

Hedin equations, GW, GW+DMFT, and all That

Karsten Held* (TU Wien)

Jülich, Oct 7th, 2011

- 1 Introduction
- 2 Hedin Equations
- 3 GW
- 4 GW+DMFT
- 5 All of That: *ab initio* D Γ A

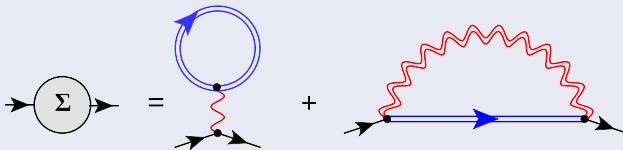
* with C. Taranto, G. Rohringer, and A. Toschi

1) Introduction

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Idea

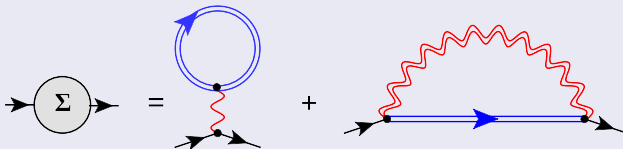
Alternative to DFT: Hedin's (1965) GW



1) Introduction

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Alternative to DFT: Hedin's (1965) GW



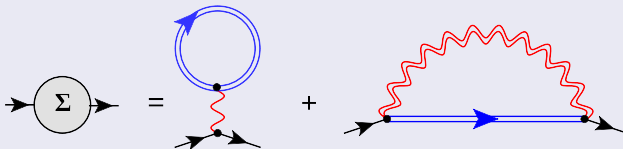
$$\Sigma^{\text{GW}}(\mathbf{r}, \mathbf{r}'; \omega) = i \int \frac{d\omega'}{2\pi} G(\mathbf{r}, \mathbf{r}'; \omega + \omega') W(\mathbf{r}, \mathbf{r}'; \omega').$$

Self energy: Fock-like term but with screened interaction W

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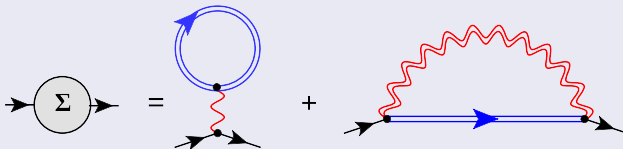
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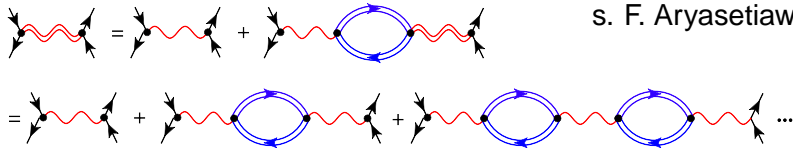
Self energy: Fock-like term but with screened interaction \mathbf{W}

- often similar as LDA
- better if non-local exchange is important
→ semiconductor band gaps

What kind of diagrams?

