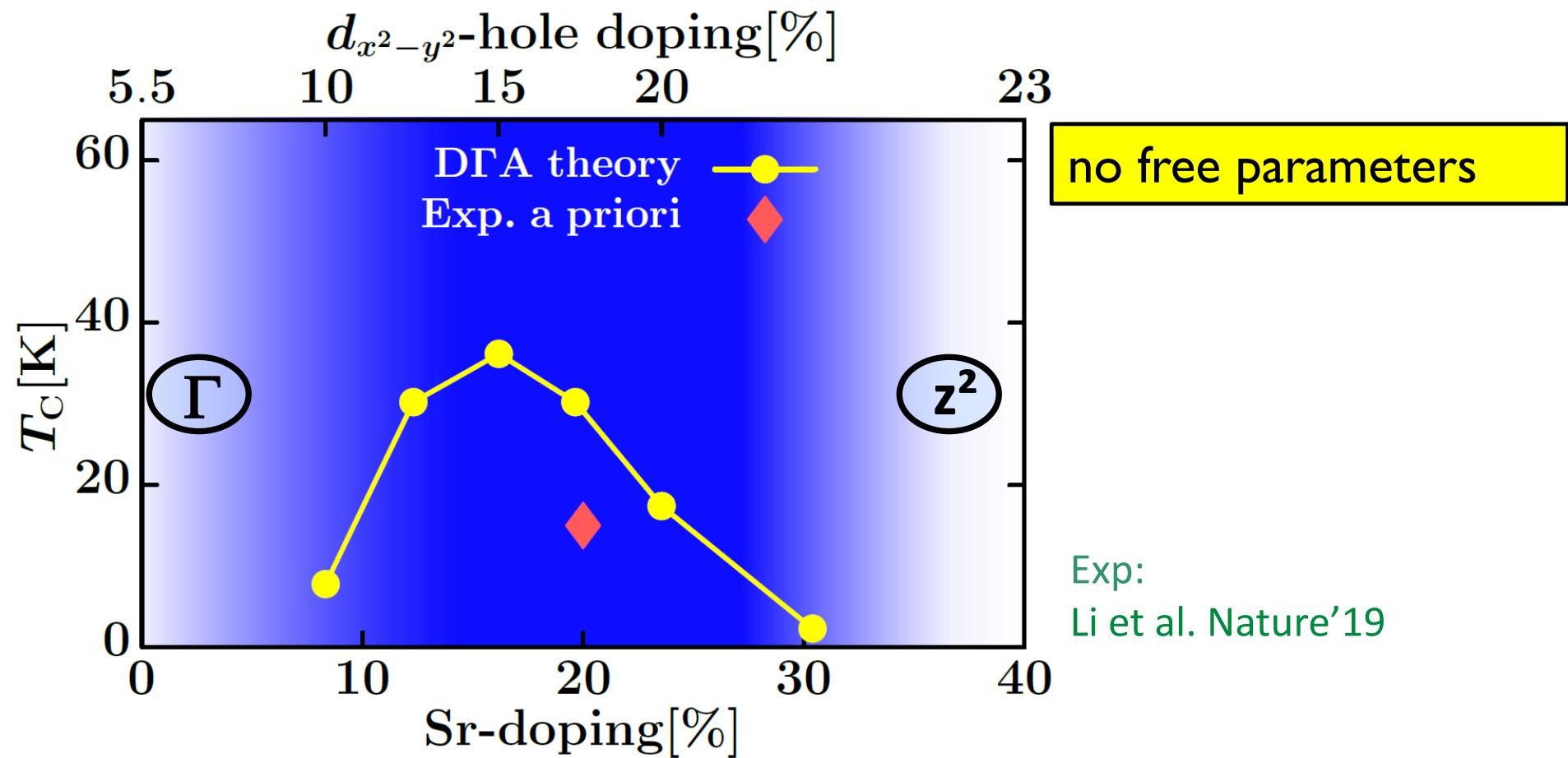
local from AIM, fully frequency dependent

Unbiased: spin and charge fluctuations, quantum critical fluct.

