

Quantum criticality and superconductivity in diagrammatic extensions of DMFT

Karsten Held (TU Wien)

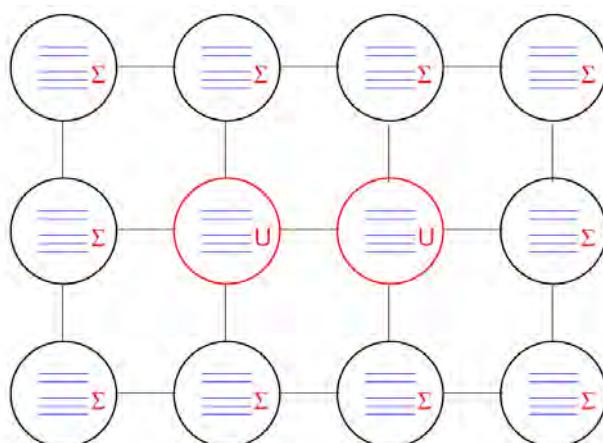
Jülich, Sept. 21st 2018

- Motivation
- Synopsis: Dynamical vertex approximation (DΓA)
- Quantum criticality in 3D Hubbard model
- Quantum criticality in 2D periodic Anderson model
- Superconductivity in 2D Hubbard model

Please ask questions

Correlations beyond DMFT:

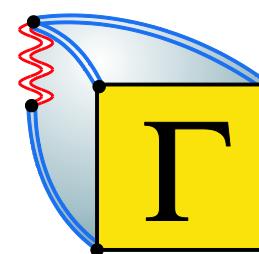
d-, p-wave superconductivity, pseudogaps,
(para-)magnons, quantum criticality



Hettler et al.'98, Lichtenstein, Katsnelson'00

Kotliar et al.'01, Potthoff et al.'03, Maier et al.'05

diagrammatic extensions



DΓA: Toschi, Katanin, KH'07

DF: Rubtsov et al.'08

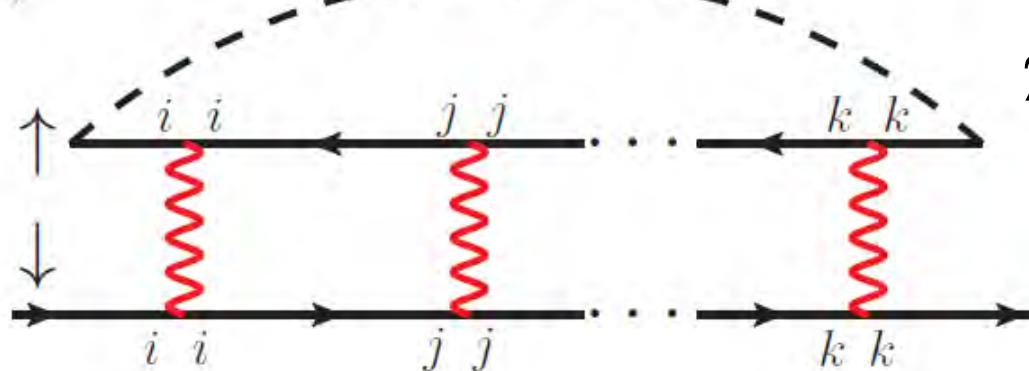
1PI, TRILEX ...

RMP'18

Take spin fluctuations

(a)

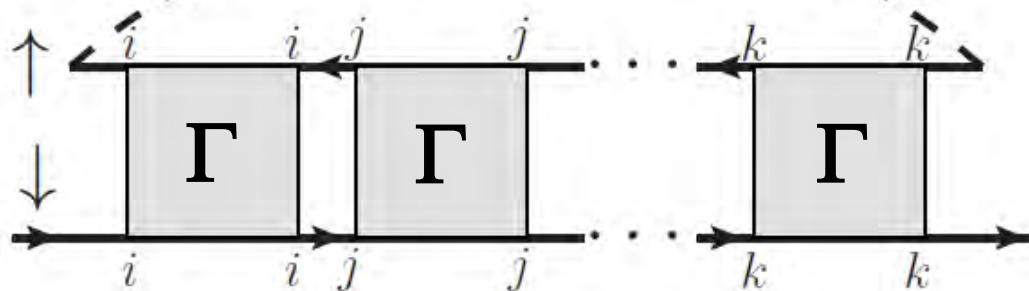
magnon self-energy Edwards-Hertz
Hertz-Millis-Moriya theory for QCP



$$\chi^{q,\omega} = \chi_0^{q,\omega} / (1 - \mathbf{U} \chi_0^{q,\omega})$$

(b)

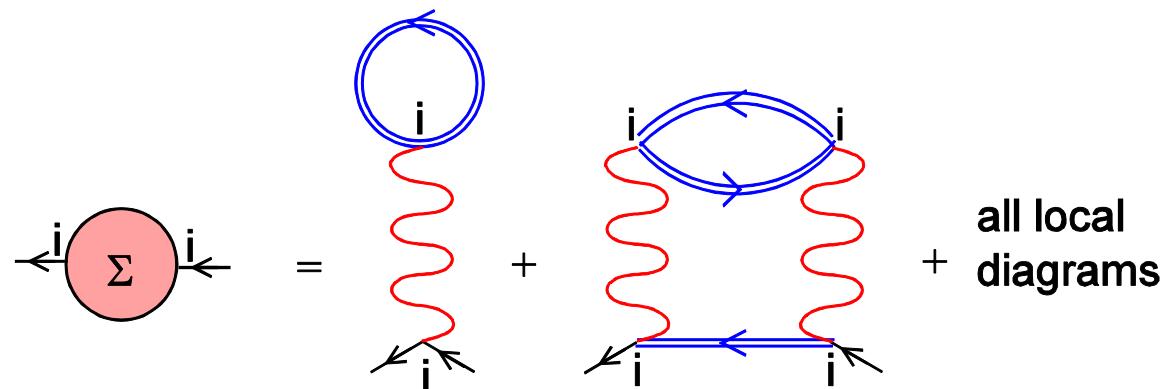
magnons but now including
all local DMFT physics



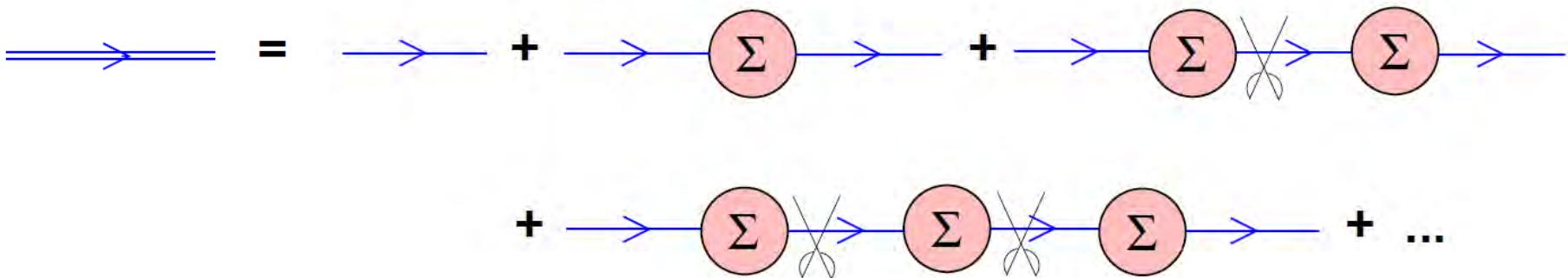
Resummation of Feynman diagrams in terms of locality
 n-particle fully irreducible vertex *approximated as local*

n=1: DMFT

- local 1-particle fully irreducible vertex Σ



Σ : one-particle irreducible one-particle vertex



Resummation of Feynman diagrams in terms of **locality**
 n-particle fully irreducible vertex ***approximated as local***

n=1: DMFT

- local 1-particle fully irreducible vertex Σ
- local 2-particle fully irreducible vertex Λ

n=2: DΓA

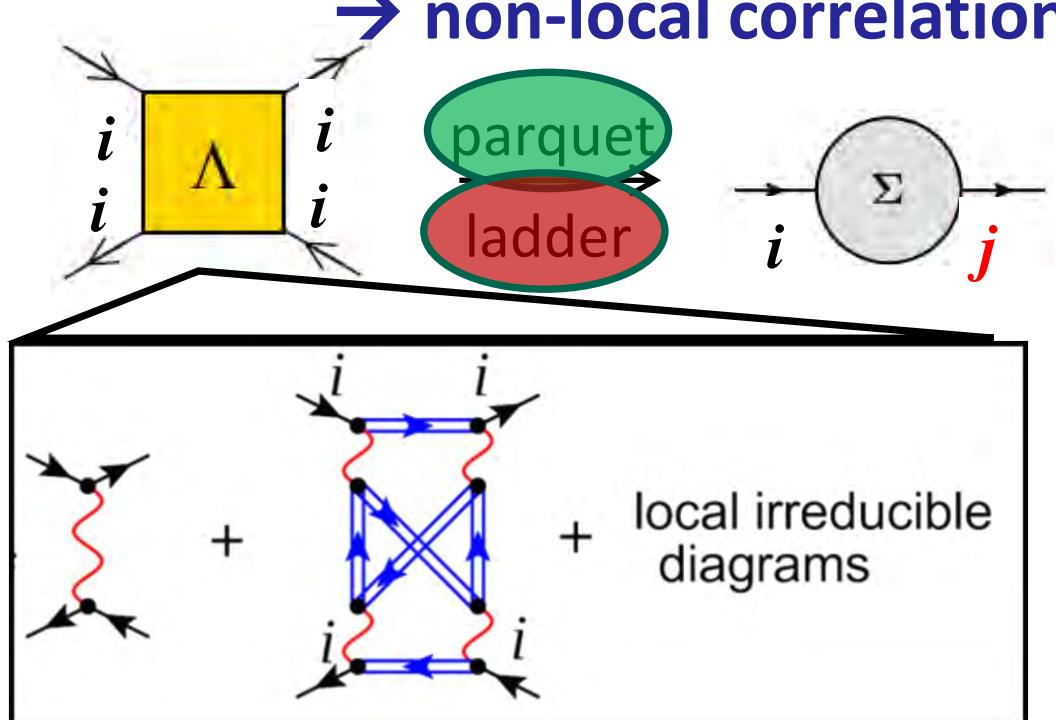
→ non-local correlations

n=3: error estimate

Ribic et al. PRB'17, 18

...

n → ∞: exact

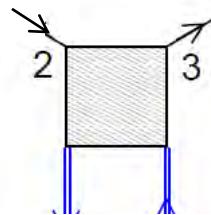


Two-particle irreducibility

$$\begin{array}{c} \text{2} \\ \text{1} \end{array} \begin{array}{c} \text{3} \\ \text{4} \end{array} \chi = - \begin{array}{c} \text{2} \\ \text{1} \end{array} \begin{array}{c} \text{3} \\ \text{4} \end{array} - \begin{array}{c} \text{2} \\ \text{1} \end{array} \begin{array}{c} \text{3} \\ \text{4} \end{array} \text{F}$$

full 2-particle vertex **F**

3 ways (channels) to cut two lines



$$\begin{array}{c} \text{2} \\ \text{1} \end{array} \begin{array}{c} \text{3} \\ \text{4} \end{array} \text{F} = \begin{array}{c} \text{2} \\ \text{1} \end{array} \begin{array}{c} \text{3} \\ \text{4} \end{array} \Lambda + \begin{array}{c} \text{2} \\ \text{1} \end{array} \begin{array}{c} \text{3} \\ \text{4} \end{array} \Phi_{\text{ph}} + \begin{array}{c} \text{2} \\ \text{1} \end{array} \begin{array}{c} \text{3} \\ \text{4} \end{array} \Phi_{\overline{\text{ph}}} + \begin{array}{c} \text{2} \\ \text{1} \end{array} \begin{array}{c} \text{3} \\ \text{4} \end{array} \Phi_{\text{pp}}$$

full vertex Λ Φ_{ph} $\Phi_{\overline{\text{ph}}}$ Φ_{pp}

parquet equation

Parquet equations

$$\begin{array}{c} 2 \\ \swarrow \searrow \\ 1 & 4 \\ \boxed{\Phi_{ph}} \\ \end{array} = \begin{array}{c} 2 \\ \swarrow \searrow \\ 1 & 4 \\ \boxed{F} & \xrightarrow{\quad} & \boxed{\Lambda} \\ \xleftarrow{\quad} & & \xleftarrow{\quad} \\ \end{array} + \begin{array}{c} 2 \\ \swarrow \searrow \\ 1 & 4 \\ \boxed{F} & \xrightarrow{\quad} & \boxed{\Phi_{ph}} \\ \xleftarrow{\quad} & & \xleftarrow{\quad} \\ \end{array} + \begin{array}{c} 2 \\ \swarrow \searrow \\ 1 & 4 \\ \boxed{F} & \xrightarrow{\quad} & \boxed{\Phi_{pp}} \\ \xleftarrow{\quad} & & \xleftarrow{\quad} \\ \end{array}$$