Screening within **random phase approx (RPA)**

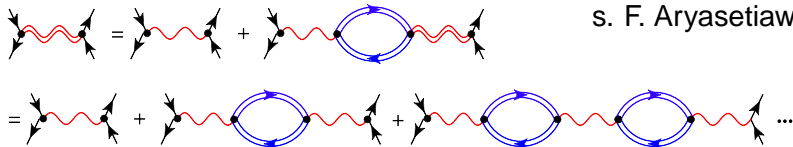
s. F. Aryasetiawan's lecture



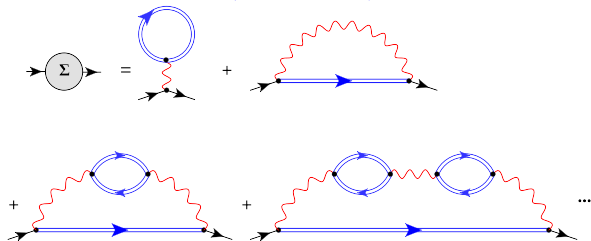
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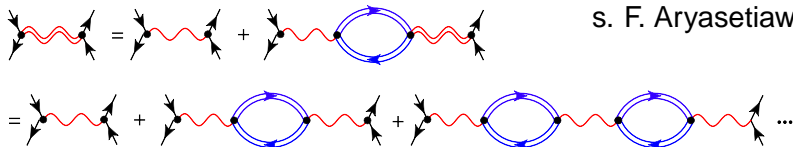
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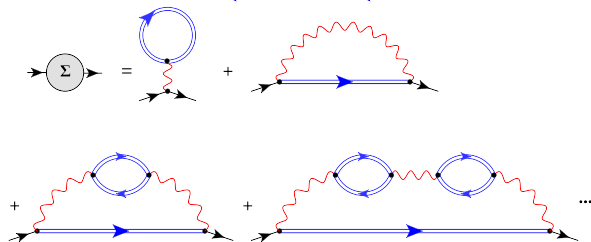
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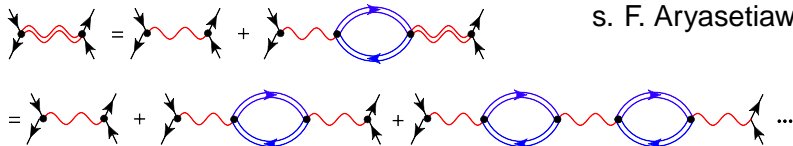


- GW exchange “in between” LDA and bare Hartree

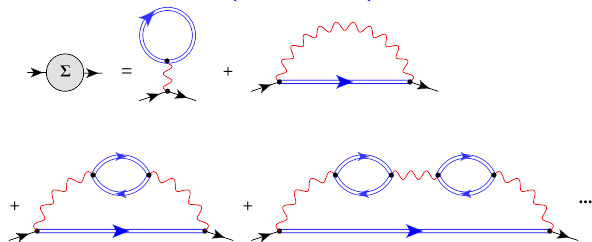
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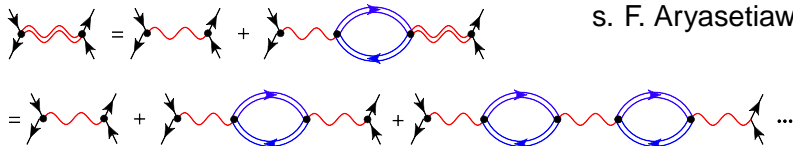


- GW exchange “in between” LDA and bare Hartree
- **good for exchange** cf. hybrid functionals

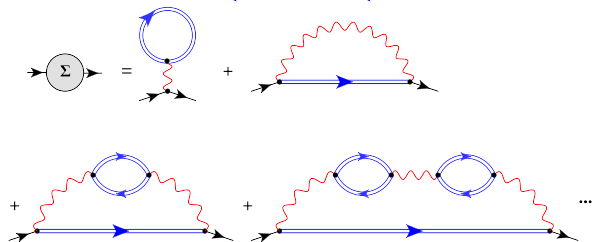
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This generates:



- GW exchange “in between” LDA and bare Hartree
- **good for exchange** cf. hybrid functionals
- **no strong correlations** (Hubbard bands) → **DMFT** or similar

2) Hedin Equations

GW is only the simplest approximation of something more general: [5 Hedin equations](#)

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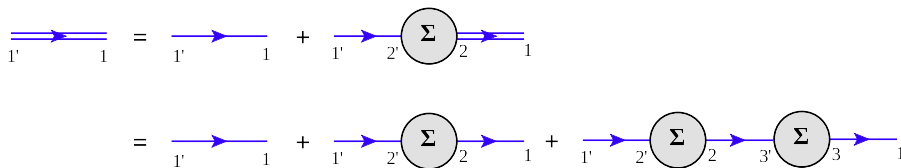
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Quote