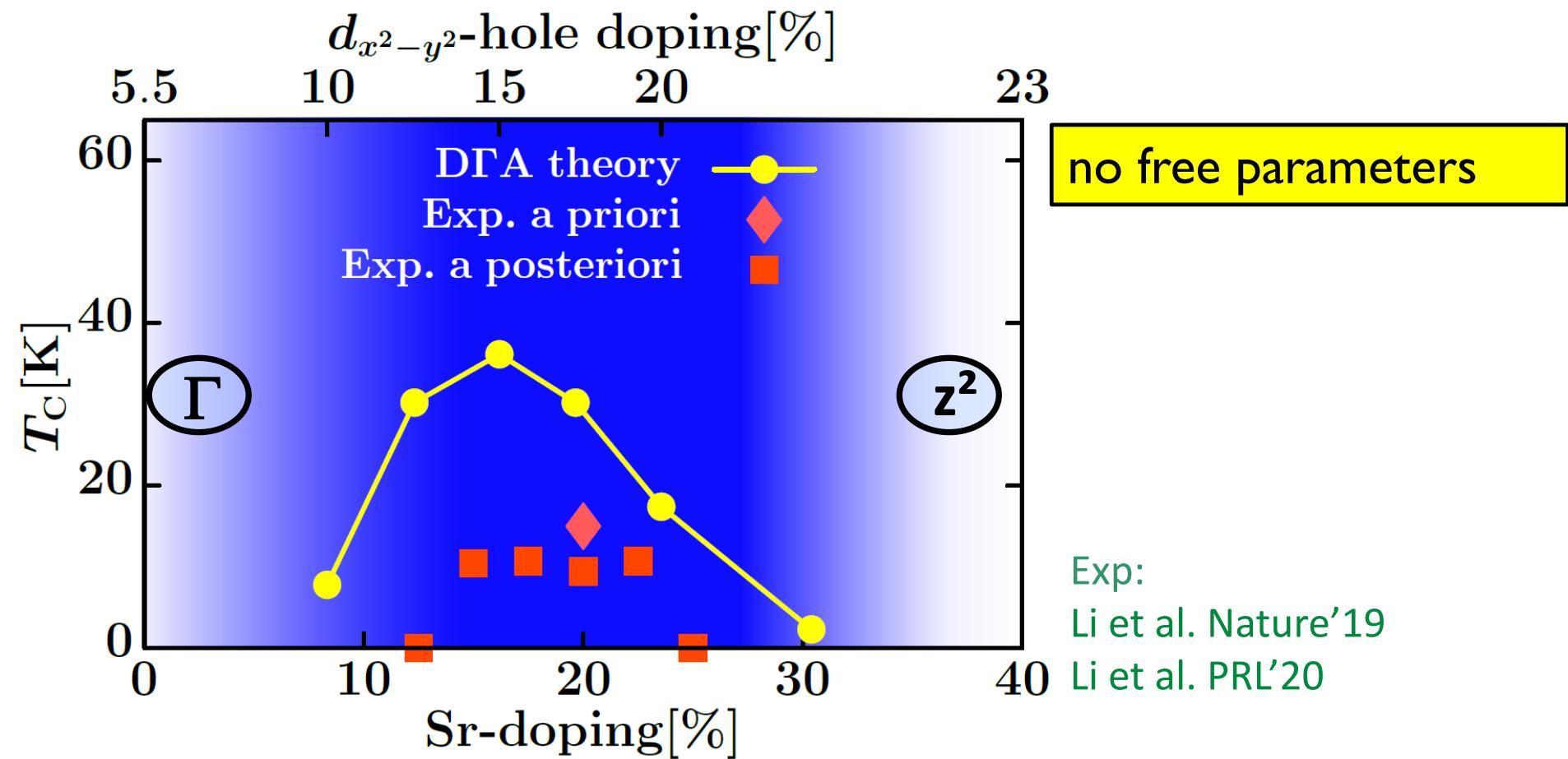
excellent agreement with diag. MC ...

Schäfer et al. PRX'21

Superconductivity $\text{Sr}_x\text{Nd}_{1-x}\text{NiO}_2$ (d-wave)