$$\begin{array}{c} 2 \\ \swarrow \searrow \\ 1 & 4 \\ \boxed{\Phi_{ph}} \\ \end{array} = \begin{array}{c} 2 \\ \swarrow \searrow \\ 1 & 4 \\ \boxed{F} \\ \xleftarrow{\quad} \end{array} + \begin{array}{c} 2 \\ \swarrow \searrow \\ 1 & 4 \\ \boxed{F} \\ \xleftarrow{\quad} \end{array} + \begin{array}{c} 2 \\ \swarrow \searrow \\ 1 & 4 \\ \boxed{F} \\ \xleftarrow{\quad} \end{array}$$

$$\begin{array}{c} 1 \\ \swarrow \searrow \\ 2 & 3 \\ \boxed{\Phi_{ph}} \\ \end{array} + \begin{array}{c} 1 \\ \swarrow \searrow \\ 2 & 3 \\ \boxed{\Phi_{ph}} \\ \end{array} + \begin{array}{c} 1 \\ \swarrow \searrow \\ 2 & 3 \\ \boxed{\Phi_{pp}} \\ \end{array}$$

$$\begin{array}{c} 2 \\ \swarrow \searrow \\ 1 & 4 \\ \boxed{\Phi_{pp}} \\ \end{array} = \begin{array}{c} 2 \\ \swarrow \searrow \\ 1 & 4 \\ \boxed{F} & \xrightarrow{\quad} & \boxed{\Lambda} \\ \xleftarrow{\quad} & & \xleftarrow{\quad} \\ \end{array} + \begin{array}{c} 2 \\ \swarrow \searrow \\ 1 & 4 \\ \boxed{F} & \xrightarrow{\quad} & \boxed{\Phi_{ph}} \\ \xleftarrow{\quad} & & \xleftarrow{\quad} \\ \end{array} + \begin{array}{c} 2 \\ \swarrow \searrow \\ 1 & 4 \\ \boxed{F} & \xrightarrow{\quad} & \boxed{\Phi_{ph}} \\ \xleftarrow{\quad} & & \xleftarrow{\quad} \\ \end{array}$$

$$\begin{array}{c} 2 \\ \swarrow \searrow \\ 1 & 4 \\ \boxed{F} \\ \end{array} = \begin{array}{c} 2 \\ \swarrow \searrow \\ 1 & 4 \\ \boxed{\Lambda} \\ \xleftarrow{\quad} \end{array} + \begin{array}{c} 2 \\ \swarrow \searrow \\ 1 & 4 \\ \boxed{\Phi_{ph}} \\ \xleftarrow{\quad} \end{array} + \begin{array}{c} 2 \\ \swarrow \searrow \\ 1 & 4 \\ \boxed{\Phi_{ph}} \\ \xleftarrow{\quad} \end{array} + \begin{array}{c} 2 \\ \swarrow \searrow \\ 1 & 4 \\ \boxed{\Phi_{pp}} \\ \xleftarrow{\quad} \end{array}$$

3 Bethe-Salpeter eq.



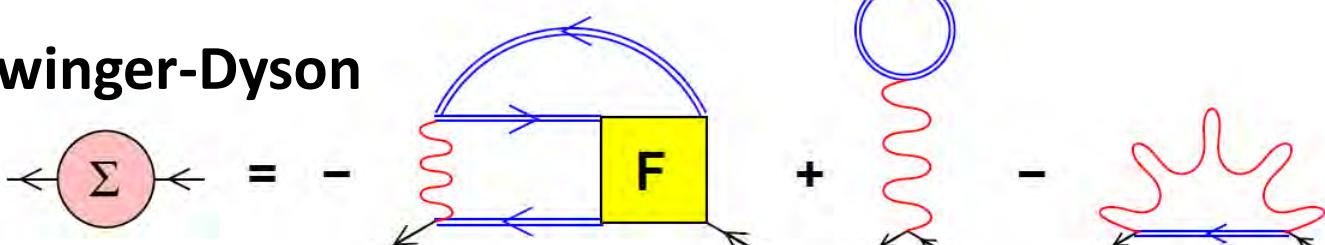
Λ given

$\rightarrow F, \Phi_{ph}, \Phi_{ph}^-, \Phi_{pp}, \Sigma, G$

$\Lambda = U$ (parquet approx.)

$\Lambda = \Lambda_{loc}$ ($D\Gamma A$)

Schwinger-Dyson

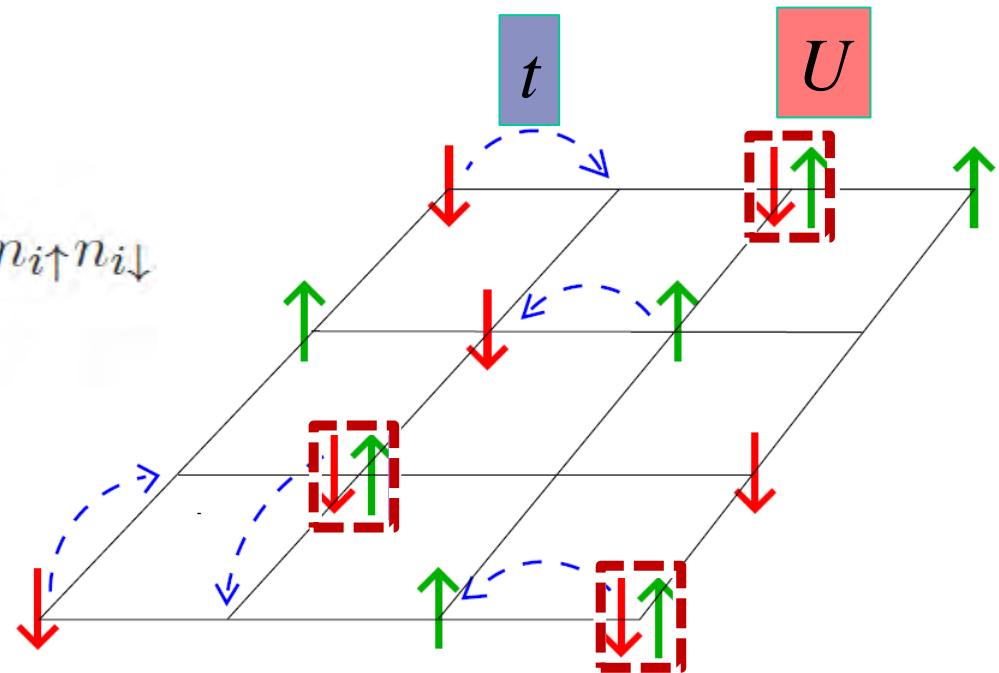


Dyson equation

$$\xrightarrow{\quad} = \xrightarrow{\quad} + \xrightarrow{\quad} \boxed{\Sigma} \xrightarrow{\quad} + \xrightarrow{\quad} \boxed{\Sigma} \diagup \diagdown \xrightarrow{\quad} \boxed{\Sigma} \xrightarrow{\quad} \dots$$

ladder DΓA results: 3D Hubbard model

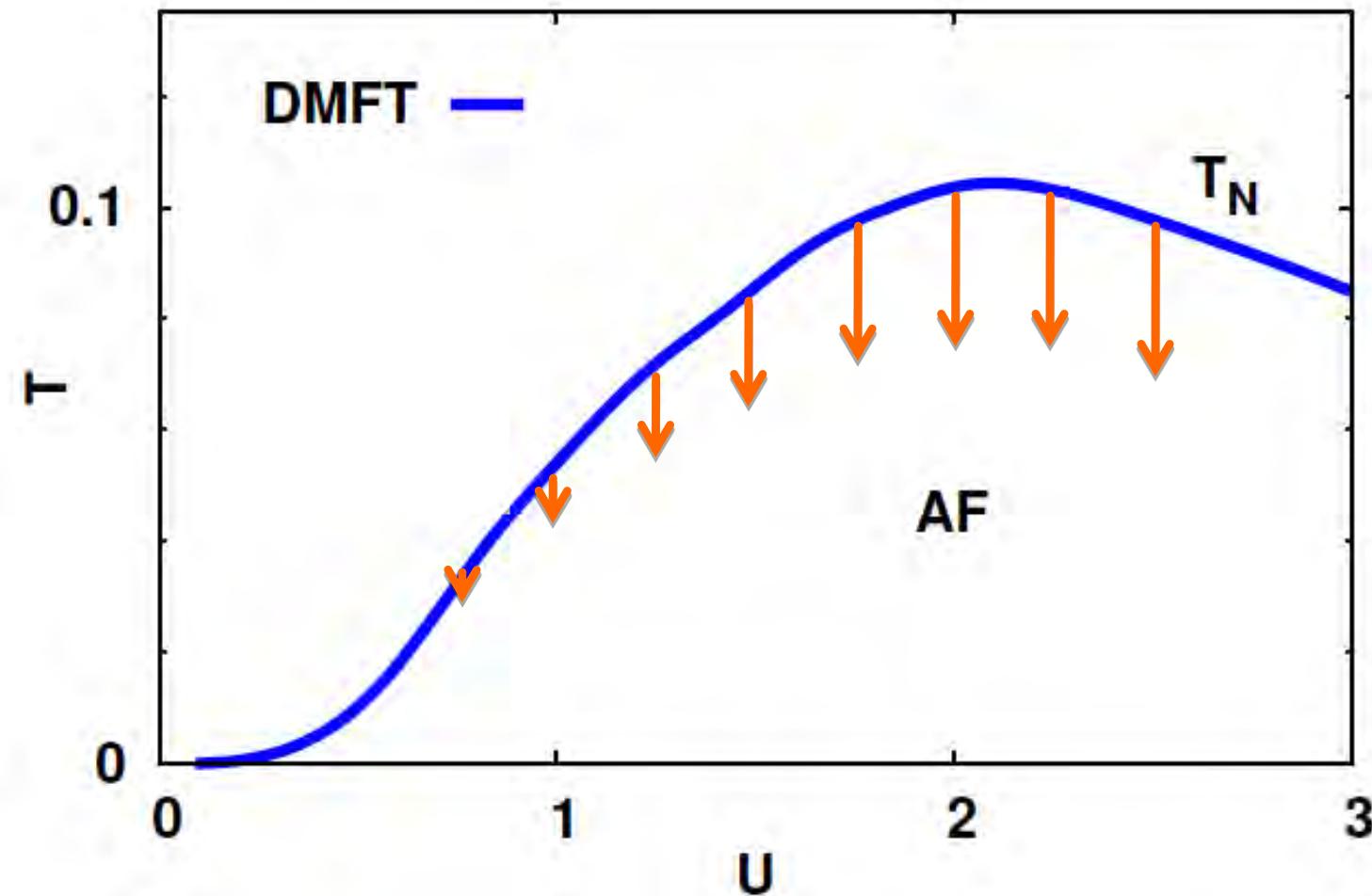
$$H = -t \sum_{\langle ij \rangle \sigma} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$



DΓA results in 3 dimensions

Phase diagram: Hubbard model in $d=3$ (cubic lattice $D=1$, half-filling)

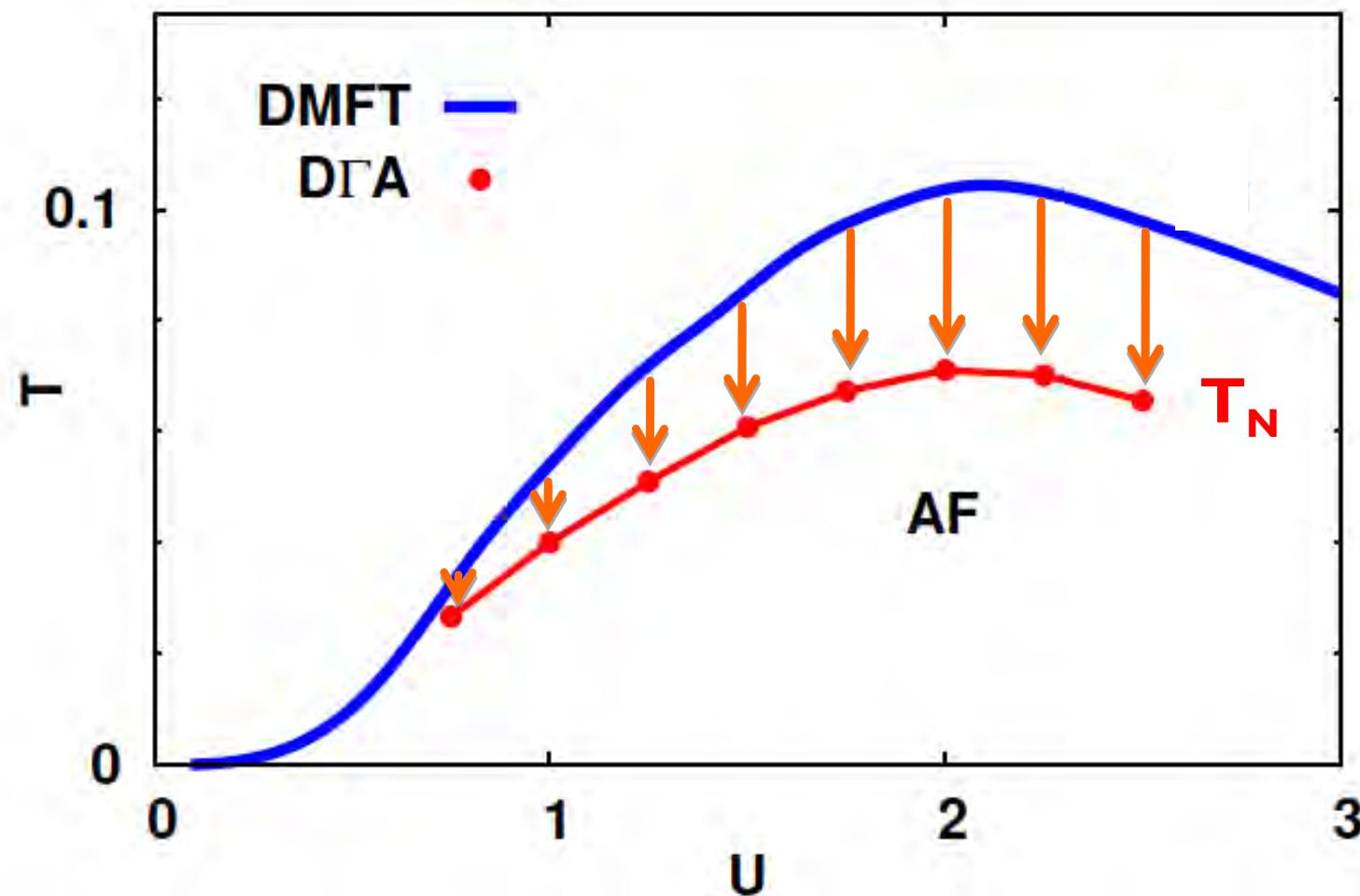
G. Rohringer et al., PRL (2011)