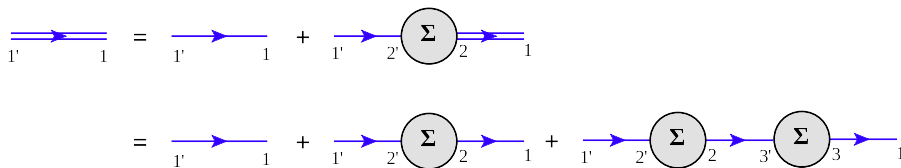
Hedin (1965): "The results [i.e., the Hedin equations] are well known to the Green function people"

1st Hedin equation: Dyson equation $G \leftrightarrow \Sigma$



$$\begin{aligned}
 G(11') &= G^0(11') + G(12)\beta\Sigma(22')G^0(2'1') \\
 &= G^0(11') + G^0(12)\beta\Sigma(22')G^0(2'1') + \\
 &\quad G^0(13)\beta\Sigma(33')G^0(3'2)\beta\Sigma(22')G^0(2'1') + \dots
 \end{aligned}$$

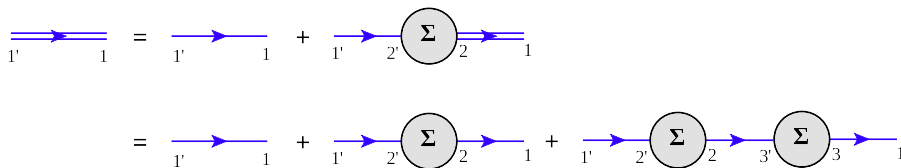
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Short hand notation $\mathbf{1}$ for $(\mathbf{r}_1, \tau_1, \sigma)$ or (i, m, ω, σ) (Bickers'04)

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All Feynman diagrams generated by connecting

1-p irreducible building blocks Σ by GF lines

1-p irreducible: cutting one GF \Rightarrow diagram still connected

2nd Hedin equation: screened interaction $W \leftrightarrow$ polarization op. P

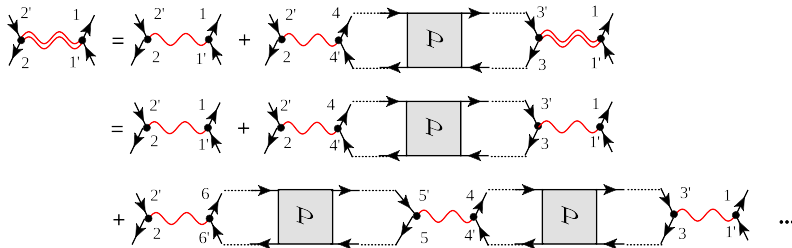
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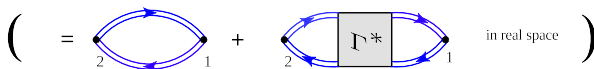
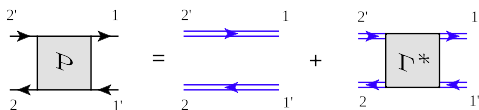
G : all Feynman diagrams connecting left and right by G^0 's



$$W(11'; 22') = V(11'; 22') + W(11'; 33') P(3'3; 4'4) V(44'; 22') .$$

P : all diagrams irreducible w.r.t. cutting one V

3rd Hedin equation: $P \leftrightarrow$ polarization op. Γ^*
 standard relation between 2-p GF (response fct.) and vertex



$$P(11'; 22') = \beta G(12')G(21') + \beta G(13)G(3'1')\Gamma^*(33'; 44')\beta G(4'2')G(24)$$

3rd Hedin equation: $P \leftrightarrow$ polarization op. Γ^*
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The diagram shows the 3rd Hedin equation for the polarization operator P and the vertex Γ^* . The first equation is:

$$P = G_0 G_0 + G_0 \Gamma^* G_0$$

where G_0 is the non-interacting Green's function. The second equation is:

$$\Gamma^* = \Gamma^* G_0 \Gamma^* + \Gamma^* G_0 \Gamma^* G_0 \Gamma^* + \dots$$

in real space.

$$P(11'; 22') = \beta G(12')G(21') + \beta G(13)G(3'1')\Gamma^*(33'; 44')\beta G(4'2')G(24)$$

- Note, P is irreducible w.r.t. cutting one V
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The diagram shows the 3rd Hedin equation for the polarization operator P and the vertex Γ^* . The top row shows the Dyson equation for P in terms of the irreducible polarization Γ^* and the two-particle Green's function G . The bottom row shows the same equation in real space, where the two-particle Green's function is represented by a loop of two G lines.