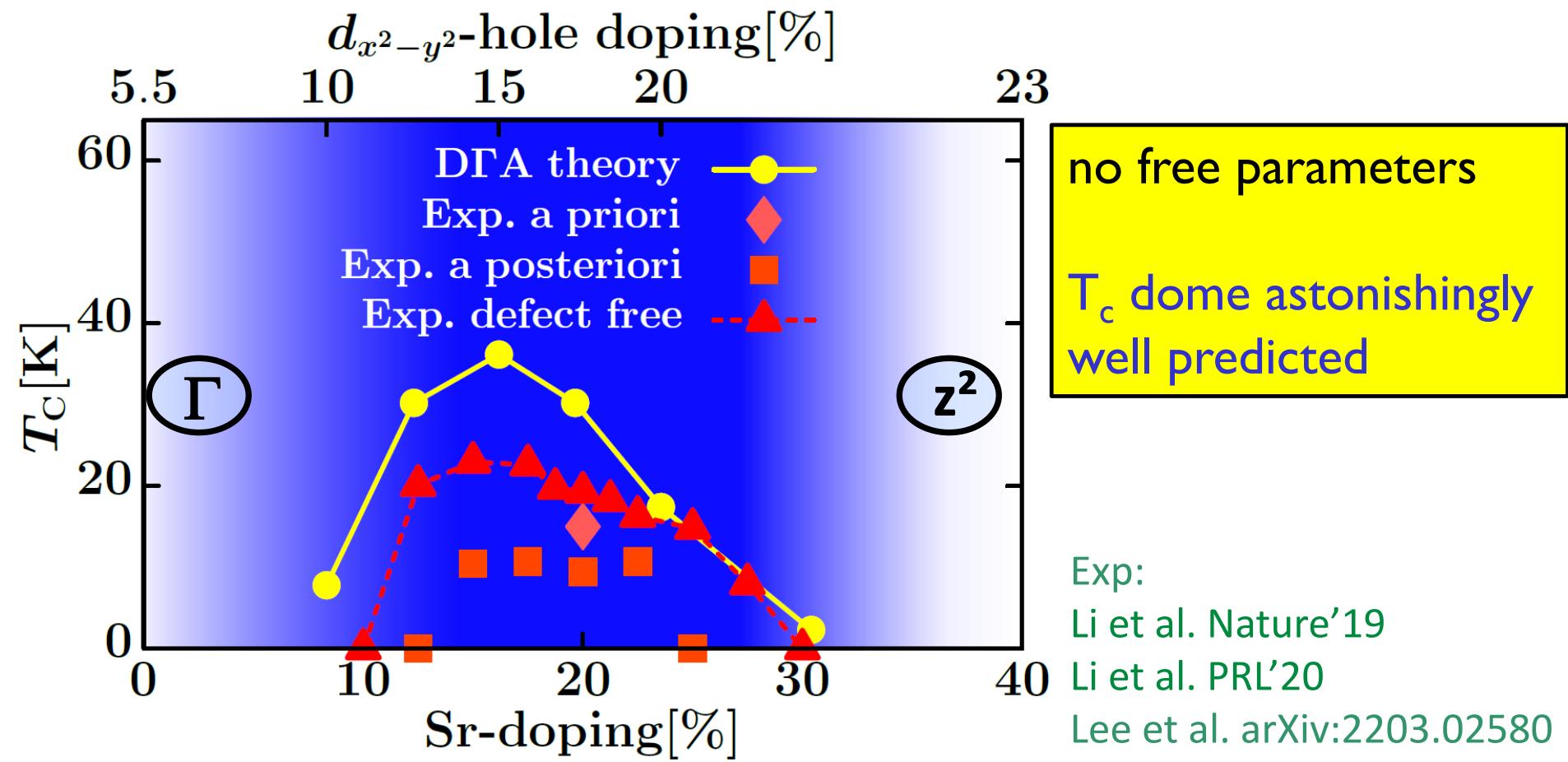


Superconductivity $\text{Sr}_x\text{Nd}_{1-x}\text{NiO}_2$ (d-wave)