DΓA results in 3 dimensions

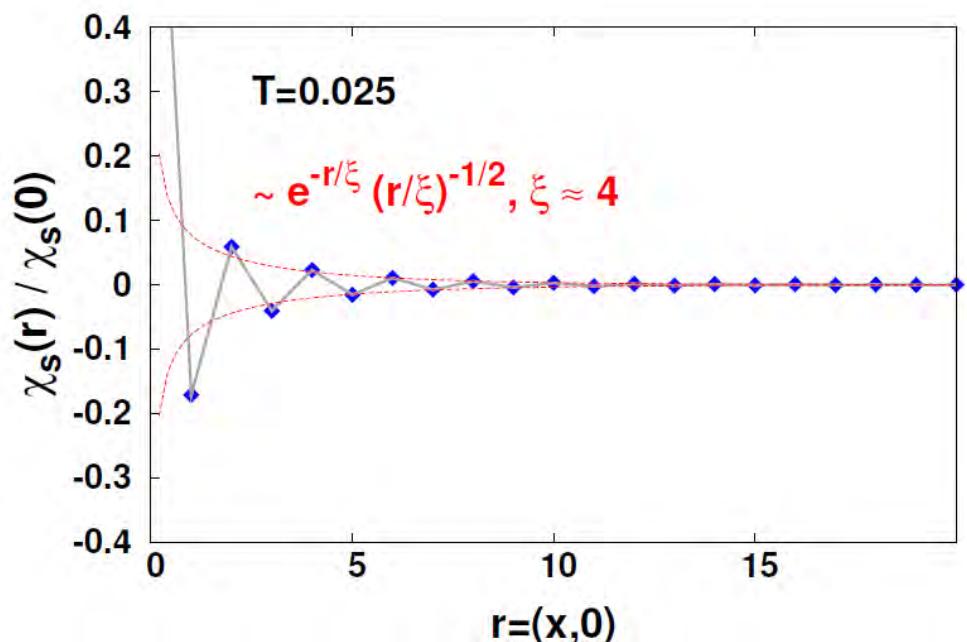
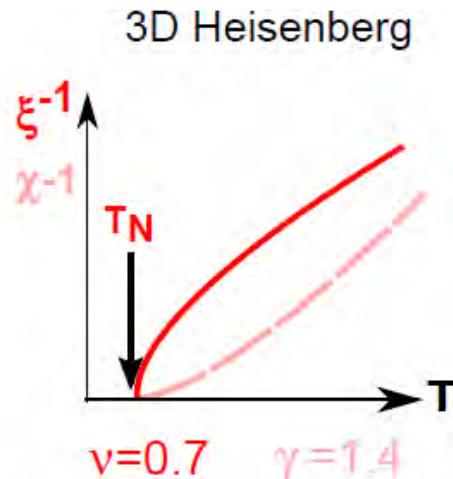
Phase diagram: Hubbard model in $d=3$ (cubic lattice $D=1$, half-filling)

G. Rohringer et al., PRL (2011)



Intermezzo: critical exponents

Critical exponents



susceptibility $\chi^{-1} \sim (T-T_c)^\gamma$

correlation length $\xi^{-1} \sim (T-T_c)^v$

$$\chi_s(r) = \langle S(r)S(0) \rangle \sim (r/\xi)^{-1/2} e^{-r/\xi}$$

Fisher'67 relation

$$\gamma/v = 2 - \eta \quad (\eta \sim 0)$$

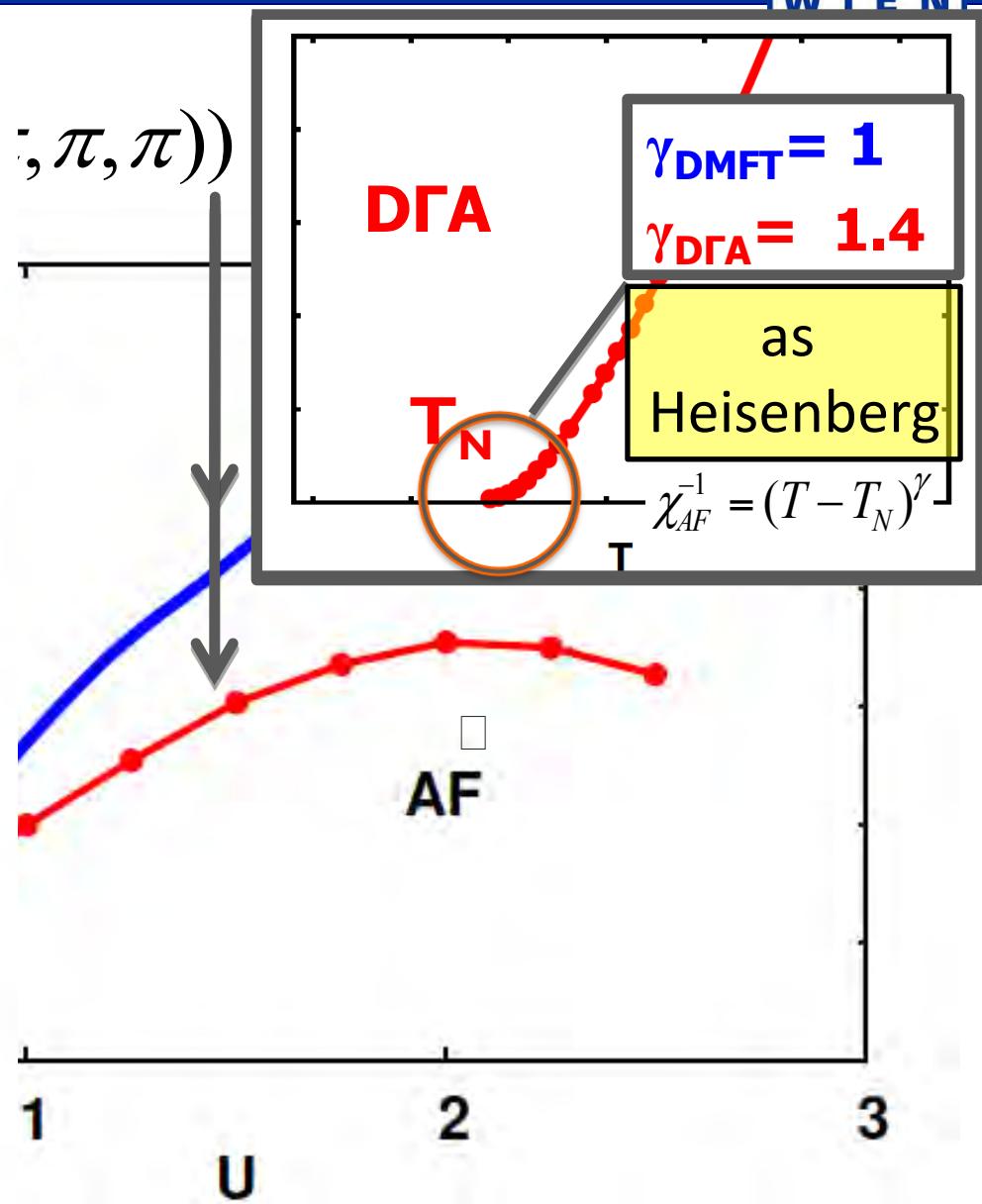
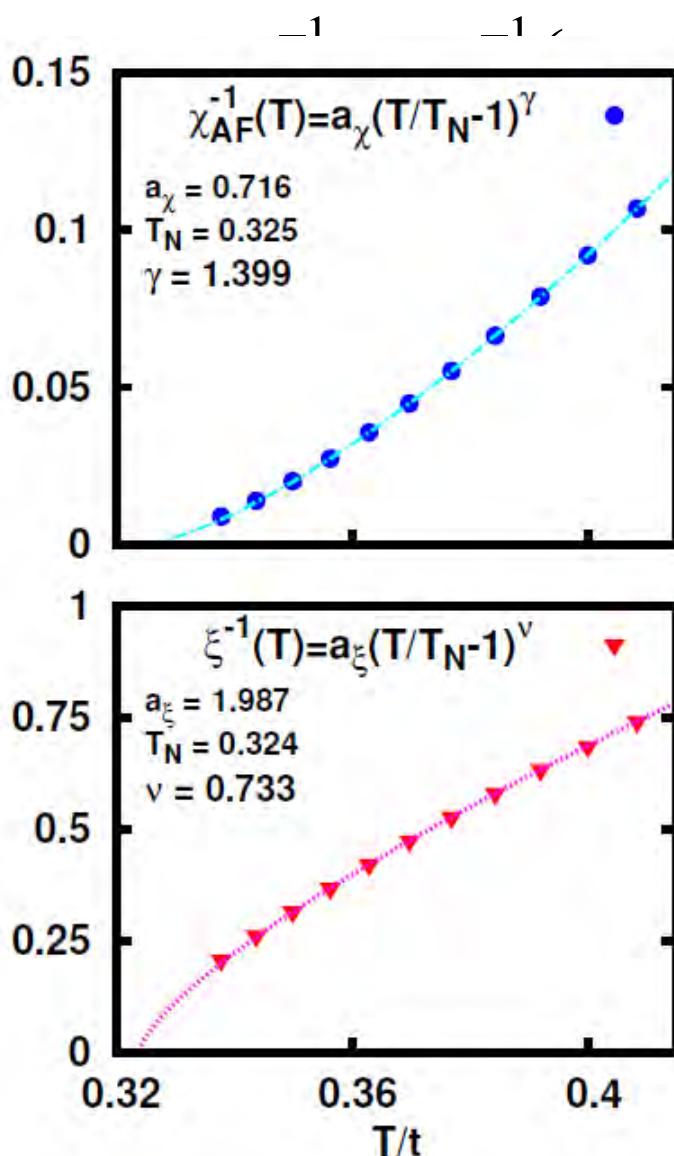
in practice: Ornstein-Zernike'16