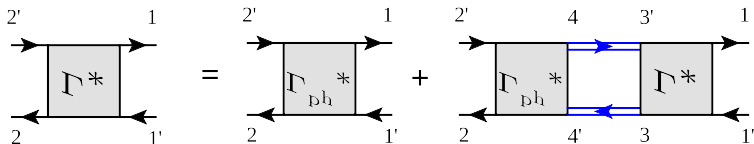
$$P = G G + \Gamma^* G G$$

(= + in real space)

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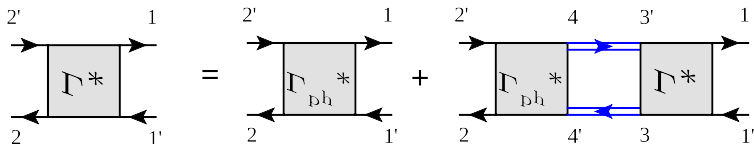
- Note, P is irreducible w.r.t. cutting one V
 $\Rightarrow \Gamma^*$ is irreducible w.r.t. cutting one V
 2nd line: simplifications in real space $2=2'$, $1=1'$

4th Hedin equation: Bethe-Salpeter eq. red. Γ^* \leftrightarrow irred. Γ_{ph}^*
 Analogon to **Dyson eq.** but for two-particles



$$\Gamma^*(11'; 22') = \Gamma_{ph}^*(11'; 22') + \Gamma^*(11'; 33') \beta G(3'4) G(4'3) \Gamma_{ph}^*(44'; 22')$$

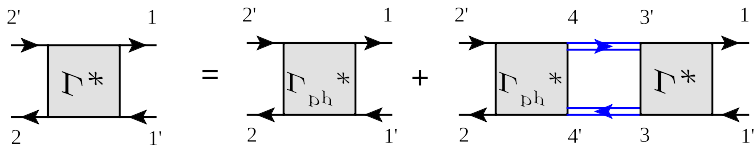
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$\Gamma_{ph}(44'; 22')$ irred. **particle-hole** vertex:
 irred. w.r.t. cutting a **particle** and a **hole** line

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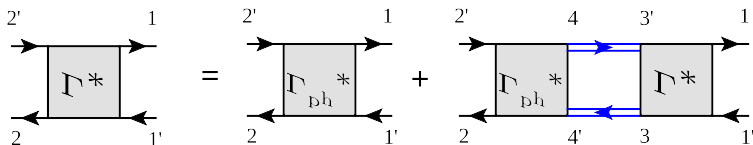
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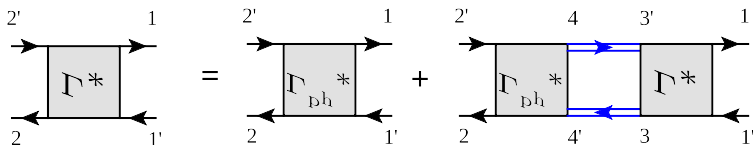
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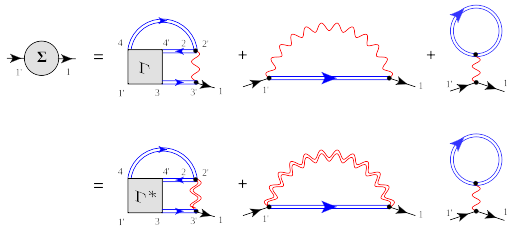
$$\text{Hedin: } \Gamma_{ph}^*(11'; 22') = \frac{\delta[\Sigma(11') - \Sigma_{\text{Hartree}}(11')]}{\delta G(2'2')}$$