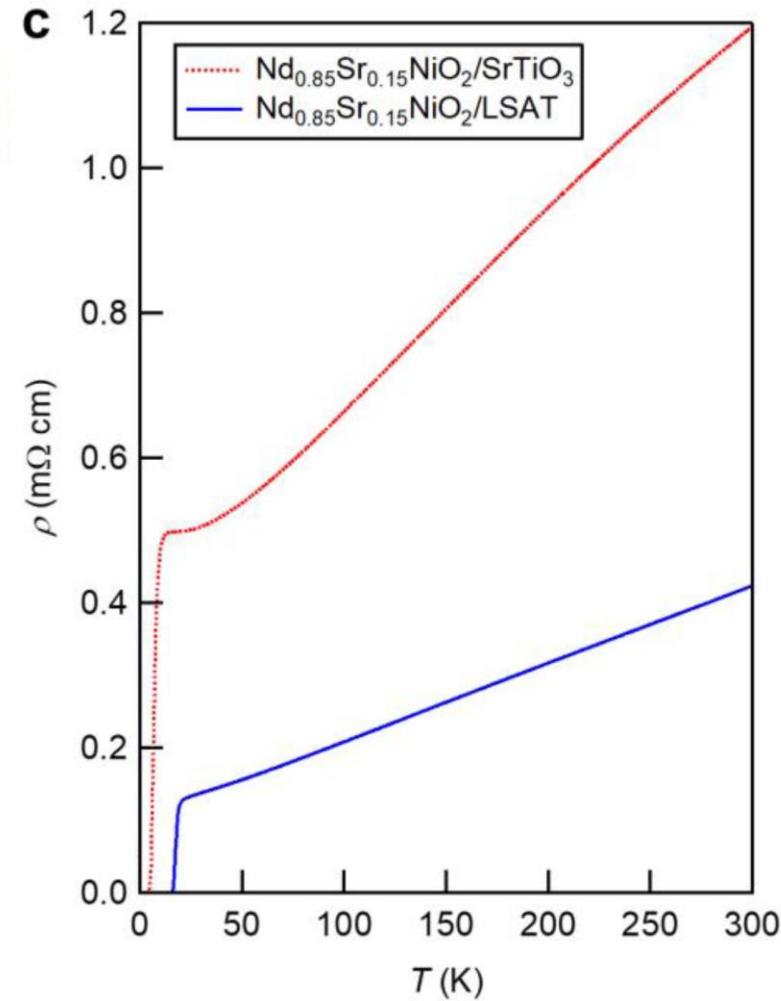
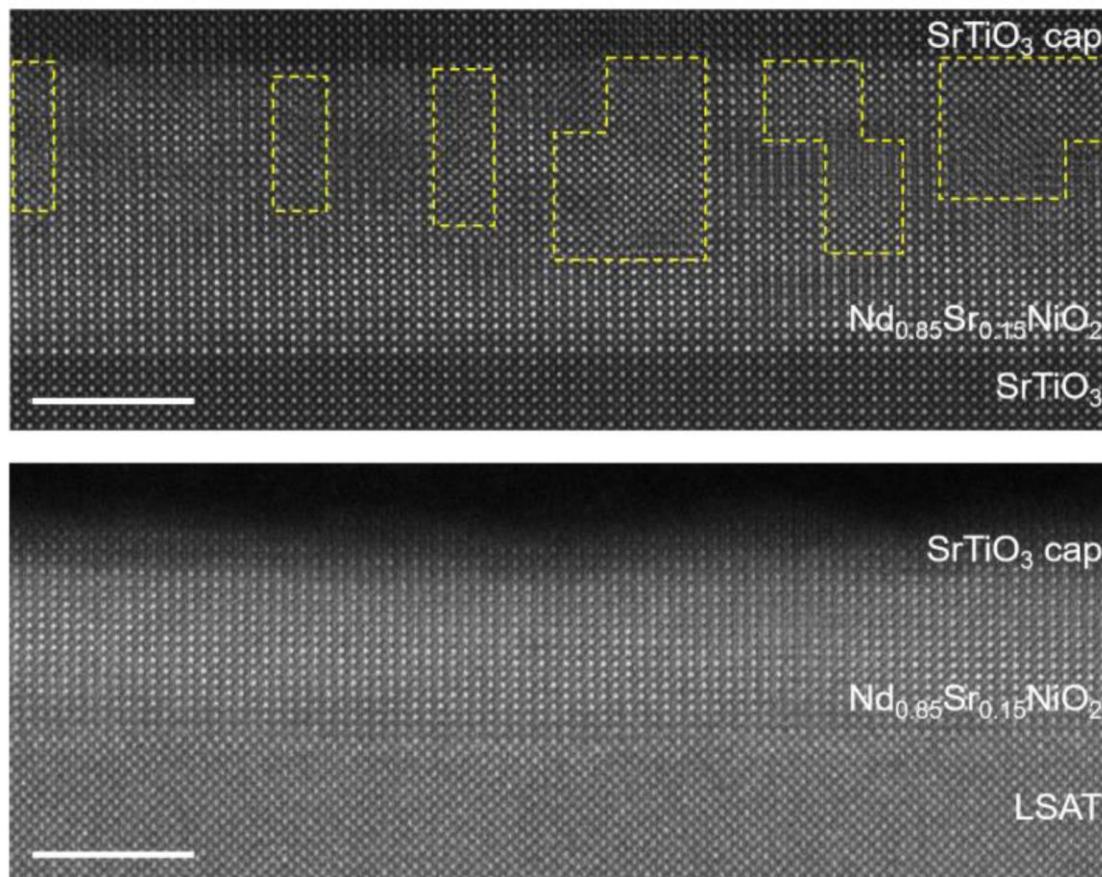


Superconductivity: nickelates

Superconductivity $\text{Sr}_x\text{Nd}_{1-x}\text{NiO}_2$ (d-wave)