$$\chi_{\omega=0}^{\omega=0} \sim 1/[(q-Q)^2 + \xi^{-2}]$$

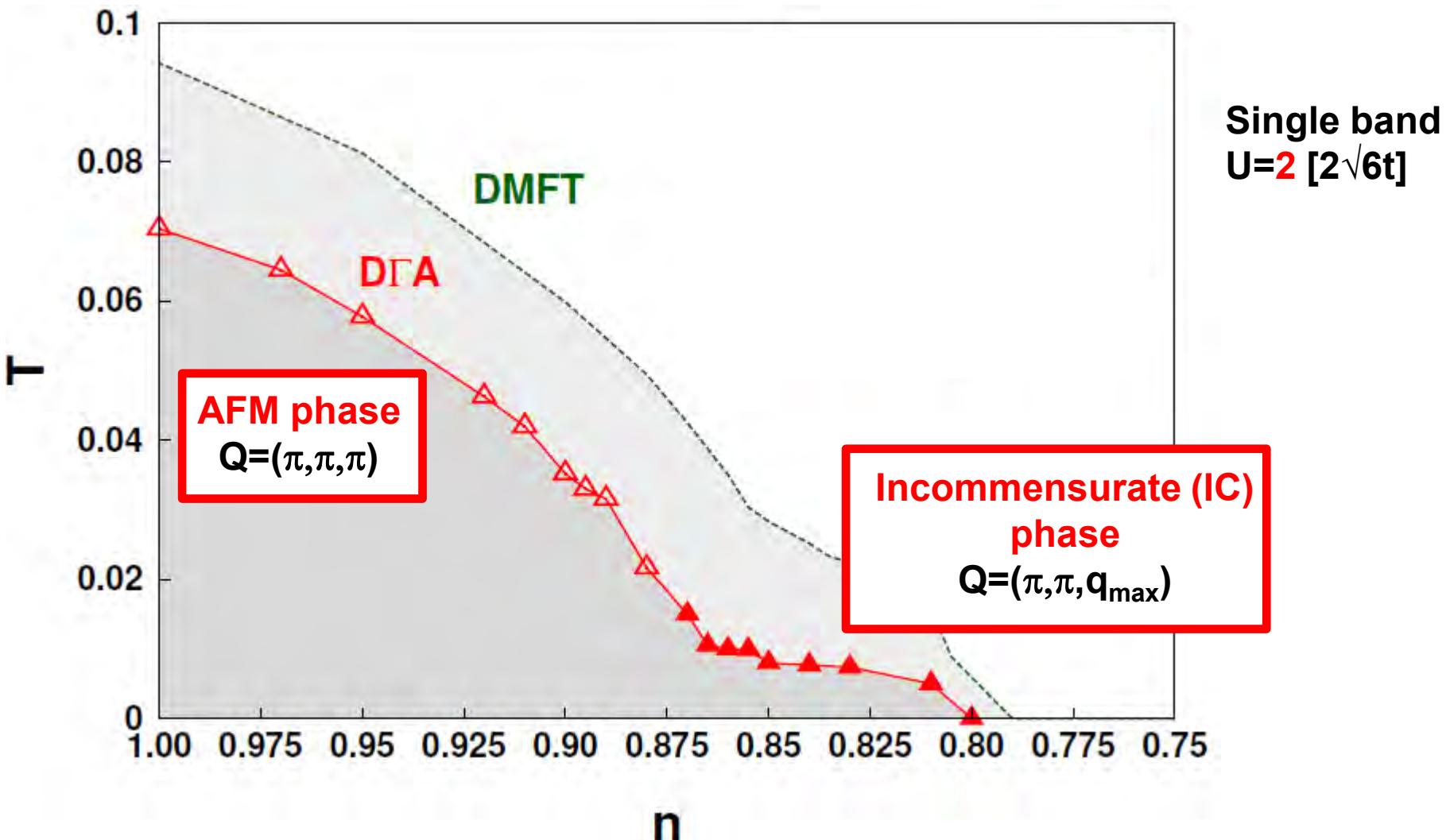
$$\rightarrow \chi = \chi_{\omega=0}^{\omega=0}, \xi$$

universality: γ, v only depend on d, symmetry order param.

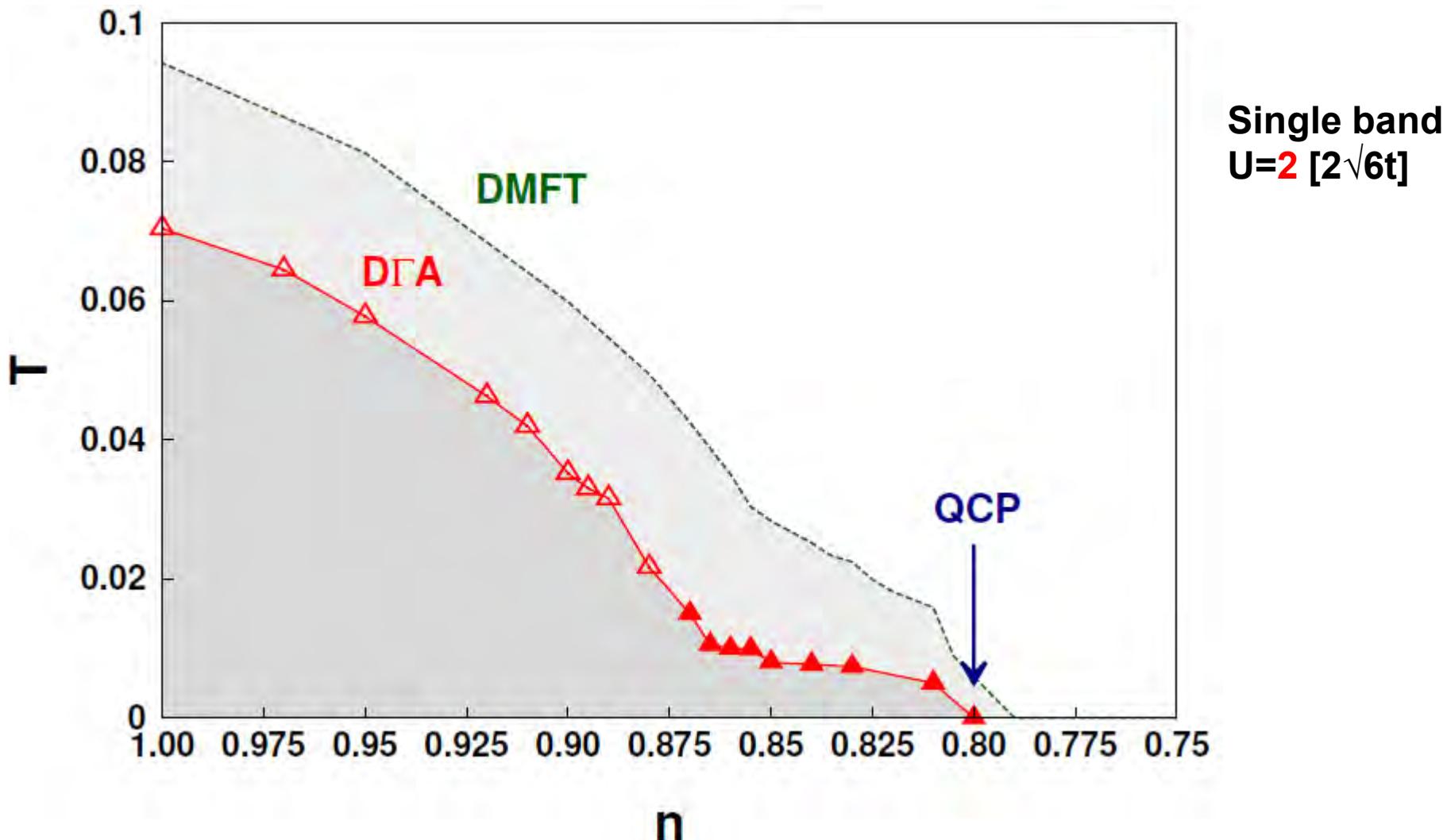
DΓA results in 3 dimensions



Quantum phase transition in 3 dimensions

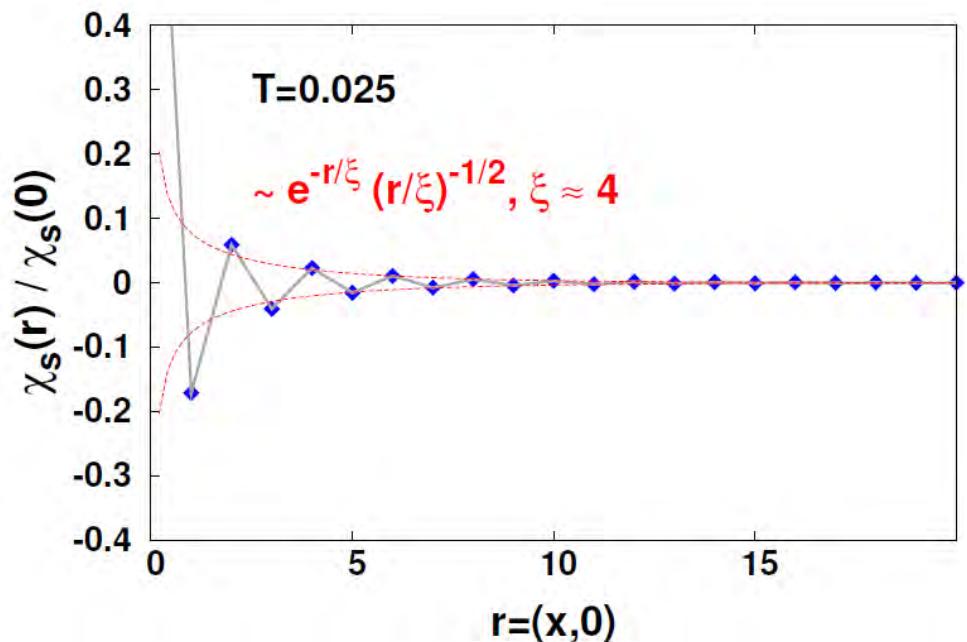
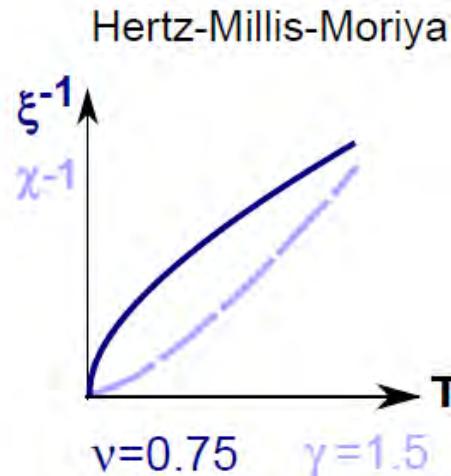