5th Hedin equation: (Heisenberg) equation of motion

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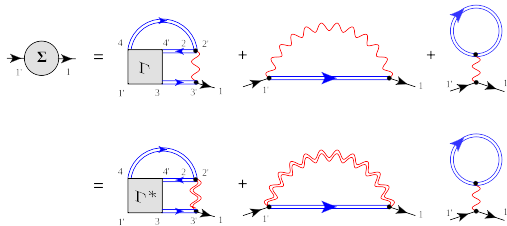


$$\begin{aligned} \Sigma(11') &= -V(13'; 22')\beta G(4'2)G(2'4)G(3'3)\Gamma(31'; 44') - V(12'; 21')G(2'2) \\ &= -W(13'; 22')\beta G(4'2)G(2'4)G(3'3)\Gamma^*(31'; 44') - W(12'; 21')G(2'2) \end{aligned}$$

for details s. Appendix of Lecture Notes

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From Hedin's eq. to GW

Complicated part is irreducible vertex

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Dyson equation yields

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Often not self-consistent $G_0 W_0$ – with LDA as starting point

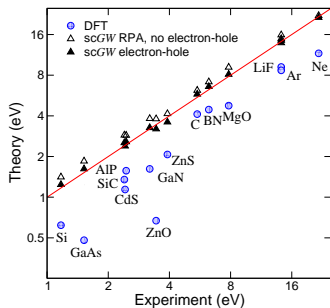
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self consistency with (Schilfgaarde, Kontani'04)

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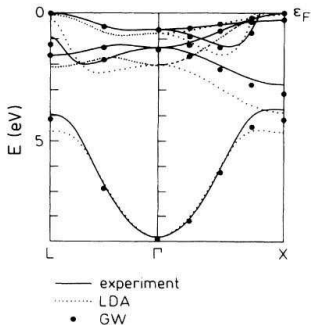
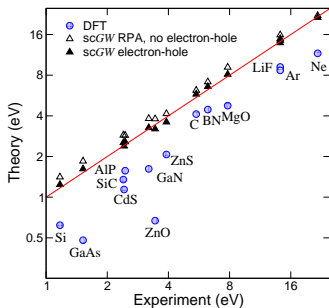
band gaps Shishkin et al.'07



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bandstructure of Ni

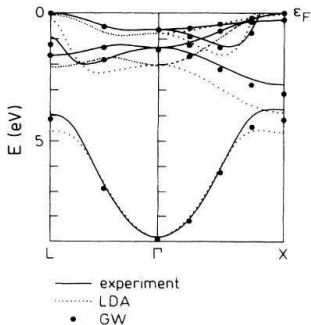
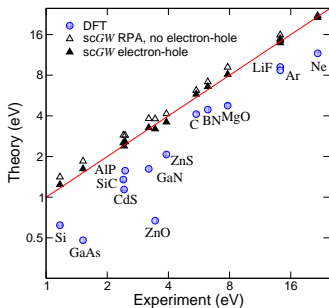
Aryasetiawan'82

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- no -6eV satellite \rightarrow *DMFT*

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Aryasetiawan'82

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- also finite life times e.g. Ag

4) GW+DMFT

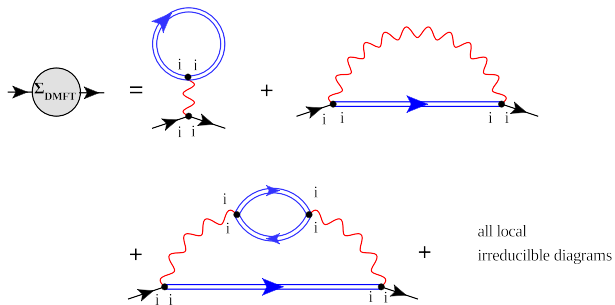
Merge GW and DMFT self energies (or energy functionals)