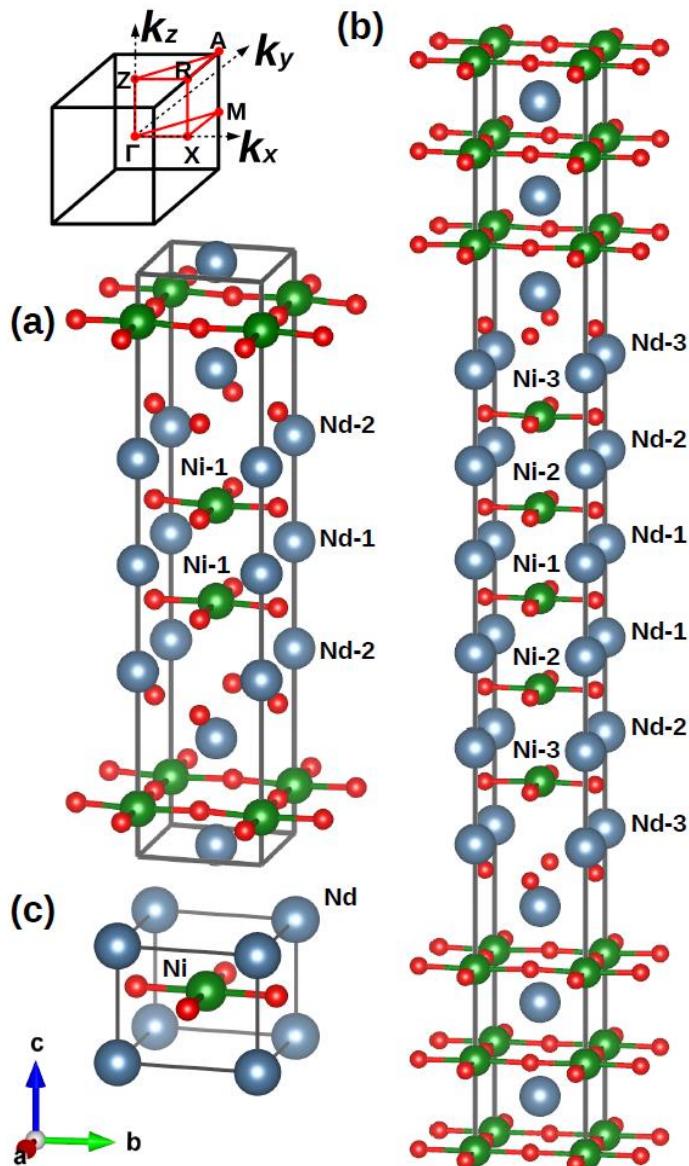
New, cleaner (NdSr) NiO_2 films



LSAT: 3.890 Å instead of STO: 3.905 Å

Lee et al. arXiv:2203.02580

From infinite layer to finite layer



Nd₆Ni₅O₁₂
Pan et al. Nature Phys.'21

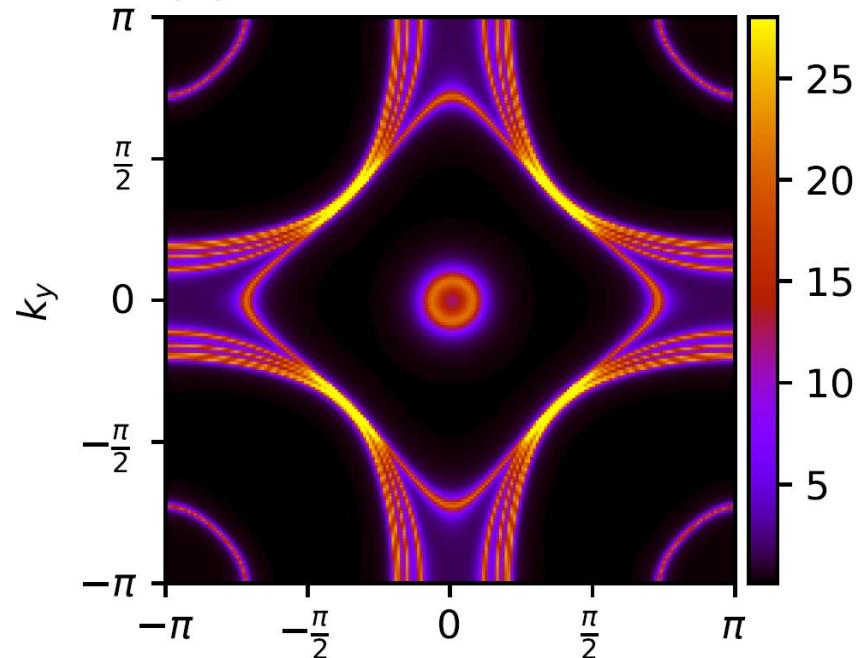
essentially same hopping
just cut-off after 5 layers

System	t	t'	t''
NdNiO ₂	-0.395	0.095	-0.047
Nd ₃ Ni ₂ O ₆	-0.414	0.092	-0.055
Nd ₆ Ni ₅ O ₁₂ : Ni-1	-0.395	0.098	-0.050
Nd ₆ Ni ₅ O ₁₂ : Ni-2	-0.392	0.097	-0.050
Nd ₆ Ni ₅ O ₁₂ : Ni-3	-0.398	0.097	-0.049