Quantum phase transition in 3 dimensions



Intermezzo: quantum critical exponents

Critical exponents



susceptibility $\chi^{-1} \sim (T^v - T_c)^y$

$$d_{\text{eff}} = z + d$$

correlation length $\xi^{-1} \sim (T^v - T_c)^v$

$z=2$, AF metal

QCP quantum fluctuations in $\tau \in [0, 1/T]$

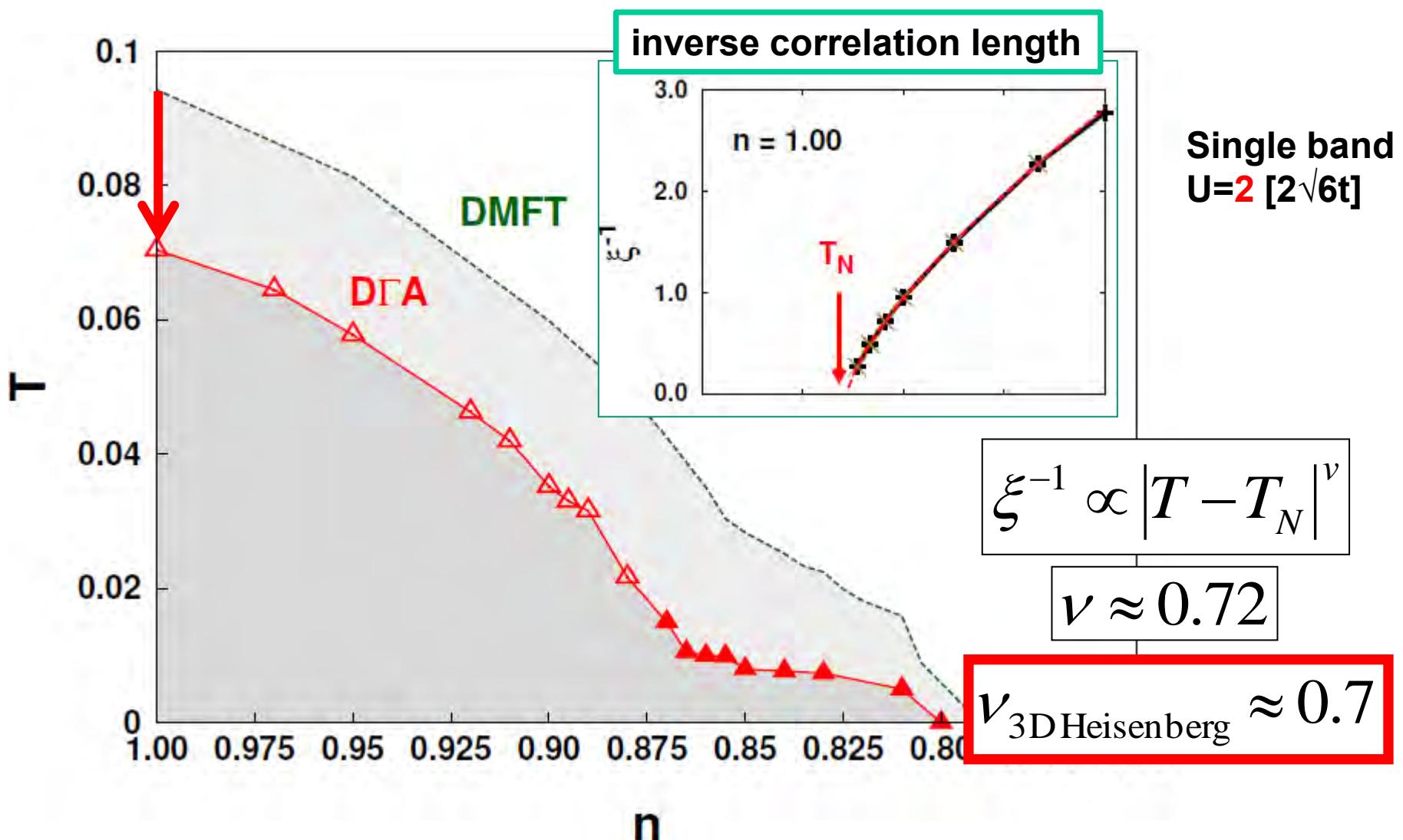
$z=1$, AF insulator

correlation length in time $\xi_\tau^{-1} \sim T^{zv}$

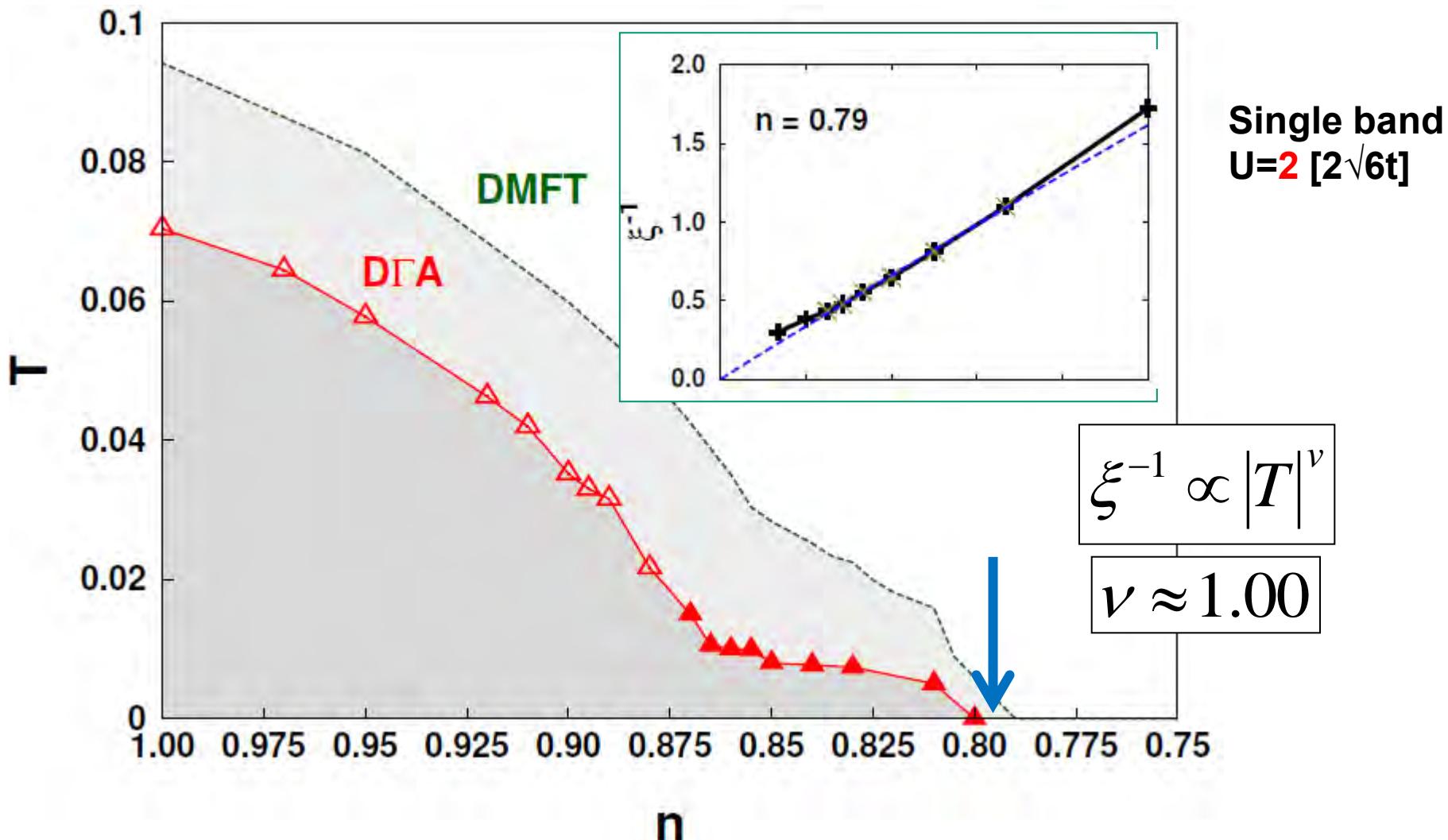
Chubukov, Sachdev, Ye'94

Troyer, Imada, Ueda'97

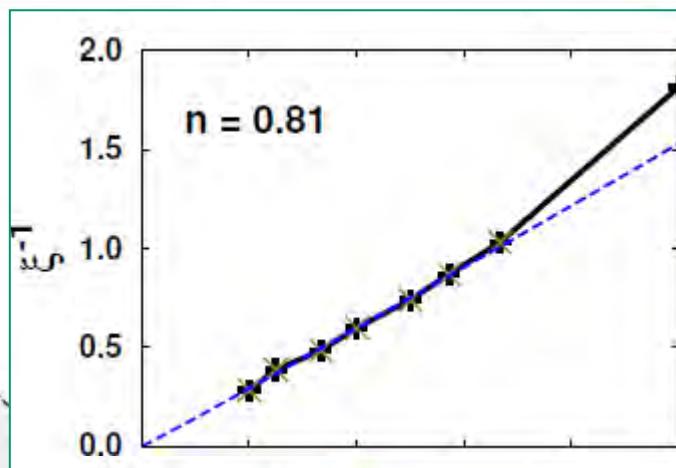
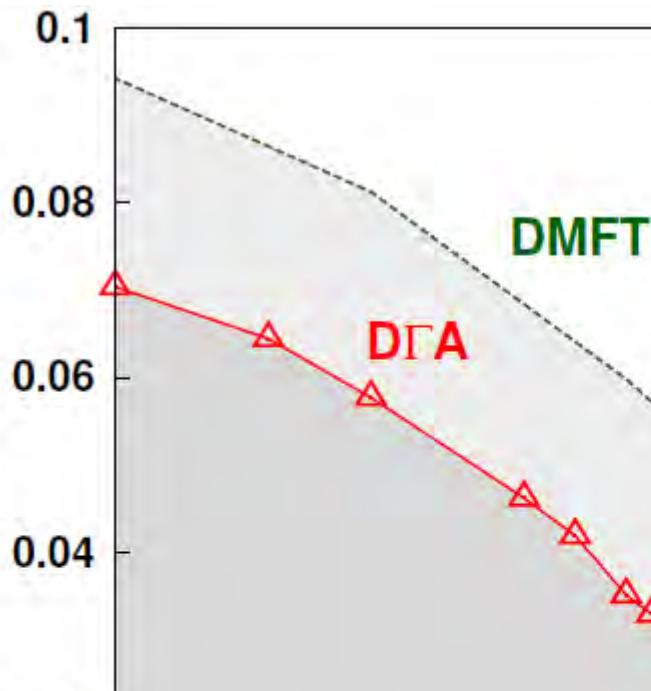
Quantum phase transition in 3 dimensions



Quantum phase transition in 3 dimensions



Quantum phase transition in 3 dimensions



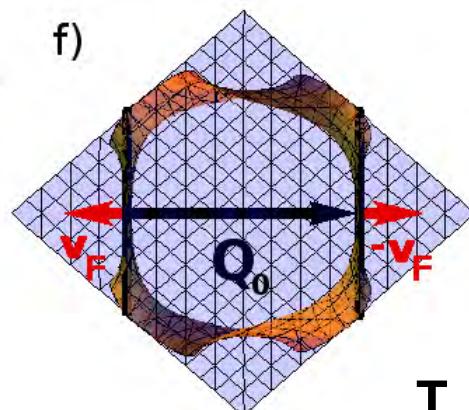
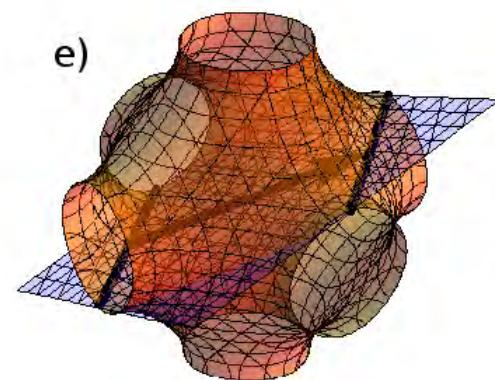
Single band
 $U=2$ [$2\sqrt{6}t$]

$$\xi^{-1} \propto |T|^\nu$$

$$\nu_{\text{Kohn/QCP}} = 1.0$$

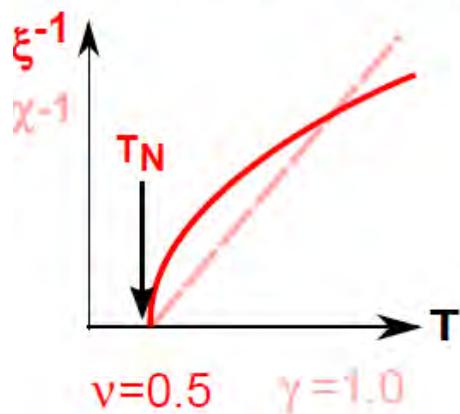
$$d_{\text{eff}} = z + d = 2 + 3 > 4$$

$$\nu_{\text{Hertz-Millis-Moriya}} = 0.75$$

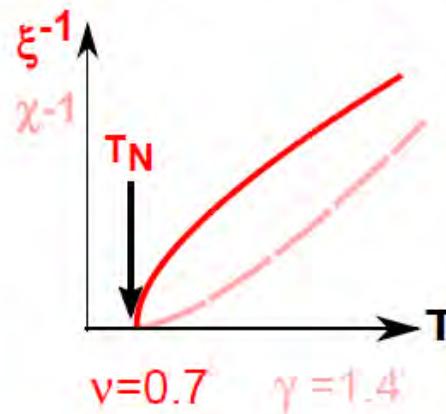


(Quantum) critical exponents

Gaussian (bosonic MF)



3D Heisenberg



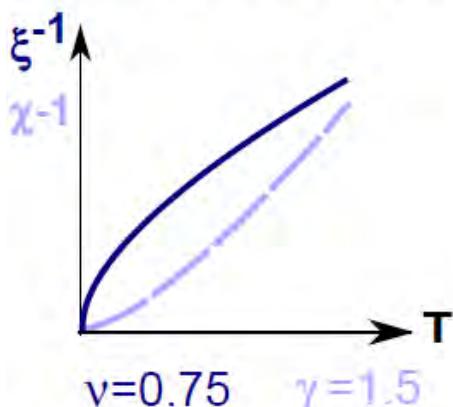
correlation length

$$\xi^{-1} \sim T^v$$

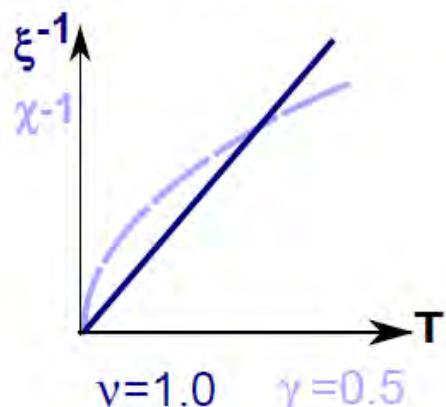
susceptibility

$$\chi^{-1} \sim T^\gamma$$

Hertz-Millis-Moriya



Kohn - QCP



Kohn-QCP:

strong deviation from

$$\gamma=2v \quad (\eta \sim 0)$$

- **Quantum criticality in 3D Hubbard model** $d_{\text{eff}} = z + d = 2 + 3$
- **Quantum criticality in 2D periodic Anderson model** $d_{\text{eff}} = 1 + 2$

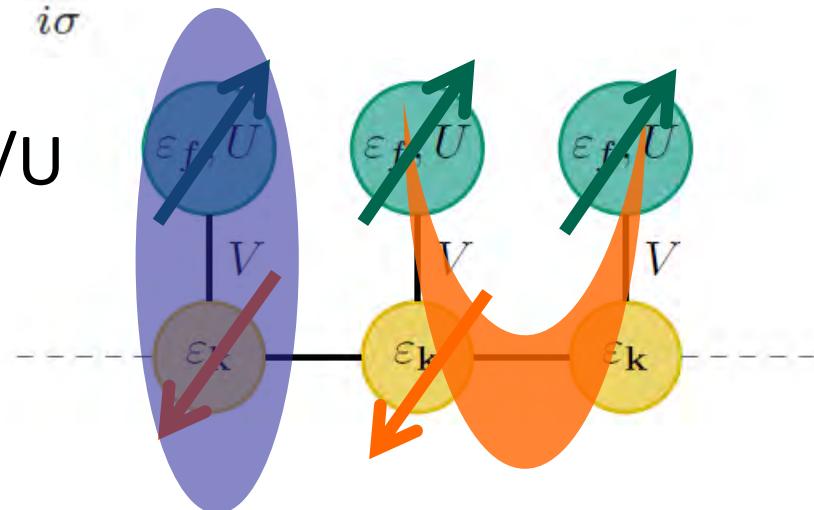
Quantum criticality 2D periodic Anderson model

$$\begin{aligned}
 H_{PAM} = & \sum_{\mathbf{k}\sigma} \varepsilon_{\mathbf{k}} a_{\mathbf{k}\sigma}^+ a_{\mathbf{k}\sigma} + \varepsilon_f \sum_{i\sigma} f_{i\sigma}^+ f_{i\sigma} \\
 & + \sum_{\mathbf{k}\sigma} V_{\mathbf{k}} (a_{\mathbf{k}\sigma}^+ f_{\mathbf{k}\sigma} + f_{\mathbf{k}\sigma}^+ a_{\mathbf{k}\sigma}) \\
 & + U \sum_i f_{i\sigma}^+ f_{i\sigma} f_{i\bar{\sigma}}^+ f_{i\bar{\sigma}} - \mu \sum_{i\sigma} (f_{i\sigma}^+ f_{i\sigma} + a_{i\sigma}^+ a_{i\sigma}),
 \end{aligned} \tag{4}$$