Biermann et al.'03

4) GW+DMFT

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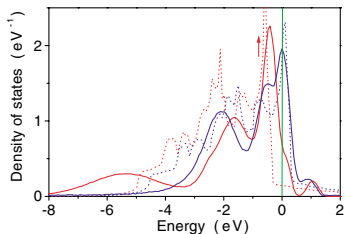


Algorithm

reproduced from
Held Adv. Phys. '07

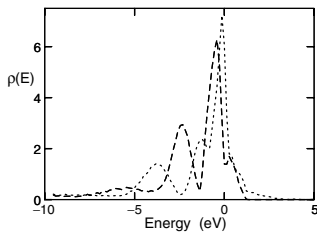
Do LDA calculation, yielding $\mathbf{G}_k(\omega) = [\omega\mathbf{1} + \mu\mathbf{1} - \epsilon^{\text{LDA}}(\mathbf{k})]^{-1}$.	
Calculate GW polarization $\mathbf{P}^{GW}(\omega) = -2i \int \frac{d\omega'}{2\pi} \mathbf{G}(\omega + \omega') \mathbf{G}(\omega')$.	
If DMFT polarization \mathbf{P}^{DMFT} is known (after the 1st iteration), include it $\mathbf{P}^{GW+\text{DMFT}}(\mathbf{k}, \omega) = \mathbf{P}^{GW}(\mathbf{k}, \omega) - \frac{1}{V_{\text{BZ}}} \int d^3k' \mathbf{P}^{GW}(\mathbf{k}, \omega) + \mathbf{P}^{\text{DMFT}}(\omega).$	
With this polarization, calculate the screened interaction: $\mathbf{W}(\mathbf{k}; \omega) = \mathbf{V}_{\text{cc}}(\mathbf{k}) [\mathbf{1} - \mathbf{V}_{\text{cc}}(\mathbf{k}) \mathbf{P}(\mathbf{k}; \omega)]^{-1}.$	
Calculate $\Sigma_{\mathbf{k}}^{\text{Hartree}} = \int \frac{d^3q}{V_{\text{BZ}}} \mathbf{G}_q(\tau=0^-) \mathbf{W}(\mathbf{k}-\mathbf{q}, 0)$ and $\Sigma_{\text{dc}}^{\text{Hartree}}$.	
Calculate $\Sigma^{GW}(\mathbf{r}, \mathbf{r}'; \omega) = i \int \frac{d\omega'}{2\pi} \mathbf{G}(\mathbf{r}, \mathbf{r}'; \omega + \omega') \mathbf{W}(\mathbf{r}, \mathbf{r}'; \omega')$.	
Calculate the DMFT self-energy Σ^{DMFT} and polarization \mathbf{P}^{DMFT} as follows:	
	From the local Green function \mathbf{G} and old self-energy Σ^{DMFT} calculate $(\mathcal{G}^0)^{-1}(\omega) = \mathbf{G}^{-1}(\omega) + \Sigma^{\text{DMFT}}(\omega) \quad \Sigma^{\text{DMFT}} = 0 \text{ in 1st iteration.}$
	Extract the local screening contributions from \mathbf{W} : $\mathbf{U}(\omega) = [\mathbf{W}^{-1}(\omega) - \mathbf{P}^{\text{DMFT}}(\omega)]^{-1}.$
	With \mathbf{U} and \mathcal{G}^0 , solve impurity problem with effective action $\mathcal{A} = \sum_{\nu\sigma} \psi_{\nu m}^{\sigma*} (\mathcal{G}_{\nu m n}^{\sigma 0})^{-1} \psi_{\nu n}^{\sigma} + \sum_{l m \sigma \sigma'} \int d\tau \psi_l^{\sigma*}(\tau) \psi_l^{\sigma}(\tau) U_{lm}(\tau - \tau') \psi_m^{\sigma'*}(\tau') \psi_m^{\sigma'}(\tau'),$ resulting in \mathbf{G} and susceptibility χ .
	From \mathbf{G} and χ , calculate $\Sigma^{\text{DMFT}}(\omega) = (\mathcal{G}^0)^{-1}(\omega) - \mathbf{G}^{-1}(\omega)$, $\mathbf{P}^{\text{DMFT}}(\omega) = \mathbf{U}^{-1}(\omega) - [\mathbf{U} - \mathbf{U}\chi\mathbf{U}]^{-1}(\omega).$
Combine this to the total GW self-energy: $\Sigma^{GW+\text{DMFT}}(\mathbf{k}, \omega) = \Sigma^{GW}(\mathbf{k}, \omega) - \int d^3k' \Sigma^{GW}(\mathbf{k}, \omega) + \Sigma_{\text{dc}}^{\text{Hartree}}(\mathbf{k}) - \Sigma_{\text{dc}}^{\text{Hartree}} + \Sigma^{\text{DMFT}}(\omega).$	
From this and \mathcal{G}^0 , calculate $\mathbf{G}_k^{\text{new}}(\omega)^{-1} = \mathcal{G}_k^0(\omega)^{-1} - \Sigma_k(\omega)$.	
Iterate with $\mathbf{G}_k = \mathbf{G}_k^{\text{new}}$ until convergence, i.e. $\ \mathbf{G}_k - \mathbf{G}_k^{\text{new}}\ < \epsilon$.	