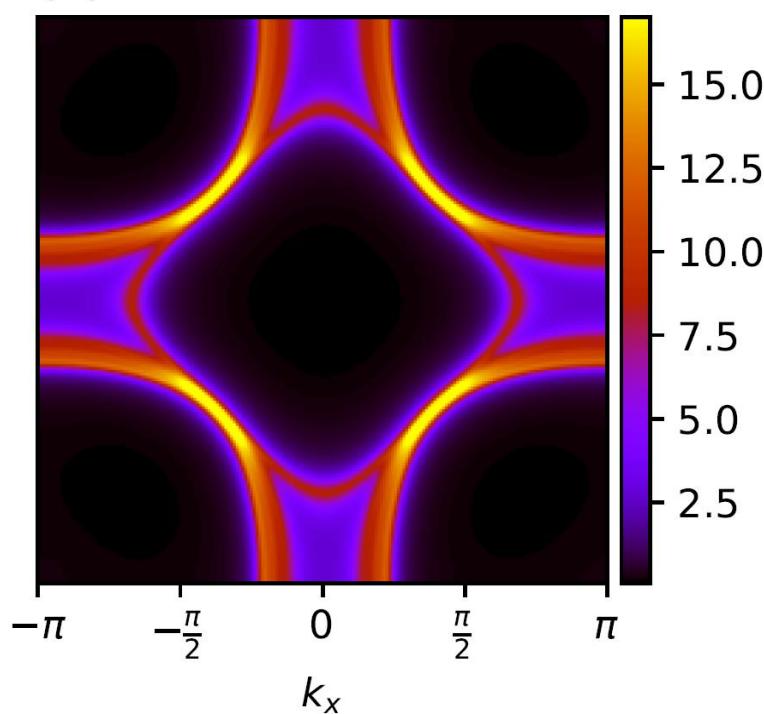
Pentalayer nickelate

$\text{Nd}_6\text{Ni}_5\text{O}_{12}$

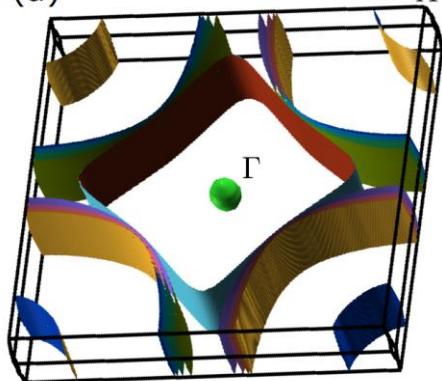
(a): DFT



(b): DMFT

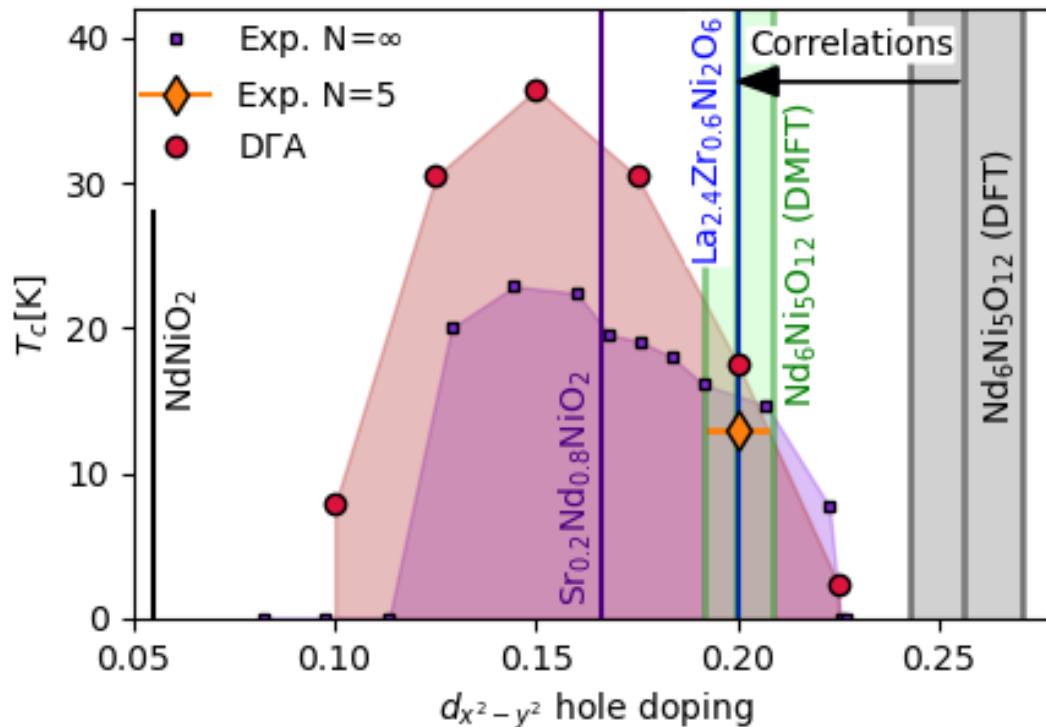


(a)