Effective spin interaction $J = 4V^2/U$

Competition between

- 1) Kondo $T_K \sim e^{-1}/(\rho_0 J)$
- 2) RKKY $T_c \sim \chi_0 J^2$



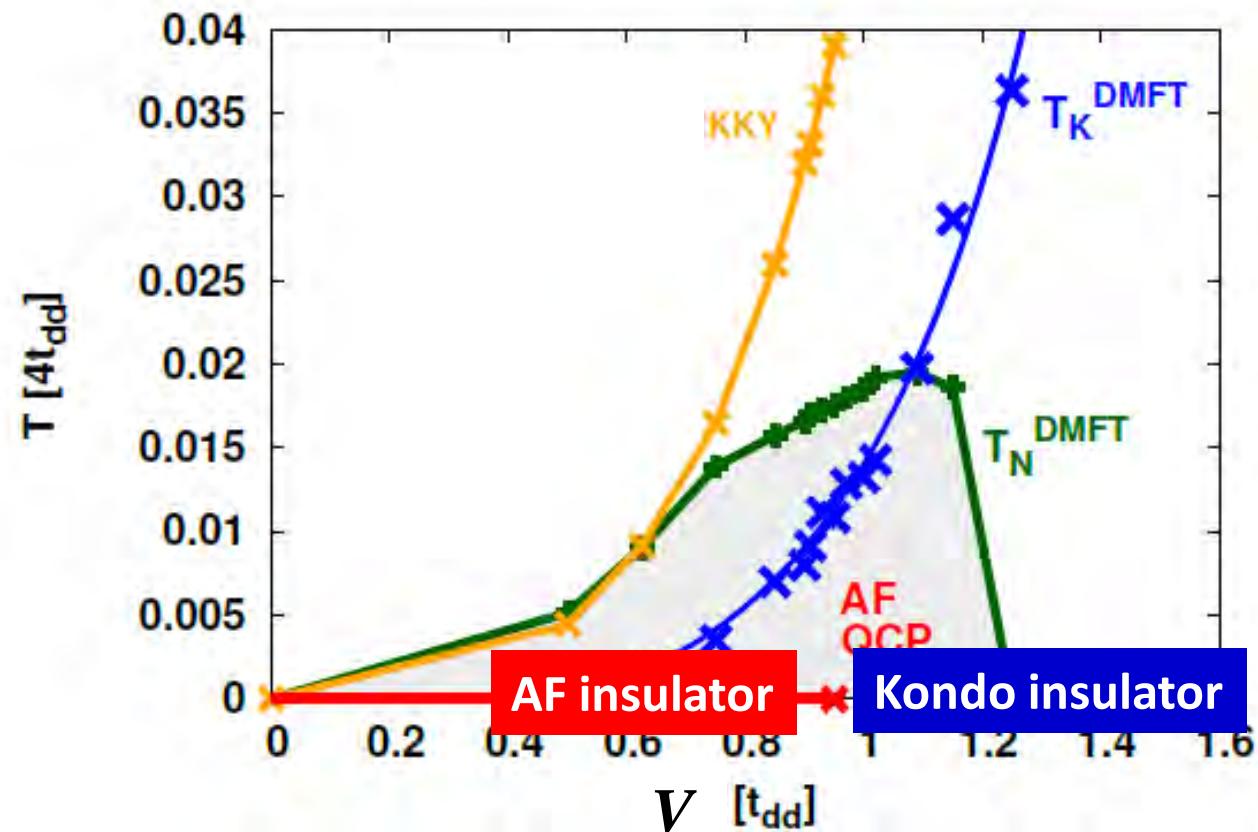
in the following: half-filling (Kondo insulator)

Doniach phase diagram in DMFT and DΓA (competition between Kondo and RKKY)

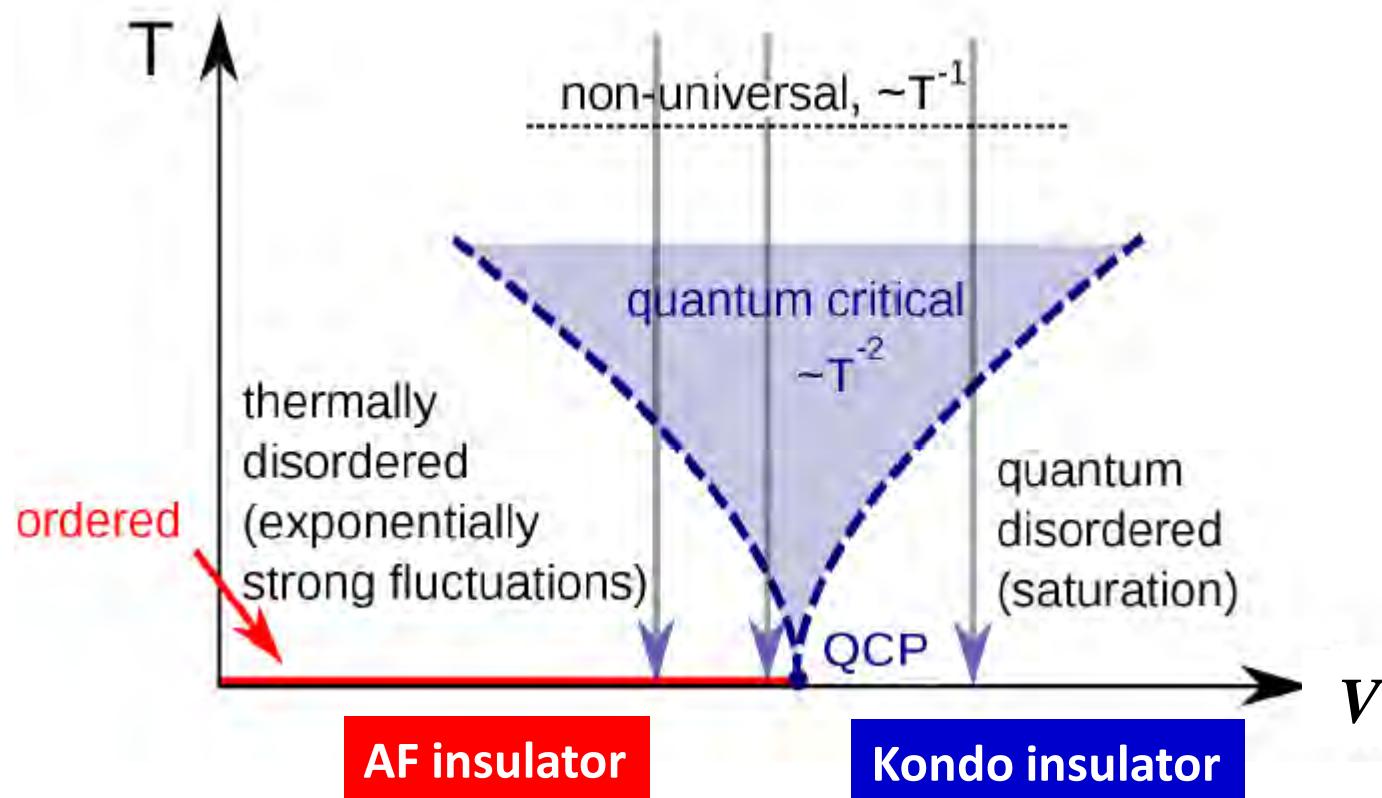
Kondo $T_K \sim e^{-1}/(\rho_0 J)$

$$J = 4V^2/U \quad U=4t_{dd}$$

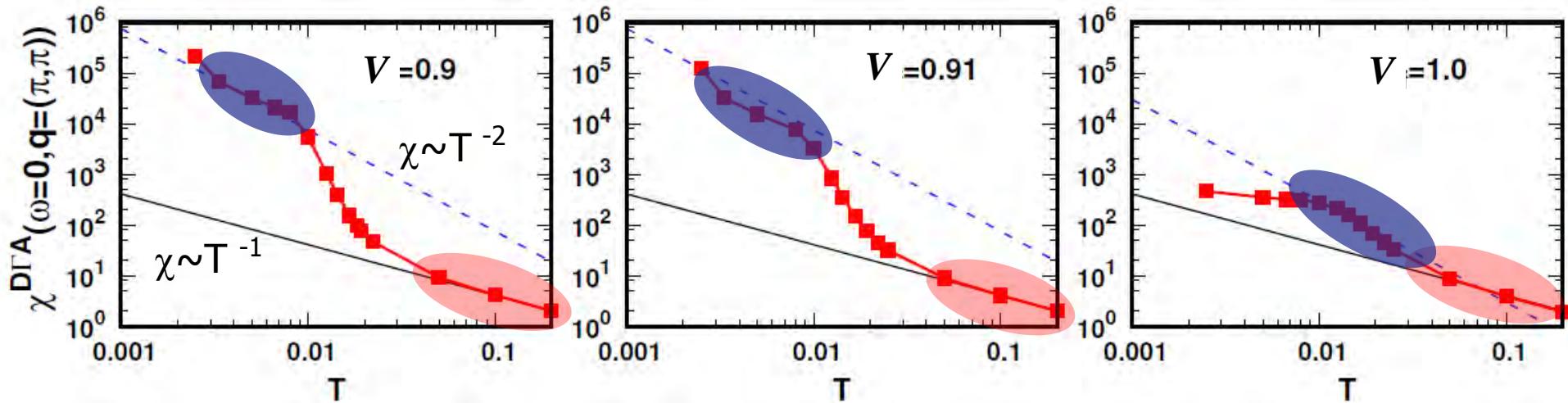
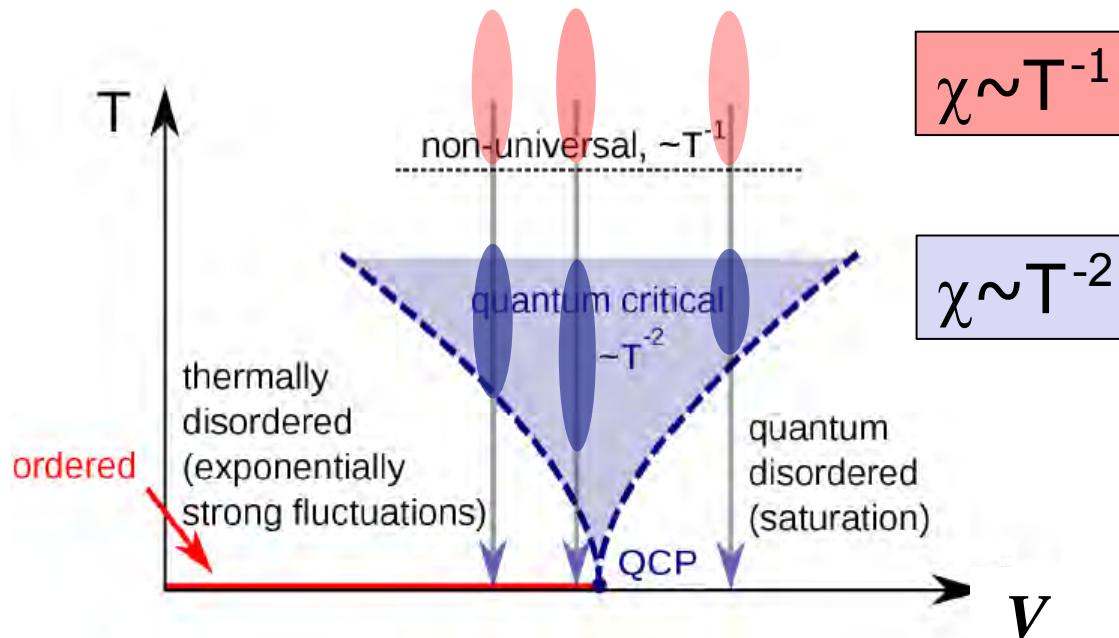
RKKY $T_c \sim \chi_0 J^2$