Results for Ni



LDA+DMFT

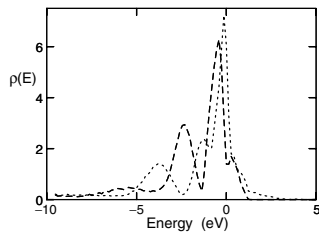
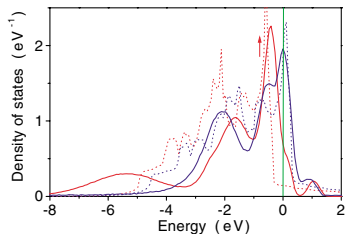
Lichtenstein et al.'01



GW+DMFT

Biermann et al.'03

Results for Ni



LDA+DMFT

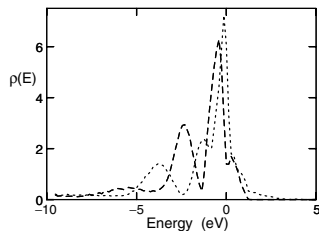
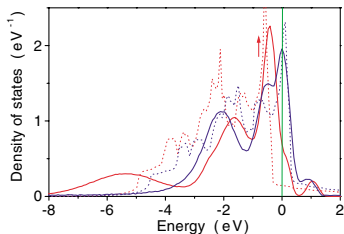
Lichtenstein et al.'01

- satellite at -6eV; both similar

GW+DMFT

Biermann et al.'03

Results for Ni



LDA+DMFT

Lichtenstein et al.'01

- satellite at -6eV; both similar
- no self consistency, P^{DMFT} ...
- $W(\omega) \rightarrow W$

GW+DMFT

Biermann et al.'03

4) All of That: *ab initio* D Γ A

Hedin's eq.: combine *GW* exchange and correlations on 2p level
instead of 1p *GW*+DMFT

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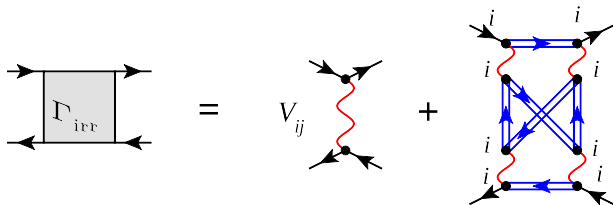
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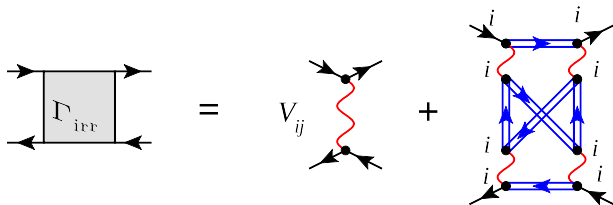
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→ starting point of Hedin eq.