x^2-y^2 -orbital system Ni 3d^{8.8}
no pockets/tubes in DMFT

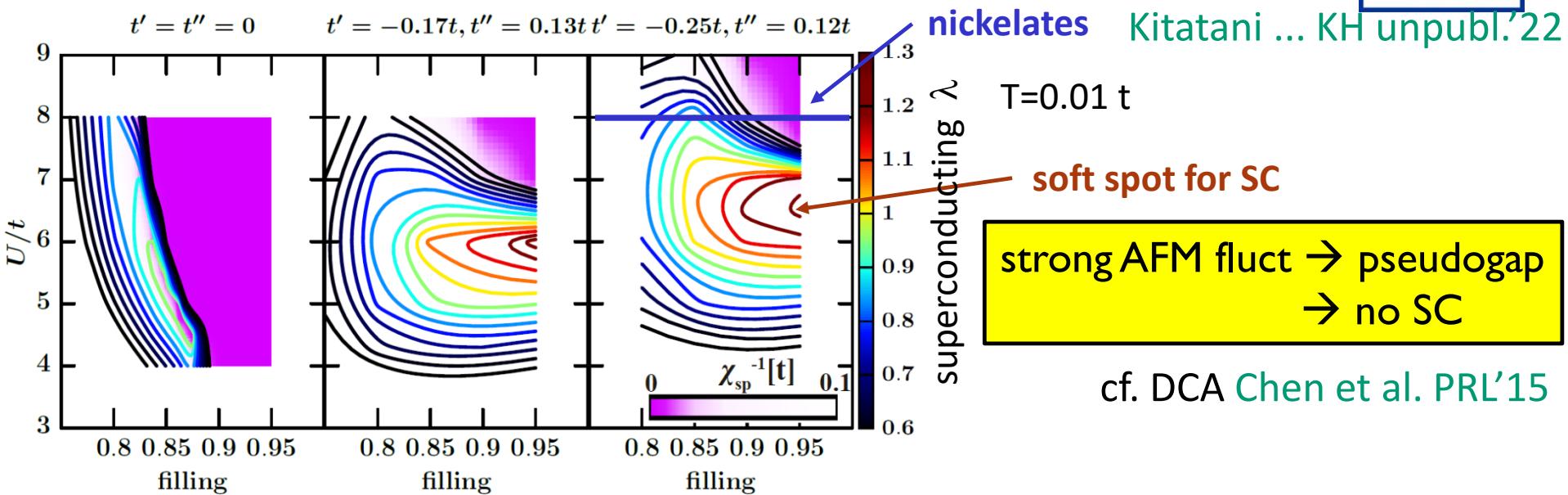
$\text{Nd}_6\text{Ni}_5\text{O}_{12}$



correlations → no pockets/tubes

T_c consistent with infinite layer nickelate at proper $d_{x^2-y^2}$ doping
exp. evidence: positive Hall coef. ✓

How to increase T_c ?



- (a) **pressure or compressive strain**
- (b) **substrate different from SrTiO_3**
- (c) **palladates with 4d instead of 3d**

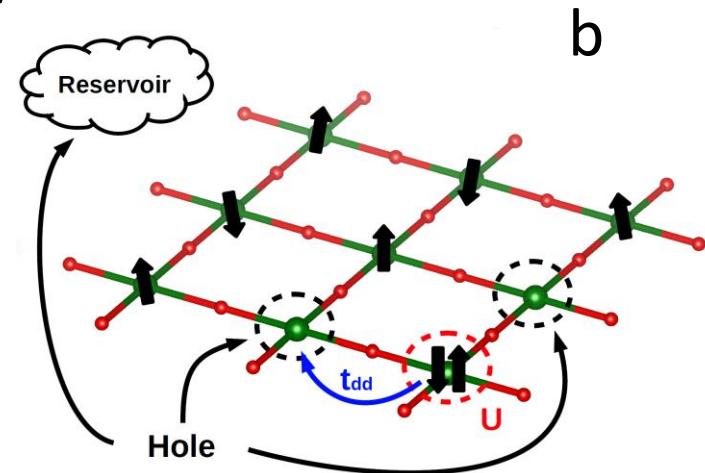
exp. confirmed: $\text{Pr}_{1-x}\text{Sr}_x\text{NiO}_2$ under 12GPa with $T_c > 30\text{K}$
Wang et al. arXiv:2109.12811