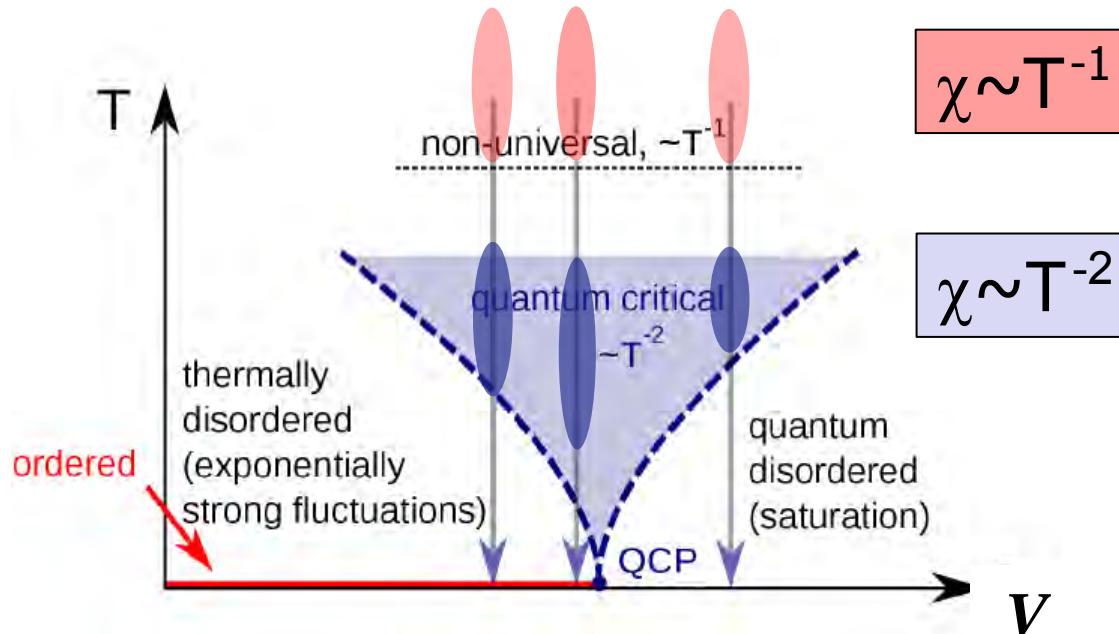
Quantum criticality 2D periodic Anderson model



Quantum criticality 2D periodic Anderson model



Quantum criticality 2D periodic Anderson model



correlation length in time $\xi_\tau^{-1} \sim T^{z\nu}$, $z=1$, correlation length $\xi^{-1} \sim T^\nu$

finite $T \rightarrow$ cut-off $\xi_\tau \sim 1/T \rightarrow \xi \sim \xi_\tau \sim 1/T$ i.e. $\nu=1$

Fisher'67 relation $\gamma/\nu=2 - \eta$ ($\eta \sim 0$) $\rightarrow \gamma = 2$ i.e. $\chi \sim T^{-2}$

same as for Heisenberg model Chakravarty et al.'88

Superconductivity in 2D



Superconductivity in 2D

2D Hubbard model (square lattice)

Kitatani et al.'2018

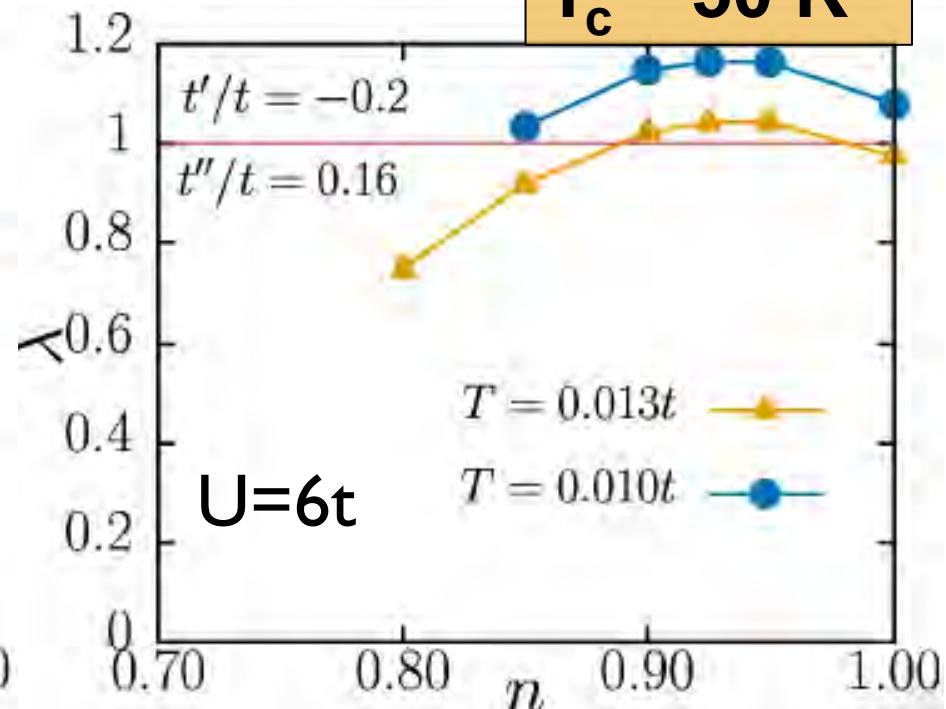
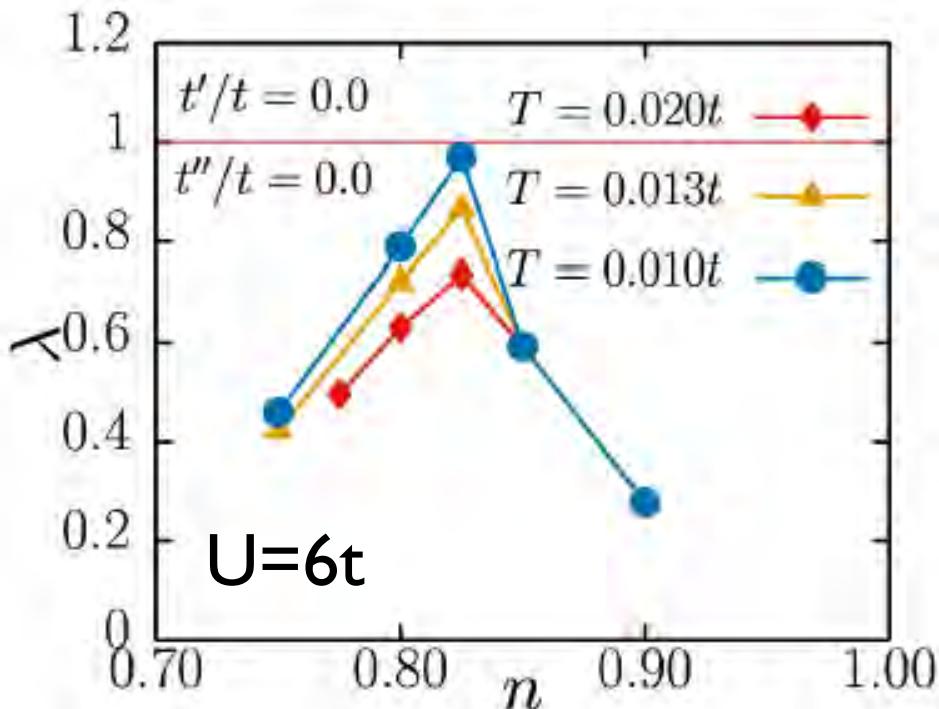
Non-local spin fluctuations in ladder DΓA (ph, $\overline{\text{ph}}$ channel)

→ remove local pp contribution → Γ_{pp}
→ Eliashberg equation (pp channel)

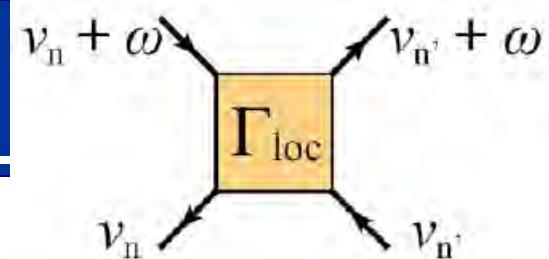
$$\chi = \chi_0 / (1 + \Gamma_{\text{pp}} \chi_0)$$

SC dome
 $\text{eigenvalue} - \lambda$
 $T_c \sim 50 \text{ K}$

Leading d-wave eigenvalue λ

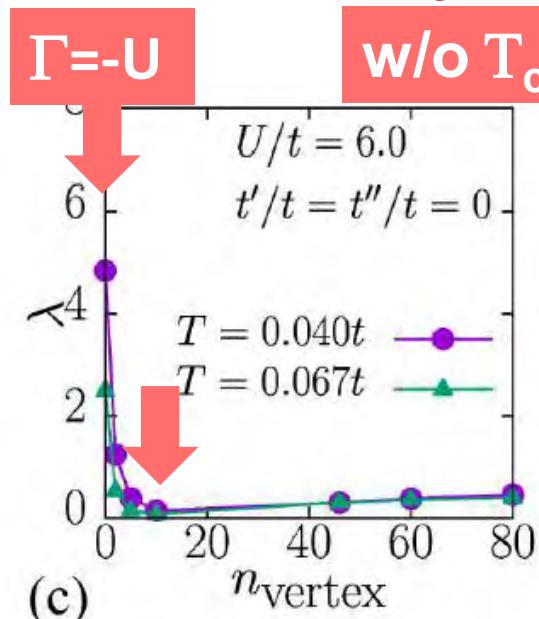


Superconductivity in 2D



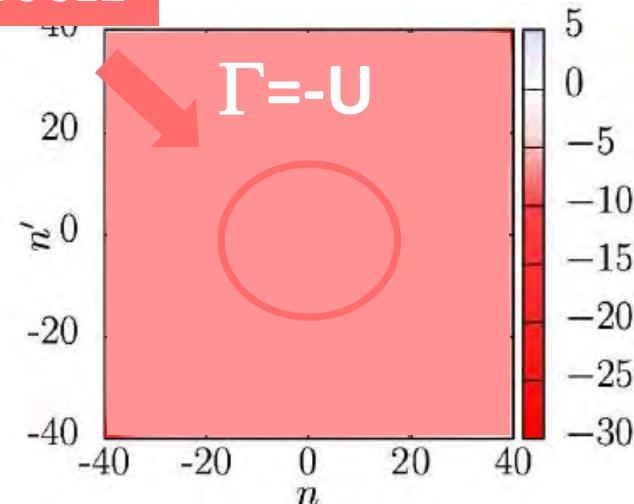
Importance of vertex dynamics!

$U=6t$

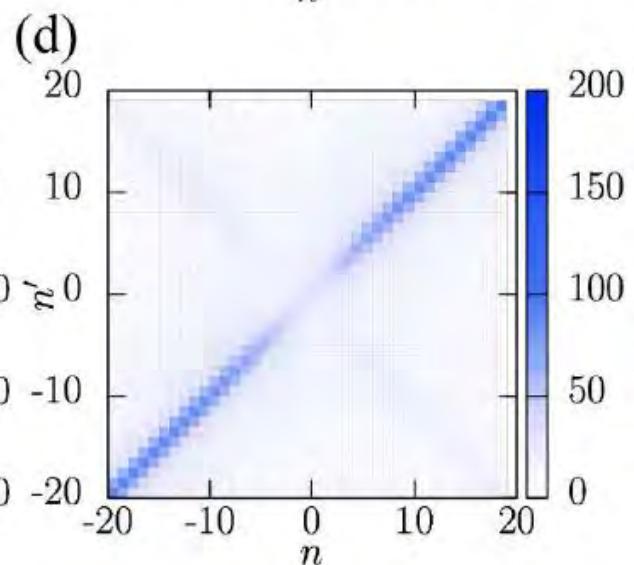
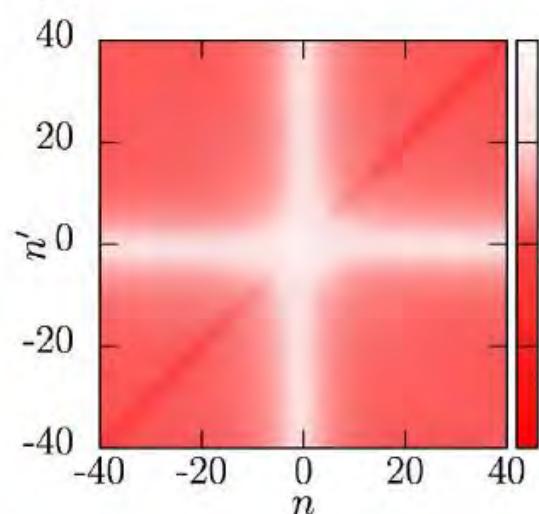


w/o $T_c \sim 400K$

local $\Gamma_m(v_n, v_{n'}, \omega)$



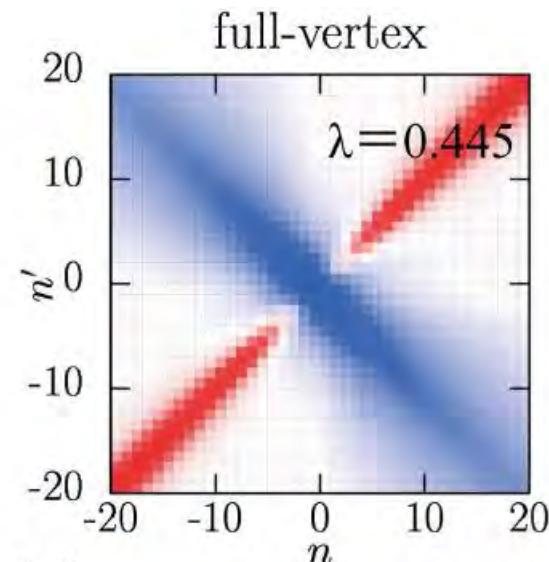
nonlocal
F
 $q = (\pi, \pi)$



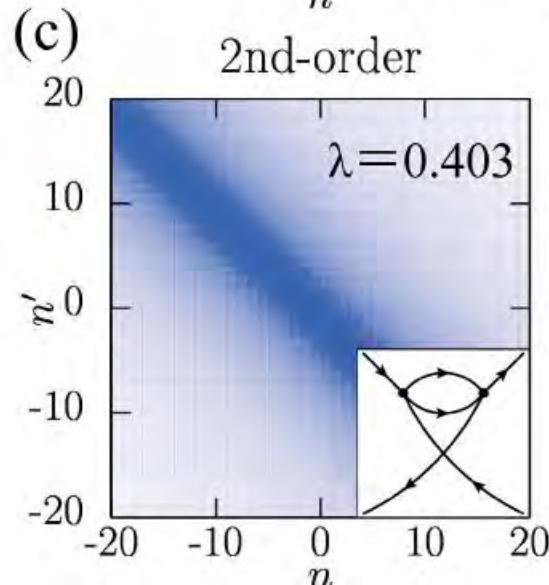
Superconductivity in 2D

Importance of **vertex dynamics!**

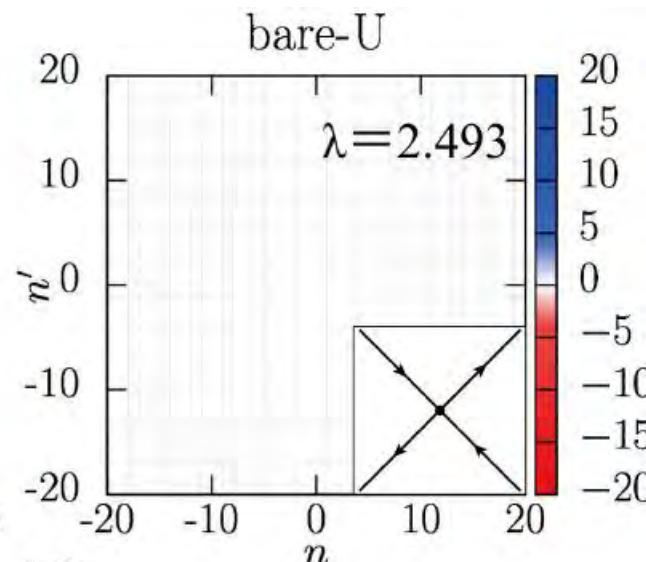
local
 $\Gamma_m(-U)$



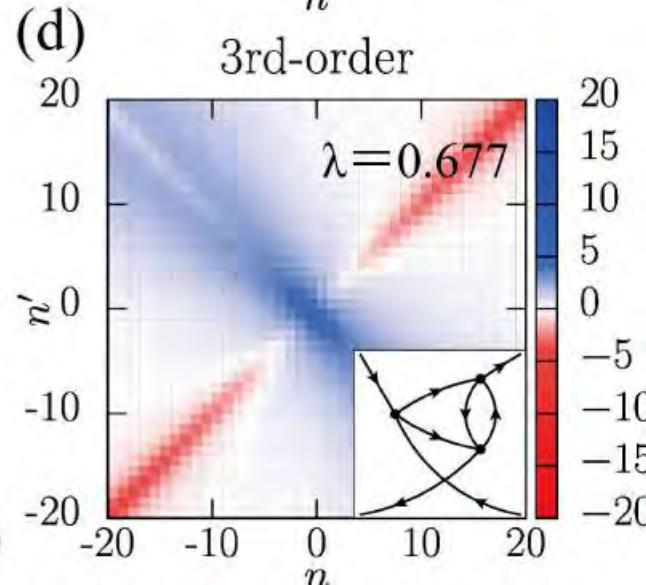
2nd order



bare U



3rd order

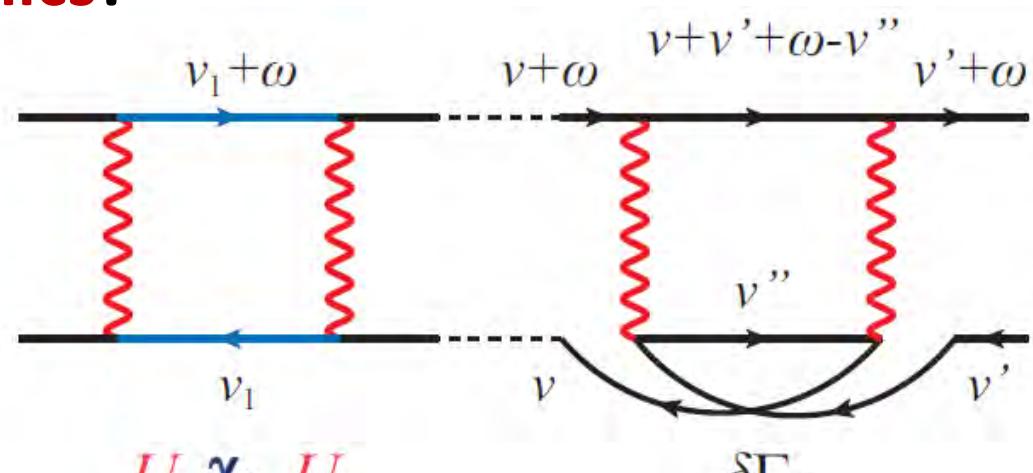


Superconductivity in 2D

Importance of **vertex dynamics!**

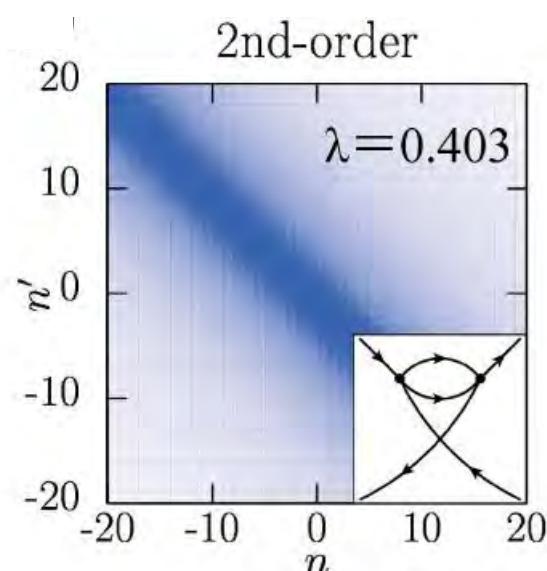
spin fluctuations

$$\chi = \chi_0 + \chi_0 U \chi_0 + \chi_0 U \chi_0 U \chi_0 + \dots$$
$$= \chi_0 / (1 - \chi_0 U)$$



suppressed by pp channel $\delta\Gamma_m$

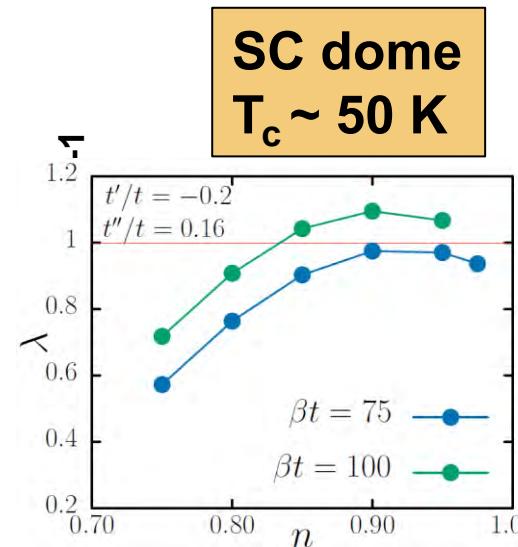
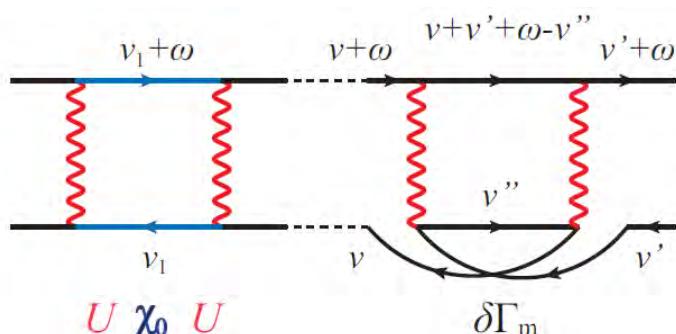
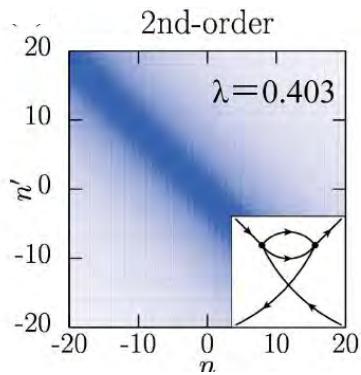
2nd order



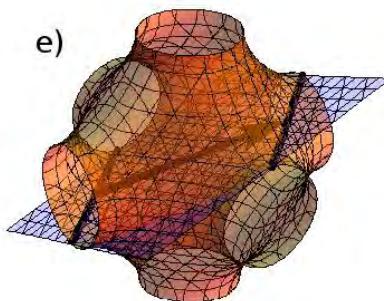
Conclusion

➤ DΓA for high T_c superconductors

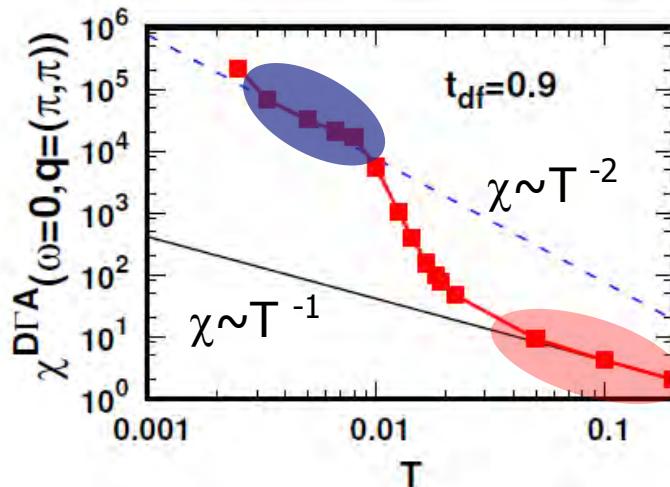
vertex dynamics suppresses T_c



➤ quantum criticality
3D Hubbard model



2D periodic Anderson model



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- **M. Kitatani, T. Schäfer, A. Toschi (TU Wien)**
- **A. Katanin (Ekaterinburg)**
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Further reading:

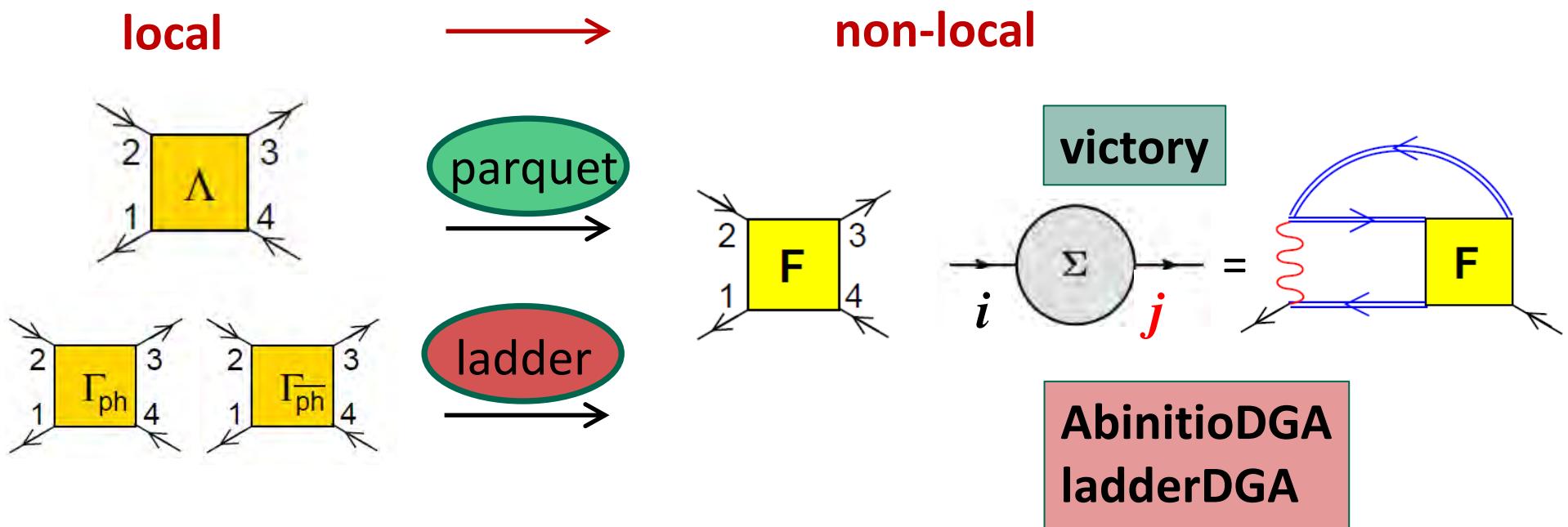
QCP Hubbard: PRL 119 0946402 (2017)
SC Hubbard: arXiv: 1801.05991
RMP 90, 025003 (2018)

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D Γ A: Resummation of Feynman diagrams in terms of locality
 n-particle fully irreducible vertex *approximated as local*



includes local **DMFT** correlations
 and **non-local** correlations
 (e.g. spin fluctuations)