Conclusion

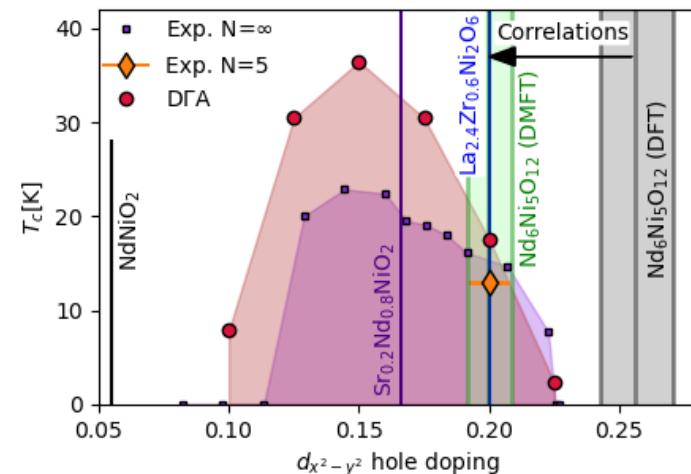
- Diagrammatic extensions of DMFT such as $D\Gamma A$ allow us to describe the physics beyond DMFT such as spin fluctuations superconductivity ...

- Nickelates

Hubbard model: Ni x^2-y^2 orbital
+ electron reservoir: A pocket



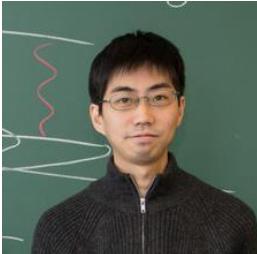
- good $D\Gamma A$ prediction of superconducting phase diagram consistent picture



Thanks to ...



Liang
Si



Motoharu
Kitatani
(Tokyo)



Paul
Worm



Josef
Kaufmann



Zhicheng
Zhong
(Ningbo)



Oleg
Janson
(Dresden)



Jan
Tomczak

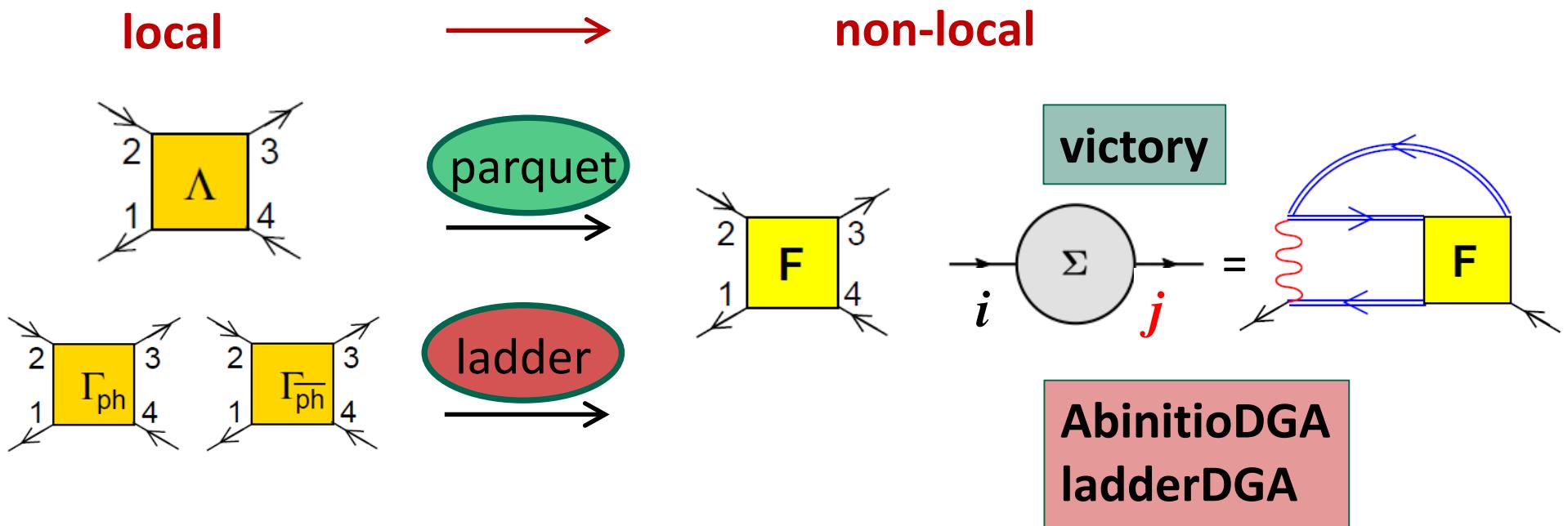


Si et al.
Kitatani et al.
Worm et al.
Held et al.

Phys. Rev. Lett. 124, 166402 (2020)
npj Qu. Materials 5, 59 (2020)
arXiv:2111.12687
Frontiers in Phys. 9, 810394 (2022)

DΓA Rohringer et al. Rev.Mod.Phys. 90, 025003 (2018)

D Γ A: Resummation of Feynman diagrams in terms of locality
 n-particle fully irreducible vertex *approximated as local*



includes local **DMFT** correlations
 and **non-local** correlations
 (e.g. spin fluctuations)

